

Conformal Symmetry and the Standard Model

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Hierarchy Problem

- **Fact: Standard Model** (= SM) of elementary particle physics is **conformally invariant** **except** for tree level mass term $-m^2\Phi^\dagger\Phi$ \rightarrow masses for vector bosons, quarks and leptons (via *Brout-Englert-Higgs* mechanism).
- Quantum corrections $m^2 \sim \Lambda^2 \Rightarrow$ why $m_H \ll M_P$? (with UV cutoff $\Lambda =$ scale of ‘new physics’)
 - *explanation of hierarchy?*
 - *stabilization of hierarchy?*
- Most popular proposal: SM \rightarrow MSSM or NMSSM: use supersymmetry to control quantum corrections via cancellation of quadratic divergences \Rightarrow

$$m^2 \sim \Lambda_{SUSY}^2 \ln(\Lambda^2/\Lambda_{SUSY}^2)$$

Coleman-Weinberg Mechanism (1973)

- Idea: spontaneous breaking of conformal symmetry by quantum corrections \implies small mass scales arise via *conformal anomaly* and *effective potential*

$$\lambda\varphi^4 \rightarrow V_{\text{eff}}(\varphi) = \lambda\varphi^4 + \frac{\lambda^2\varphi^4}{64\pi^2} \left(\ln \left(\frac{\lambda\varphi^2}{v^2} \right) - \frac{1}{2} \right)$$

- In its original form this proposal does not work:
 - Higgs mass too small ($\sim \mathcal{O}(10 \text{ GeV})$), or
 - Scalar self-couplings too large \rightarrow Landau poles

[See e.g.: Sher, Phys.Rep.179(1989)273; Ford, Jones, Stephenson, Einhorn, Nucl.Phys.B395(1993)17]

- Also, must accommodate:
 - $m_H > 115 \text{ GeV}$ and $m_{top} = 174 \text{ GeV}$
 - $m_\nu < 1 \text{ eV}$ \rightarrow *large intermediate scale?*

Our Proposal [= hep-th/0612165]

- **Classical conformal symmetry** (i.e. no tree level mass terms) in SM
- plus right-chiral neutrinos
- plus enlarged scalar sector: Φ and φ
- All mass scales from effective (CW) potential
- all coupling constants small and positive up to M_{Pl}
- No large intermediate scales \implies
 - no grand unification (GUTs)
 - no new scales required to explain $m_\nu < 1$ eV
- No low energy SUSY

The Model

- Start from conformally invariant (and therefore renormalizable) Lagrangian $\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}'$ with:

$$\begin{aligned} \mathcal{L}' := & \left(\bar{L}^i \Phi Y_{ij}^E E^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^D D^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^U U^j + \right. \\ & \left. + \bar{L}^i \epsilon \Phi^* Y_{ij}^\nu \nu_R^j + \varphi \nu_R^{iT} \mathcal{C} Y_{ij}^M \nu_R^j + \text{h.c.} \right) - \\ & - \frac{\lambda_1}{4} (\Phi^\dagger \Phi)^2 - \frac{\lambda_2}{2} \varphi^2 (\Phi^\dagger \Phi) - \frac{\lambda_3}{4} \varphi^4 \end{aligned}$$

- Besides usual $SU(2)$ doublet Φ : new scalar field φ (here taken to be *real*, but could be complex and in non-trivial representation of family symmetry)
- No mass terms, all coupling constants dimensionless
- $Y_{ij}^U, Y_{ij}^E, Y_{ij}^M$ real and diagonal
 Y_{ij}^D, Y_{ij}^ν complex \rightarrow parametrize family mixing

