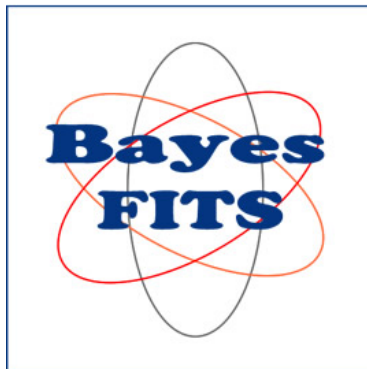


Implications of CMS searches for the Constrained MSSM - A Bayesian approach



Y-L Sming Tsai

BayesFITS Group

(A. Fowlie, K. Kowalska, S. Munir, L. Roszkowski, E. Sessolo,
S. Trojanowski

and

A. Kalinowski, M. Kazana, K. Nawrocki, P. Zalewski)

National Center for Nuclear Research (NCBJ)

Outline

1. Constrained MSSM (CMSSM)
2. Bayesian statistics in a nutshell
3. Impact of CMS α_T 1.1/fb and XENON100 limits
4. Summary

Fowlie, Kalinowski, Kazana, Roszkowski, Tsai (arXiv:1111.6098)

Roszkowski, Sessolo, Tsai (arXiv:1202.1503)

Constrained Minimal Supersymmetric Standard Model

G. L. Kane, C. F. Kolda, L. Roszkowski and J. D. Wells, Phys. Rev. D 49 (1994) 6173

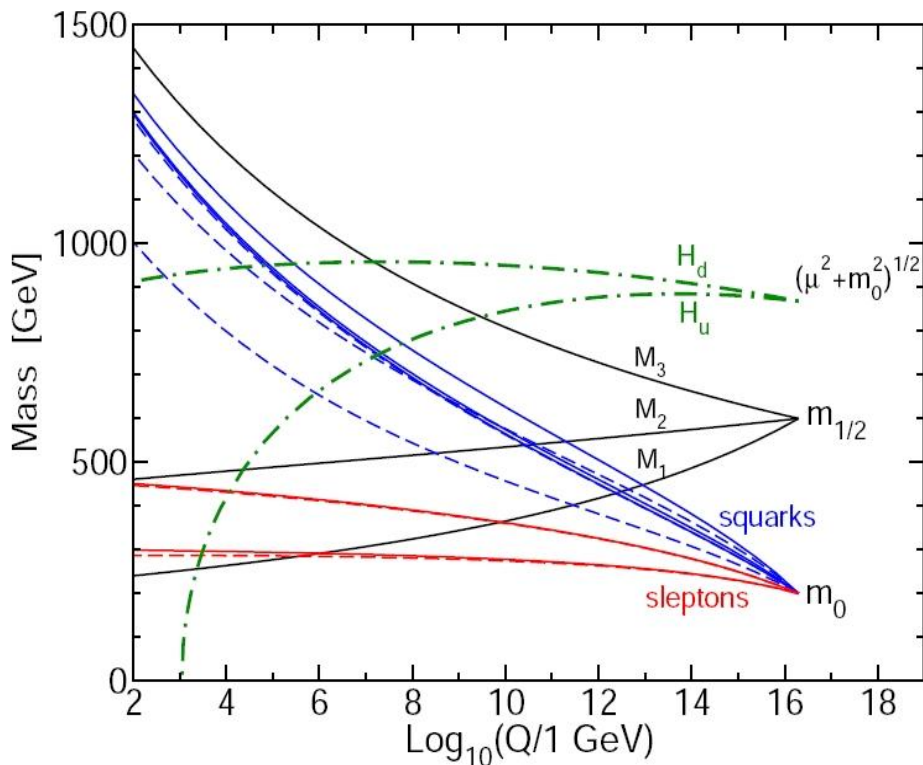


figure from hep-ph/9709356

At $M_{\text{GUT}} \simeq 2 \times 10^{16}$ GeV:

- gauginos $M_1 = M_2 = m_{\tilde{g}} = m_{1/2}$
- scalars $m_{\tilde{q}_i}^2 = m_{\tilde{l}_i}^2 = m_{H_b}^2 = m_{H_t}^2 = m_0^2$
- 3-linear soft terms $A_b = A_t = A_0$
- radiative EWSB
$$\mu^2 = \frac{m_{H_b}^2 - m_{H_t}^2 \tan^2 \beta}{\tan^2 \beta - 1} - \frac{m_Z^2}{2}$$
- five independent parameters: $m_{1/2}, m_0, A_0, \tan \beta, \text{sgn}(\mu)$
- well developed machinery to compute masses and couplings
- neutralino χ mostly bino

Comparison of experimental data and model prediction

Traditional chi-square method:

$$\chi^2 = \sum \frac{(\text{prediction} - \text{measurement})^2}{(\text{error})^2}.$$

As long as experimental central value and error are given, one can use chi-square method to compare the experimental data and model predictions.

Alternatively...

Bayesian Statistics...

Bayes theorem:

$$\text{Posterior} = \frac{\text{Prior} \times \text{Likelihood}}{\text{Evidence}}.$$

- **Likelihood**: the probability of obtaining data if hypothesis is true.
 - **Prior**: what we know about hypothesis BEFORE seeing the data.
 - **Evidence**: normalization constant, crucial for model comparison.
 - **Posterior**: the probability about hypothesis AFTER seeing the data.
-

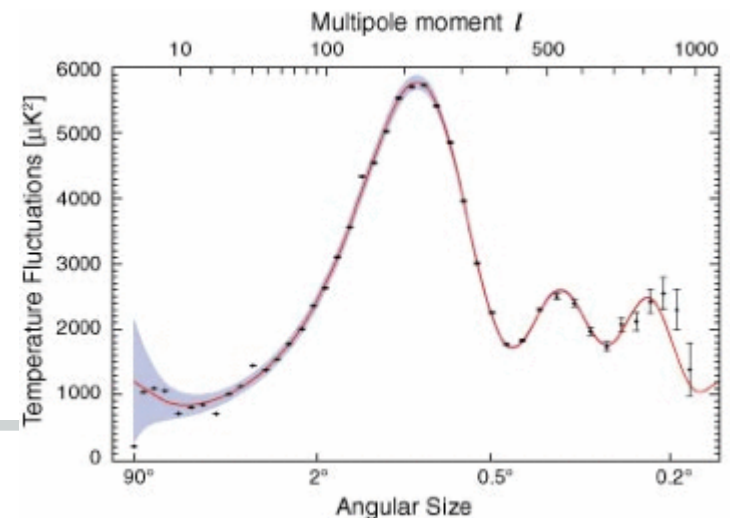
There is no single, ``right'' statistics...

