Hunting SUSY: Dark Matter at the LHC

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Outline



Introduction

- Evidence for Dark Matter
- Calculating the Density
- Dark Matter and the LHC
- 2 SUSY Dark Matter• SUSY at the LHC
 - **3** EW Scale MSSM Parameters
 - 4 GUT Scale Parameters
- 5 Conclusions and open questions

There are three primary categories of evidence for dark matter:

Rotation Curves of Galaxies

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- Gravitational Lensing

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- Gravitational Lensing
- The Cosmic Microwave Background

Rotation Curves of Galaxies



(From Klypin, Zhao and Somerville, 2002)

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Lensing: Dark Matter in the Bullet Cluster



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The Cosmic Microwave Background



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Neutral

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Current Limits

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$$\Omega_{CDM} h^2 = 0.1126 \pm 0.0081.$$

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In contrast, the density of baryonic matter is measured as:

$$\Omega_b h^2 = 0.0224$$

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- As the universe cools, the dark matter falls out of equilibrium.
- It's annihilation rate during this time determines the resulting relic abundance.
- The relic density can be approximated as:

$$\Omega_{CDM} h^2 \approx \frac{3 \times 10^{-27} \mathrm{cm}^{3} \mathrm{s}^{-1}}{<\sigma \nu >}$$

What has this got to do with the LHC?

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Particle dark matter should have a mass within the reach of the LHC.

Exotic matter produced at the LHC will decay to the lightest stable particle - a dark matter candidate.

Therefore the LHC should be a dark matter factory.

The LHC is critical to the study of dark matter.

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Dark matter allows us to differentiate theories with the same mass spectrum.

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- What precision does this then give for the relic density calculation?

Candidates

- Neutralinos
- Kaluza-Klein excitations
- Massive neutrinos
- Axions
- Sneutrinos
- Axinos
- Gravitinos
- Q-balls
- Mirror particles
- WIMPZILLAS
- Shadow world matter
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- Cryptons
- ... but here we stick to candidates from Supersymmetry.

The MSSM Particle Spectrum

Standard Model

quarks q

leptons /

gauge bosons $A^{(a)}_{\mu}$

Higgs boson H

The MSSM Particle Spectrum



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R-parity

The MSSM contains a discrete symmetry called R-parity.

$$P_R = (-1)^{3(B-L)+2S}$$

This is required to prevent fast proton decay.

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Under R-parity, Standard Model and supersymmetric matter have the following assignments:

Due to R-parity, decays of the form:

sparticle \rightarrow particle + particle

are not allowed.

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R-parity conserving supersymmetry requires a relic density of sparticles.

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(M ₁	0	$-m_Z \cos\beta\sin\theta_w$	$m_Z \sin \beta \sin \theta_w$
0	M_2	$m_Z \cos \beta \cos \theta_w$	$-m_Z \sin\beta\cos\theta_w$
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• β is related to the ratio of the Higgs VeVs, tan β .

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Providing a candidate is not enough. It must occur in the right abundance.

Bino dark matter primarily annihilates via the process:



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Image: A math a math

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The cross-section is generally far too small, giving $\Omega_{CDM}h^2 \gg \Omega_{CDM}^{WMAP}h^2$.

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The cross-section is generally far too small, giving $\Omega_{CDM}h^2 \gg \Omega_{CDM}^{WMAP}h^2$. Such models are ruled out as they would overclose the universe (which is bad...).

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If we have a Wino LSP, $M_2 \ll M_1, \mu$,

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If we have a Wino LSP, $M_2 \ll M_1, \mu$, or a higgsino LSP, $\mu \ll M_1, M_2$,

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In general such processes are extremely efficient resulting in $\Omega_{CDM} h^2 \ll \Omega_{CDM}^{WMAP} h^2.$

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SUSY has a problem

• Bino Dark Matter: Generally gives $\Omega_{CDM}h^2 \gg \Omega_{CDM}^{WMAP}h^2$

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Therefore SUSY naturally gives the wrong dark matter density!

To account for the observed dark matter density, we need different annihilation channels.