Effective Potential

$$\begin{aligned} V_{\text{eff}}(H, \varphi) = & \frac{\lambda_1 H^4}{4} + \frac{\lambda_2 H^2 \varphi^2}{2} + \frac{\lambda_3 \varphi^4}{4} \\ & + \frac{3}{256\pi^2} (\lambda_1 H^2 + \lambda_2 \varphi^2)^2 \ln \left[\frac{\lambda_1 H^2 + \lambda_2 \varphi^2}{v^2} \right] \\ & + \frac{1}{64\pi^2} F_+^2 \ln \left[\frac{F_+}{v^2} \right] + \frac{1}{64\pi^2} F_-^2 \ln \left[\frac{F_-}{v^2} \right] \\ & - \frac{6}{32\pi^2} g_t^4 (H^2)^2 \ln \left[\frac{H^2}{v^2} \right] - \frac{1}{32\pi^2} g_M^4 \varphi^4 \ln \left[\frac{\varphi^2}{v^2} \right] \end{aligned}$$

$g_M^4 \equiv \text{Tr } Y_M^4$, v dimensionful scale (**conformal anomaly**)
Not included: contributions from $SU(2)_w \times U(1)_Y$ gauge fields (because respective gauge couplings are small), nor from $SU(3)_c$ (because it is a two-loop effect)

Search for minima: must perform *numerical analysis!*

Numerics

Choice of parameters is *strongly constrained* by experimental data and RGE analysis \rightarrow ‘trial and error’ method leads to following choice of parameters

$$\lambda_1 = 3.4, \quad \lambda_2 = 2.6, \quad \lambda_3 = 3.3, \quad g_t = 1, \quad g_M^2 = 0.4$$

Minimum lies at

$$\langle H \rangle = 0.415 \cdot 10^{-5} v, \quad \langle \varphi \rangle = 2.506 \cdot 10^{-5} v$$

Normalize this by setting $\langle H \rangle = 174 \text{ GeV} \Rightarrow$

$$H' = H \cos \beta + \varphi \sin \beta, \quad \varphi' = -H \sin \beta + \varphi \cos \beta$$

$$m_{H'} = 217 \text{ GeV}, \quad m_{\varphi'} = 439 \text{ GeV}; \quad \sin \beta = 0.119$$

‘Higgs mixing’: only the components along H of the mass eigenstates couple to the usual SM particles.

NB: Not (yet?) a definitive prediction.

Renormalization Group Equations

Effective couplings (\equiv 4th derivatives at minimum):

$$\lambda_1^{\text{eff}} = 1.463, \quad \lambda_2^{\text{eff}} = 0.348, \quad \lambda_3^{\text{eff}} = 0.626$$