Frequentist: "probability is the number of times the event occurs over the total number of trials, in the limit of an infinite series of equiprobable repetitions"

Bayesian: "probability is a measure of the degree of belief about a proposition"

Bayesian statistics is very popular in many branches of science (astronomy, cosmology, etc.).

For example, The Wilkinson Microwave Anisotropy Probe (WMAP) analysis of cosmic microwave background (CMB) spectrum:

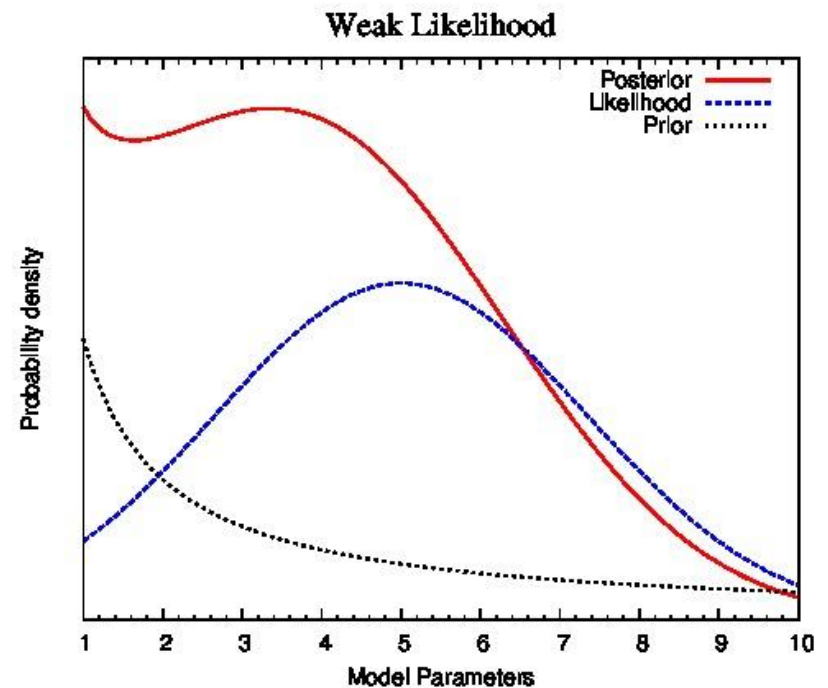
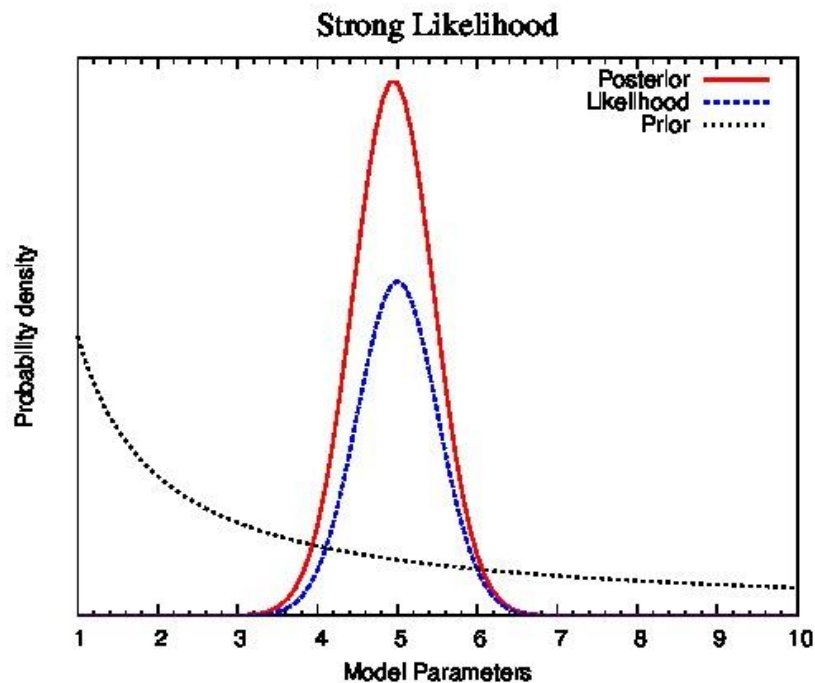


Prior dependence

Prior dependence is two-fold:

- Prior range.
- Prior distribution.

$$\text{Posterior} = \frac{\text{Prior} \times \text{Likelihood}}{\text{Evidence}}$$



- If the Likelihood is well-peaked, the posterior follows the Likelihood
- Otherwise, it follows the prior.

$$\text{Posterior} = \frac{\text{Prior} \times \text{Likelihood}}{\text{Evidence}}$$

Gaussian Likelihood

Take a single observable $\xi(m)$ that has been measured

(e.g., M_W)

• c – central value, σ – standard exptal error

• define

$$\chi^2 = \frac{[\xi(m) - c]^2}{\sigma^2}$$

• assuming Gaussian distribution ($d \rightarrow (c, \sigma)$):

$$\mathcal{L} = p(\sigma, c | \xi(m)) = \frac{1}{\sqrt{2\pi}\sigma} \exp\left[-\frac{\chi^2}{2}\right]$$