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To account for the observed dark matter density, we need different annihilation channels. Firstly if:

 $2m_{\tilde{\chi}_1^0} \approx m_{h^0,H^0,A^0,Z^0}$

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In this case, annihilation is usually **too efficient**: $\Omega_{CDM}h^2 \ll \Omega_{CDM}^{WMAP}h^2$ except on the edges of the resonance.

Coannihilation

A second exception is if there is another sparticle close in mass to the LSP, then we must adapt our calculation of the dark matter density as there will be a significant number density of these particles at freeze-out.

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In this case we must take into account processes of the form:



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An example - the CMSSM parameter space



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The neutralino will be seen as missing energy.
We always need to know $m_{\widetilde{\chi}_1^0}$.

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• Wino/Higgsino dark matter

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• $m_{\tilde{\chi}_1^+}$, $(m_{\tilde{\chi}_3^0}, m_{\tilde{\chi}_2^+})$

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 - Composition of $\tilde{\chi}_1^0$.
 - $m_{\tilde{\chi}^0_2}$
 - $m_{\tilde{\chi}_{1}^{+}}^{*}, (m_{\tilde{\chi}_{2}^{0}}, m_{\tilde{\chi}_{2}^{+}})$
- Resonant annihilation

•
$$\Delta m = m_{Z^0, h^0, H^0, A^0} - 2m_{\tilde{\chi}_1^0}$$

The MSSM has roughly 128 different parameters. How much progress can we make?

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But...

The other parameters determine the rest of the spectrum and therefore, much of the LHC phenomenology.

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We need to quantify the sensitivity:

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Sensitivity

We need to quantify the sensitivity:

$$\Delta_{a}^{\Omega} = \left| \frac{\partial \ln \left(\Omega_{CDM} h^{2} \right)}{\partial \ln \left(a \right)} \right|$$

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Sensitivity

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$$\Delta_{a}^{\Omega} = \left| \frac{\partial \ln \left(\Omega_{CDM} h^{2} \right)}{\partial \ln \left(a \right)} \right|$$

With these tools we can quantify the sensitivity of $\Omega_{CDM}h^2$ to SUSY inputs.

For a typical Bino dark matter region, the sensitivity to low energy masses is:

Parameter a	Δ^{Ω}_{a}
$m_{ ilde{ au}, ilde{ extbf{e}}_R, ilde{\mu}_R}$	0.6
$m_{ ilde{\chi}_1^0}$	0.09

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A caveat: here we only consider sensitivity to the masses, not the compositions. This decay channel cares about the composition of the lightest $\tilde{\tau}$, $\tilde{\mu}_R$ and \tilde{e}_R .

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Resonant annihilation

Now consider a typical A^0 resonance.

Parameter a	Δ^{Ω}_{a}
m _A	49
$m_{ ilde{\chi}_1^0}$	48

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Therefore to calculate the relic density with a precision of $\mathcal{O}(1\%)$, we would need to know the mass of the A^0 and $\tilde{\chi}_1^0$ to $\mathcal{O}(0.02\%)$.



In the case of coannihilation we get:

Parameter a	Δ^{Ω}_{a}
$m_{ ilde{ au}}$	41
$m_{{ ilde \chi}^0_1}$	24

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Once again we need $\mathcal{O}(0.02\%)$ precision to achieve a $\mathcal{O}(1\%)$ precision in $\Omega_{CDM}h^2$.

 There are a finite number of ways SUSY dark matter can account for the observed relic density.

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Can the LHC(ILC) achieve the precision required to compete with WMAP in these regions?

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The low energy masses are not believed to be independent. The GUT scale mass structure is determined by the mechanism of SUSY breaking.

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But no-one knows what this is yet.

Hunting SUSY: Dark Matter at the LHC

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So what changes?

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- The LHC phenomenology becomes more tractable.

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Hunting SUSY: Dark Matter at the LHC

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 - - Only the lightest mass particles are specified.

Open Questions

Jonathan Roberts IFT, University of Warsaw

Hunting SUSY: Dark Matter at the LHC

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In terms of the low energy MSSM:

• Study the LHC phenomenology in the WMAP regions to assess the LHC precision and understand the important signals.