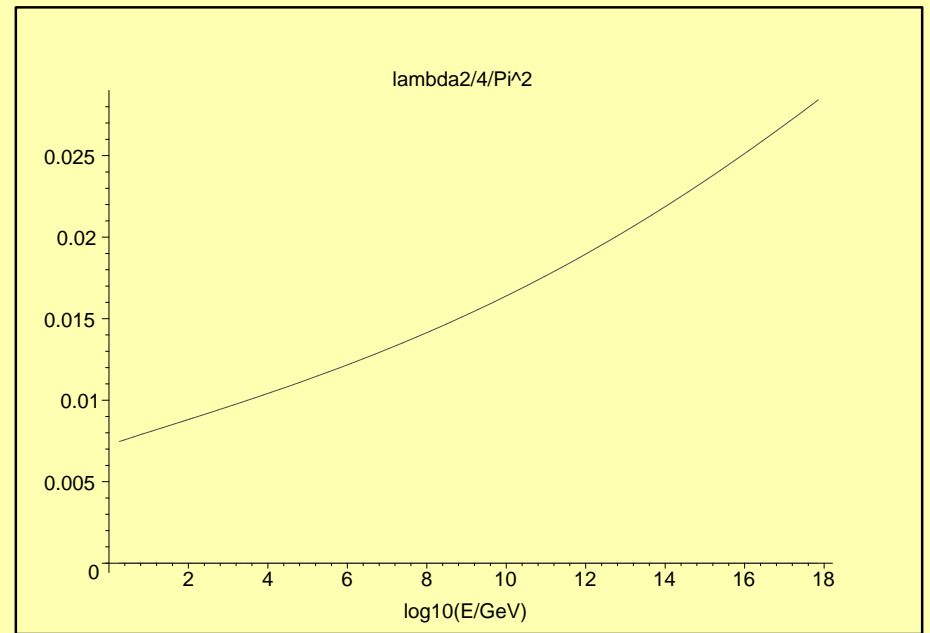
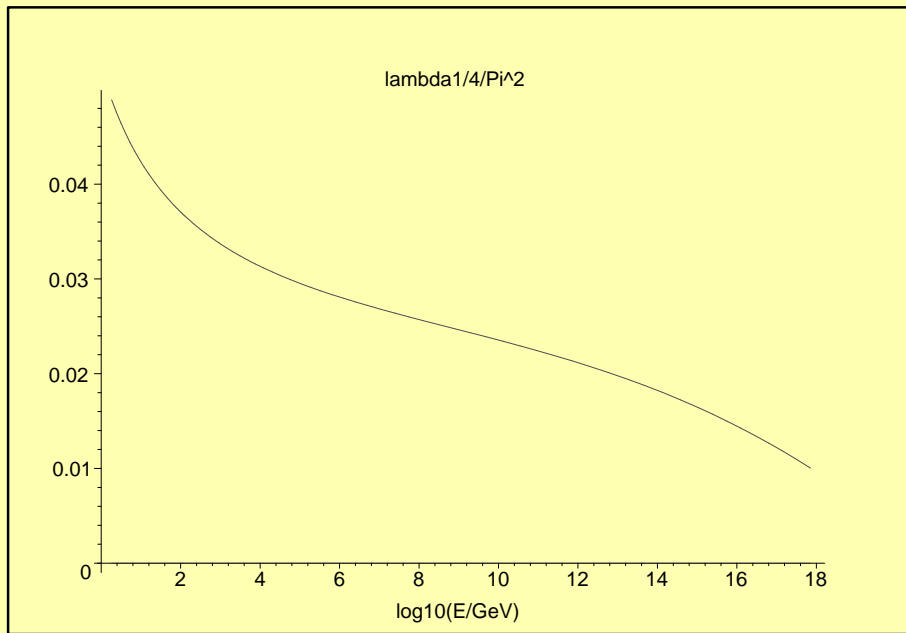
With

$$y_1 = \frac{\lambda_1^{\text{eff}}}{4\pi^2}, \quad y_2 = \frac{\lambda_2^{\text{eff}}}{4\pi^2}, \quad y_3 = \frac{\lambda_3^{\text{eff}}}{4\pi^2}, \quad x = \frac{g_t^2}{4\pi^2}, \quad u = \frac{g_M^2}{4\pi^2}, \quad z = \frac{\alpha_s}{\pi}$$