• when include theoretical error estimate τ (assumed Gaussian):

$$\sigma \rightarrow s = \sqrt{\sigma^2 + \tau^2}$$

TH error “smears out” the EXPTAL range

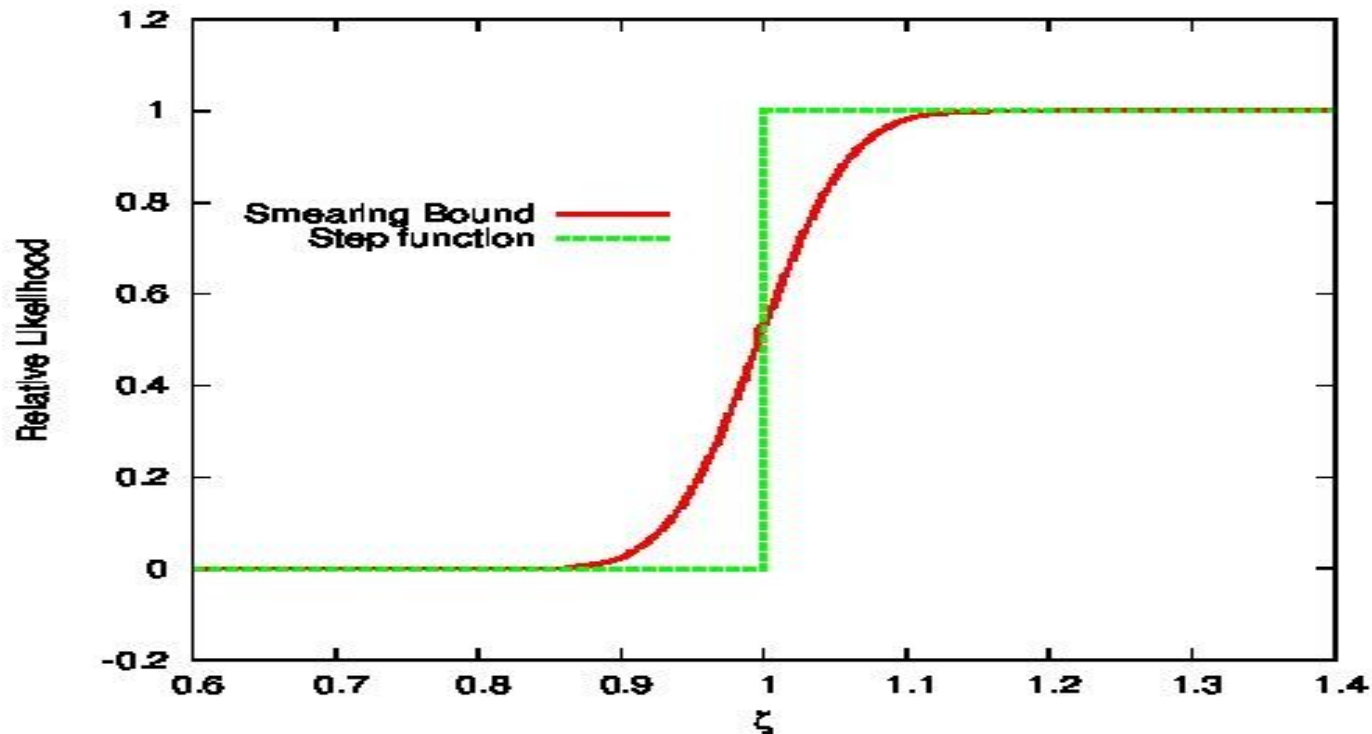
• for several uncorrelated observables (assumed Gaussian):

$$\mathcal{L} = \exp\left[-\sum_i \frac{\chi_i^2}{2}\right]$$

$$\text{Posterior} = \frac{\text{Prior} \times \text{Likelihood}}{\text{Evidence}}$$