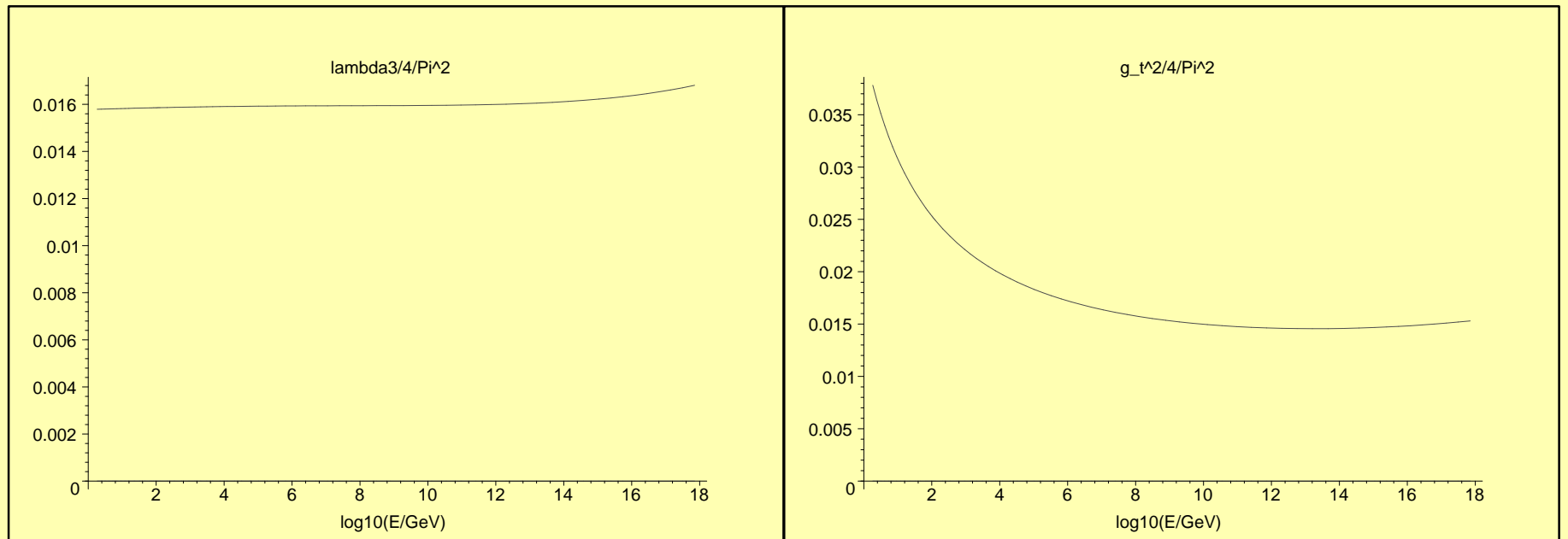
we get

$$\begin{aligned} \mu \frac{dy_1}{d\mu} &= \frac{3}{2}y_1^2 + \frac{1}{8}y_2^2 - 6x^2, \\ \mu \frac{dy_2}{d\mu} &= \frac{3}{8}y_2(2y_1 + y_3 + \frac{4}{3}y_2), \\ \mu \frac{dy_3}{d\mu} &= \frac{9}{8}y_3^2 + \frac{1}{2}y_2^2 - u^2, \quad \mu \frac{du}{d\mu} = \frac{3}{4}u^2 \\ \mu \frac{dx}{d\mu} &= \frac{9}{4}x^2 - 4xz, \quad \mu \frac{dz}{d\mu} = -\frac{7}{2}z^2. \end{aligned}$$