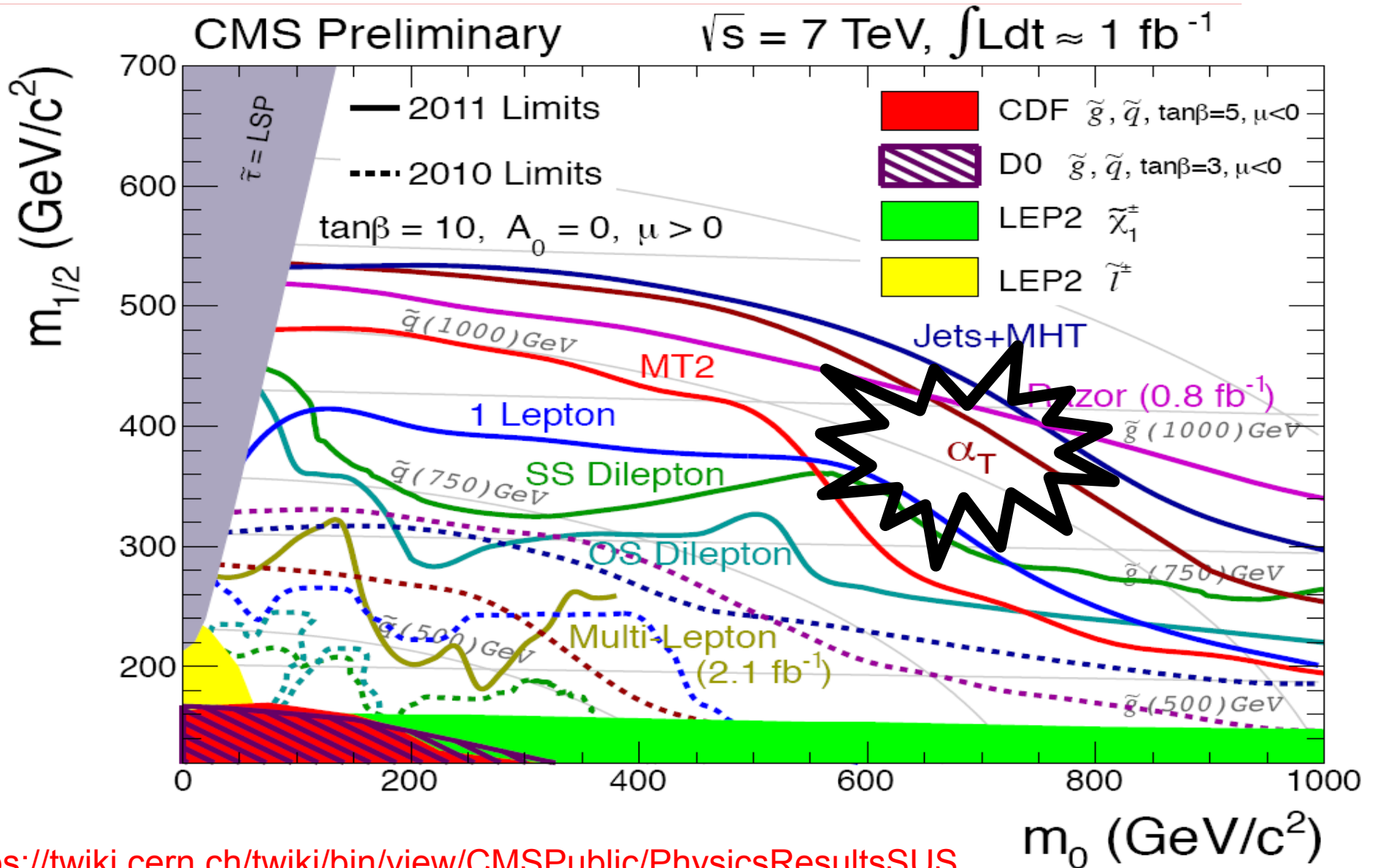
Likelihood from limits

Upper/lower limit for null result: exclusion.



- Use error function to smear the bound!
- Can add theory error as well.

SUSY search at CMS



How could we go beyond the CMS published α_T m_0 and $m_{1/2}$ range?

How do we include information from CMS α_T (1.1 /fb) analysis into our likelihood?

$$\text{Posterior} = \frac{\text{Prior} \times \text{Likelihood}}{\text{Evidence}}$$

Poisson Likelihood

Poisson distribution to characterize counting experiments.

$$\mathcal{L} = \prod_i \frac{e^{-(s_i + b_i)} (s_i + b_i)^{o_i}}{o_i!}$$

o_i : observed events in LHC.

b_i : expected SM background events.

s_i : $s_i = \epsilon_i \times \sigma \times \int L$.

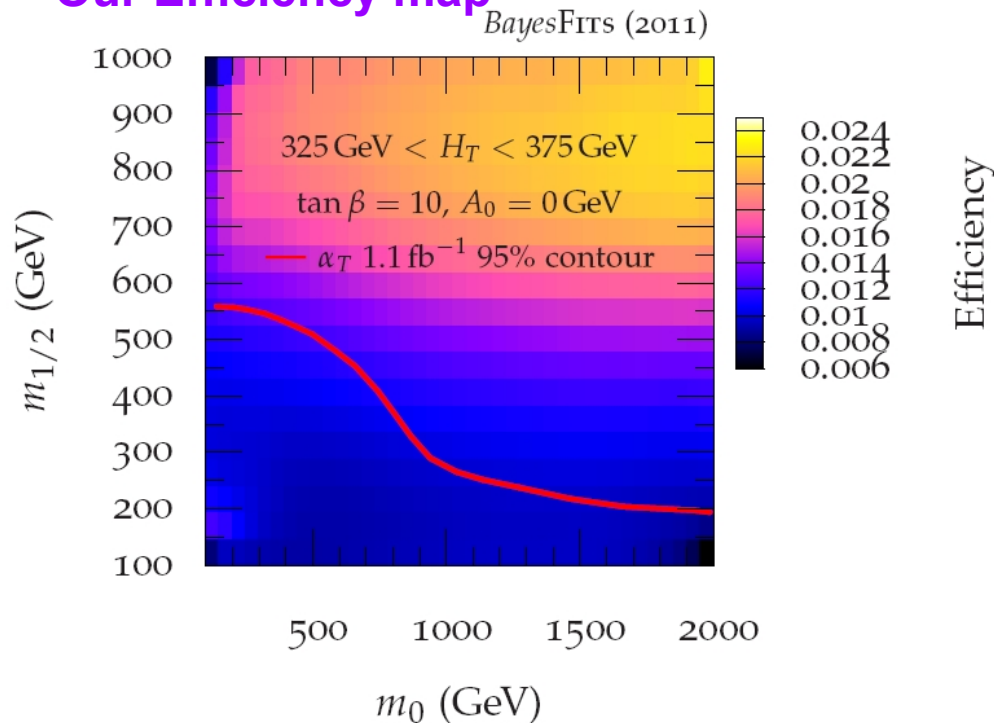
ϵ_i : $N_i(\alpha_T > 0.55) / N_{\text{total}}$

$i = 1, 2, 3, \dots, 8$.