Evolution of Coupling Constants

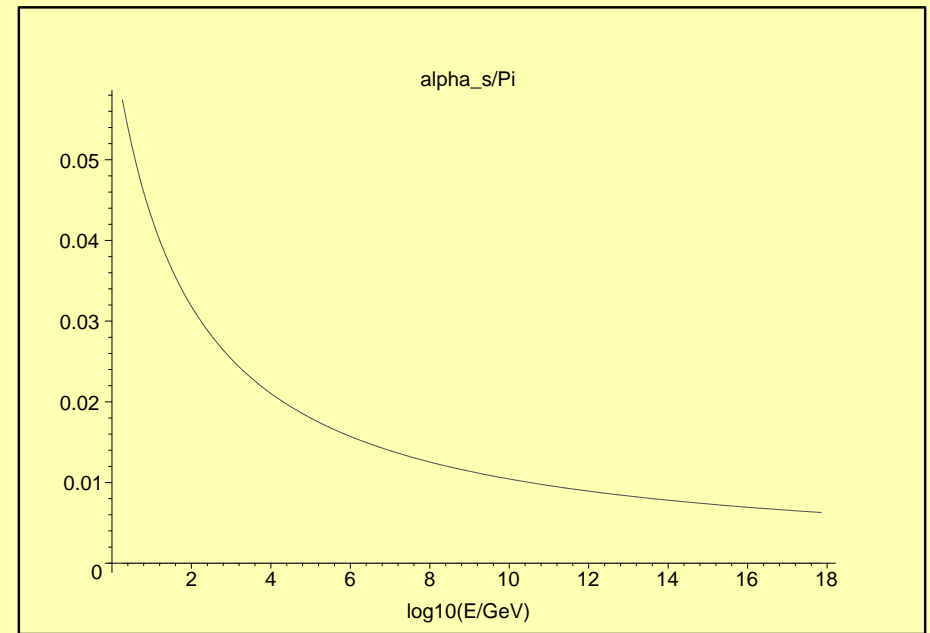
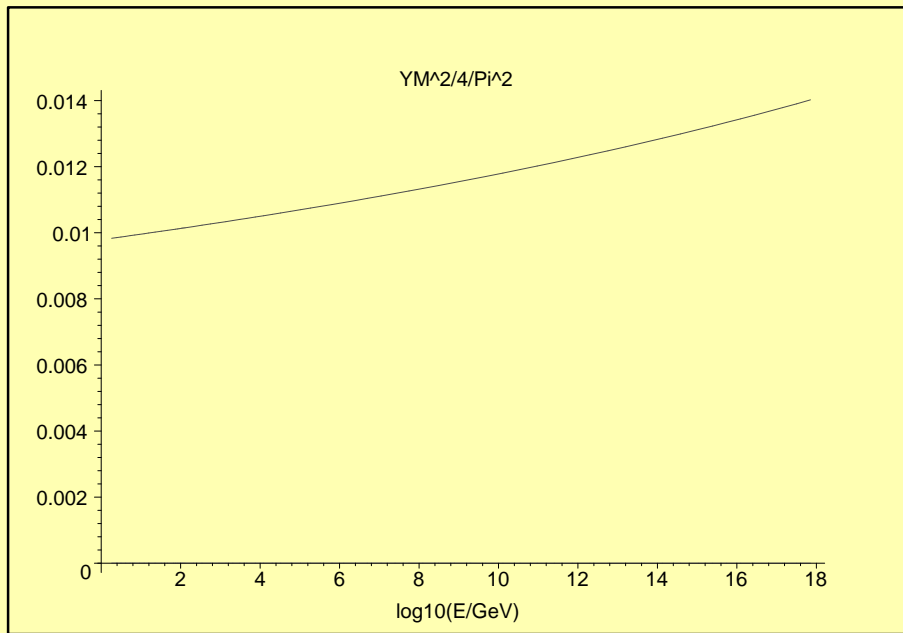


Evolution of Coupling Constants



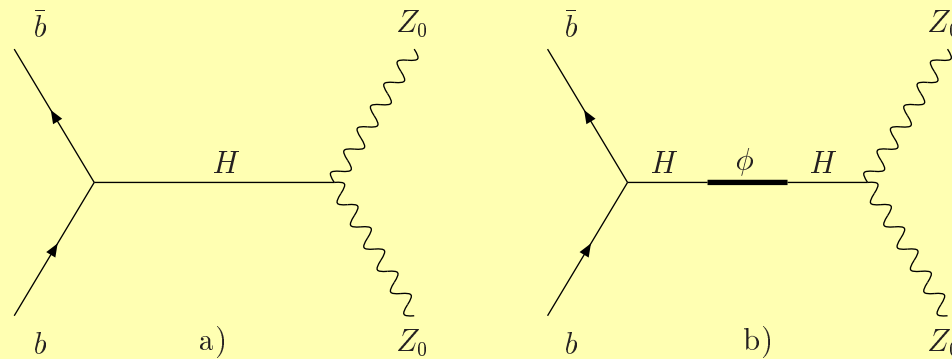
Important: top Yukawa stays bounded up to M_{Pl} due to large α_s !

Evolution of Coupling Constants



Phenomenology

- New scalar is a 'heavier brother' of the SM Higgs (the same BRs except for the mass difference and lower cross sections) **unique signature**



- if $m_\phi > 2m_Z$ a clear signal in LHC
- no other new particles at LHC except standard Higgs and a new scalar (scalars?)
- if $|\sin(\beta)| < 0.1$ then no contradiction with LEP precision data

Discussion and outlook

- conformal symmetry is a different (much simpler than SUSY...) explanation of hierarchy of scales
- SM with massive neutrinos and the new scalar can be viable up to the Planck scale so LHC may see just the SM Higgs and (several?) new scalar particles...
- LHC will tell...