Our approximate Efficiency maps

$$\epsilon_i : N_i(\alpha_T > 0.55) / N_{\text{total}}$$

Our Efficiency map

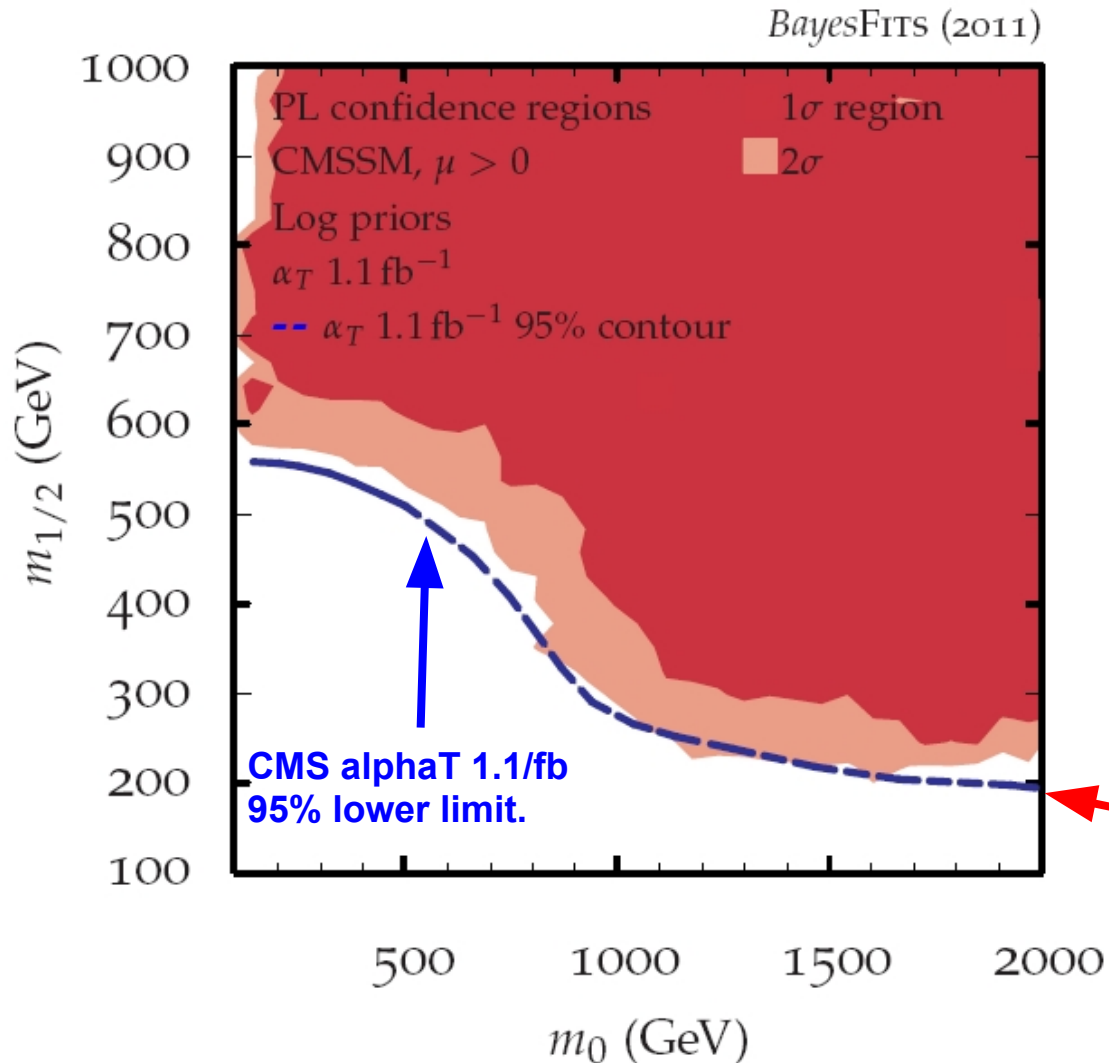


More selected events at large $m_{1/2}$!

(b) The case of 325 GeV < H_T < 375 GeV.

Limit from CMS alphaT 1.1/fb

Fowlie, Kalinowski, Kazana, Roszkowski, Tsai (arXiv:1111.6098)



$$\mathcal{L} = \prod_i \frac{e^{-(s_i+b_i)} (s_i + b_i)^{o_i}}{o_i!}$$

$$s_i = \epsilon_i \times \sigma \times \int L.$$

**VERY GOOD AGREEMENT WITH
 CMS 95% LIMIT!**

We can extend the range of m_0
 and $m_{1/2}$ to whatever we want.

SUSY: Summary of constraints

Observable	Mean	Exp. Error	Theor. Error	Likelihood	Distribution
Non-LHC:					
$\Omega_\chi h^2$	0.1120	0.0056	10%		Gaussian
$\sin^2 \theta_{\text{eff}}$	0.231160	0.00013	15.0×10^{-5}		Gaussian
M_W	80.399	0.023	0.015		Gaussian
$\delta(g-2)_\mu^{SUSY} \times 10^{10}$	30.5	8.6	1		Gaussian
$\mathcal{BR}(\bar{B} \rightarrow X_s \gamma) \times 10^4$	3.6	0.23	0.21		Gaussian
$\mathcal{BR}(B_u \rightarrow \tau \nu) \times 10^4$	1.66	0.66	0.38		Gaussian
ΔM_{B_s}	17.77	0.12	2.4		Gaussian
$\mathcal{BR}(B_s \rightarrow \mu^+ \mu^-)$	$< 1.5 \times 10^{-8}$	0	14%	Upper limit	Error fn
LEP — 95% Limits					
m_h	> 114.4	0	3	Lower limit	Error fn
ζ_h^2	$< f(m_h)$	0	0	Upper limit	Step fn
m_χ	> 50	0	5%	Lower limit	Error fn
$m_{\chi_1^\pm}$	$> 103.5 (92.4)$	0	5%	Lower limit	Error fn
$m_{\tilde{e}_R}$	$> 100 (73)$	0	5%	Lower limit	Error fn
$m_{\tilde{\mu}_R}$	$> 95 (73)$	0	5%	Lower limit	Error fn
$m_{\tilde{\tau}_1}$	$> 87 (73)$	0	5%	Lower limit	Error fn
$m_{\tilde{\nu}}$	$> 94 (43)$	0	5%	Lower limit	Error fn
LHC CMS α_T 1.1/fb analysis					
α_T	See text	See text	0		Poisson
XENON100					
$\sigma_p^{\text{SI}}(m_\chi)$	$< f(m_\chi)$ — see text	0	1000%	Upper limit	Error fn
Nuisance					
$1/\alpha_{em}(M_Z)^{\overline{MS}}$	127.916	0.015	0		Gaussian
m_t^{pole}	172.9	1.1	0		Gaussian
$m_b(m_b)^{\overline{MS}}$	4.19	0.12	0		Gaussian
$\alpha_s(M_Z)^{\overline{MS}}$	0.1184	0.0006	0		Gaussian

Dark Matter relic density

anomalous magnetic moment of the muon

Favour physics

We include all important constraints.

Dark matter search

Including all the constraints
into likelihood,
we can conduct a random
scan with prior range:

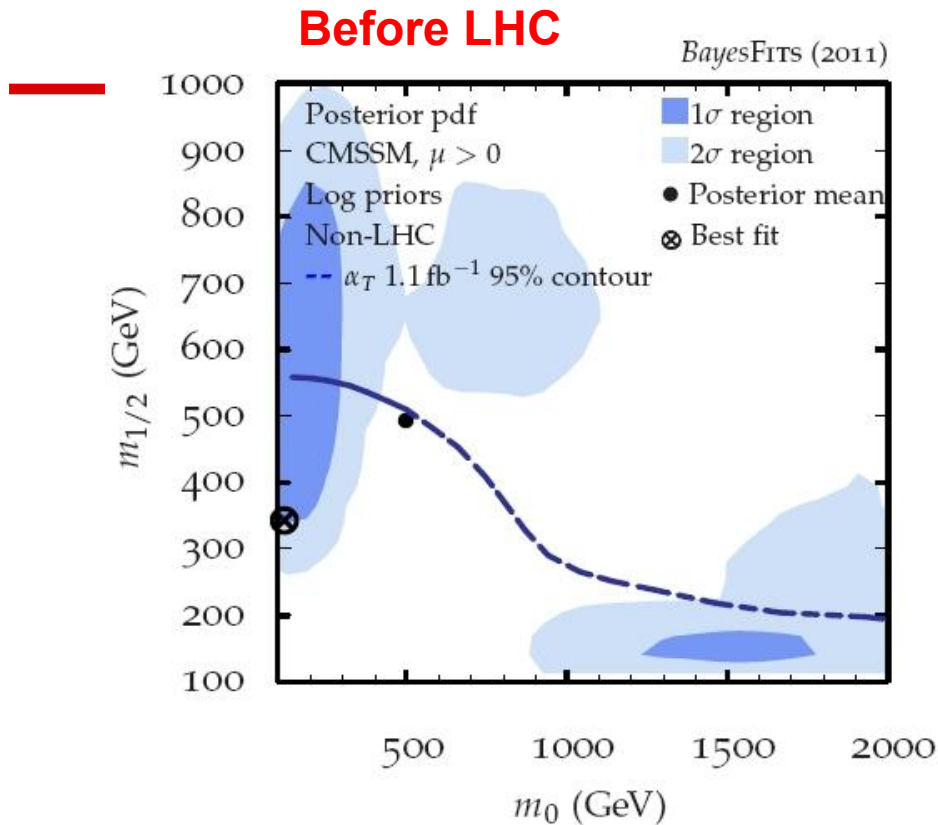
$$100 \text{ GeV} \leq m_0 \leq 4000 \text{ GeV}$$

$$100 \text{ GeV} \leq m_{1/2} \leq 2000 \text{ GeV}$$

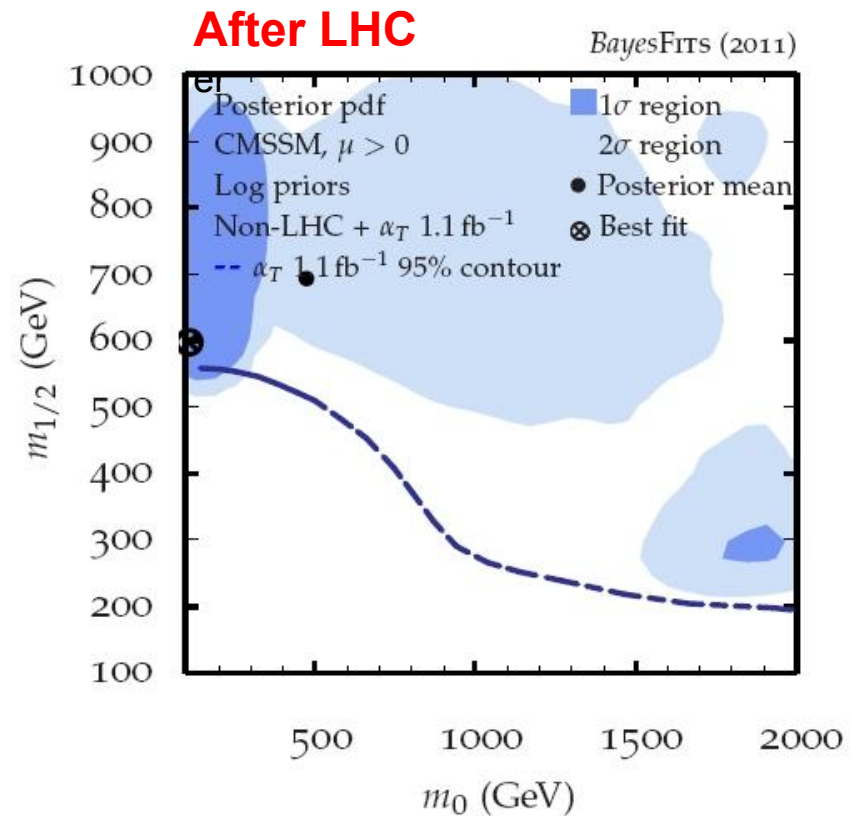
$$-2000 \text{ GeV} \leq A_0 \leq 2000 \text{ GeV}$$

$$3 \leq \tan \beta \leq 62 .$$

Impact of the alphaT limit on CMSSM



(a) The CMSSM's parameters constrained by the non-LHC experiments only.



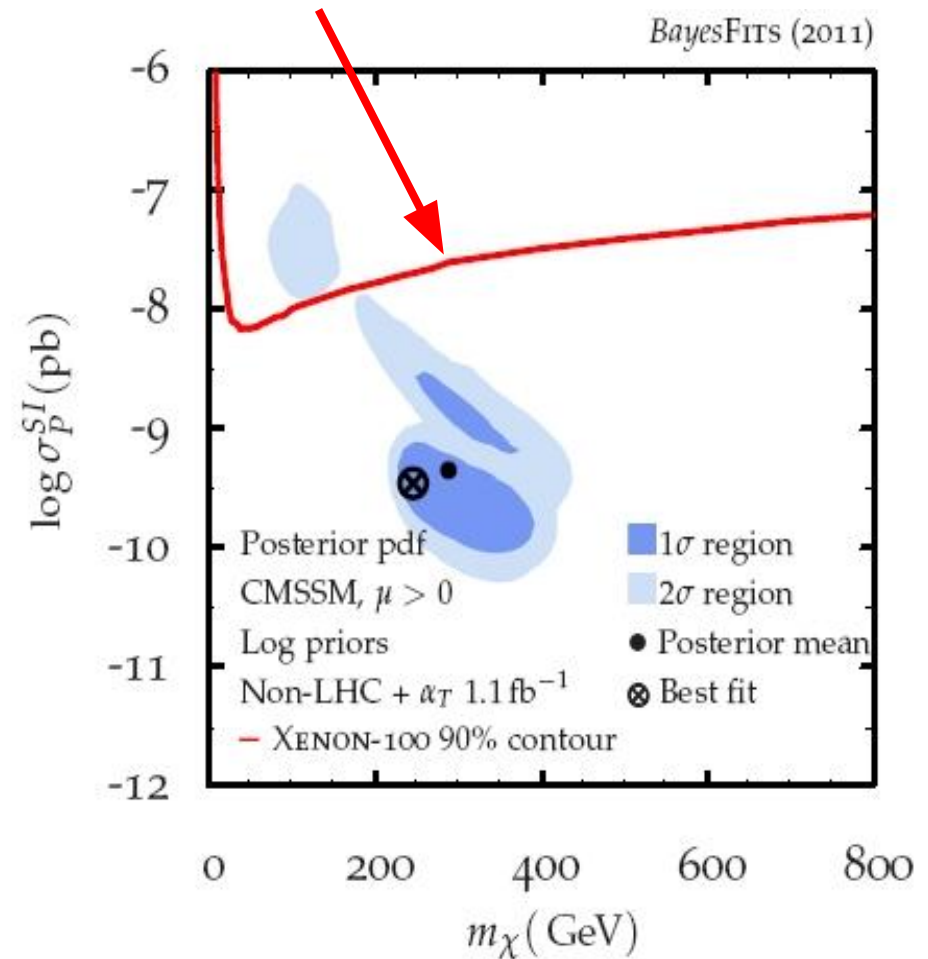
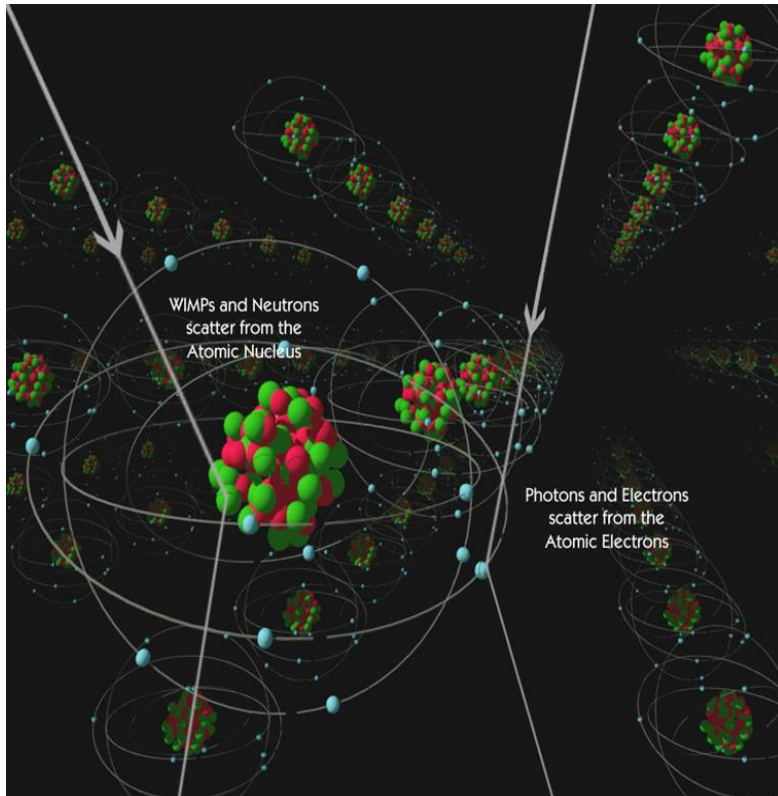
(b) The CMSSM's parameters constrained by the CMS $\alpha_T 1.1/\text{fb}$ analysis and by the non-LHC experiments.

General trend:

- favoured ranges are now pushed up.
- poorer fit, but
- best fit point remains just above the CMS 95% CL limit

Direct searches of dark matter

currently best limit from XENON100

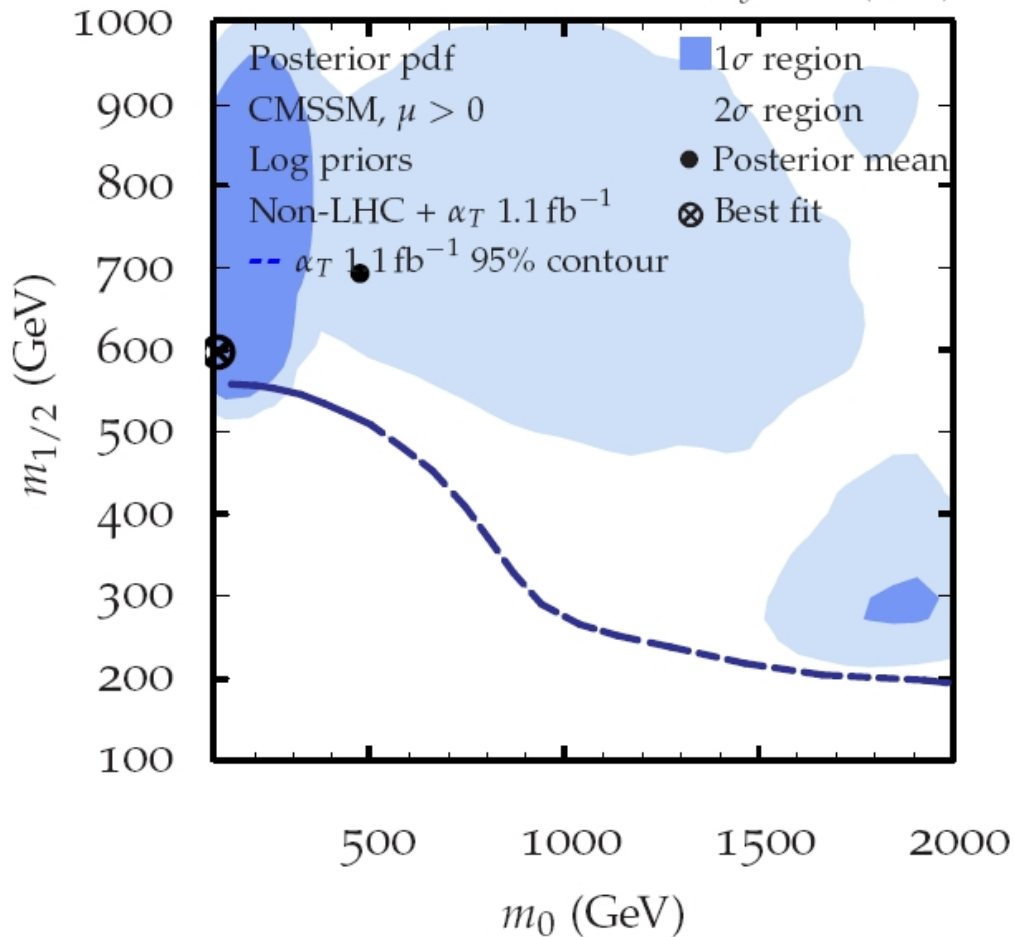


XENON100 limit not applied here

Impact on CMSSM

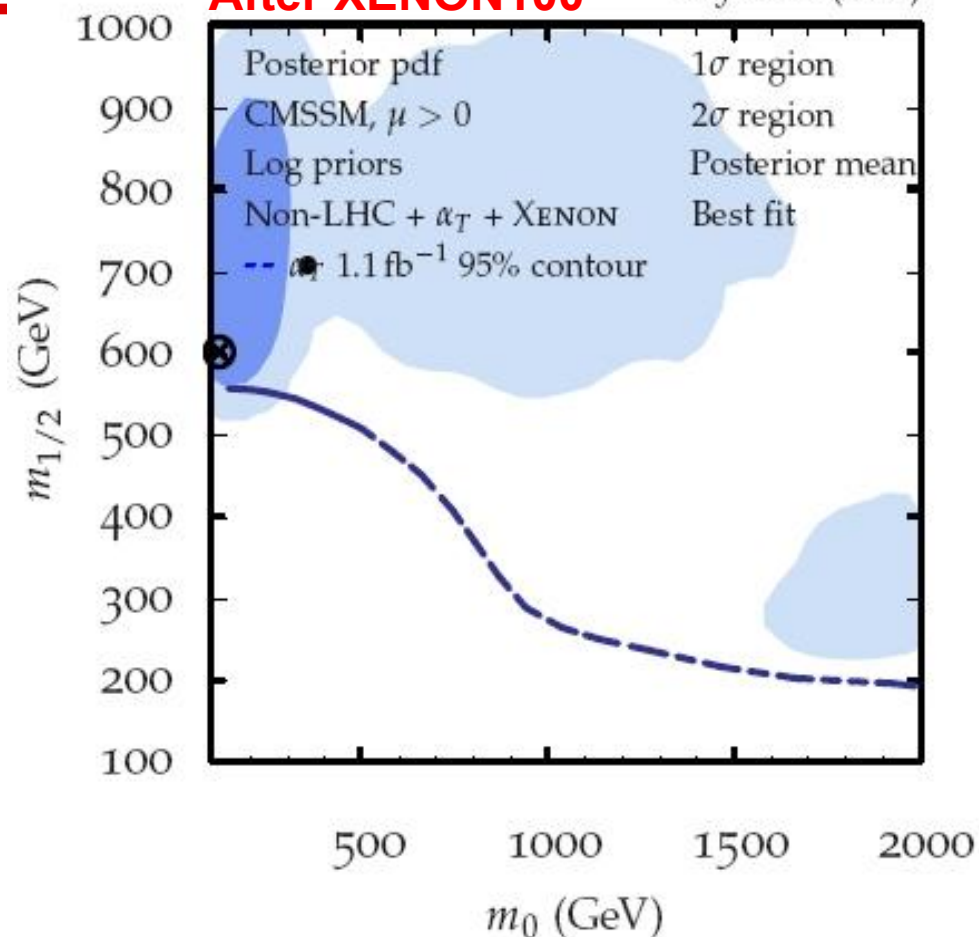
Before XENON100

BayesFITS (2011)



After XENON100

BayesFITS (2011)



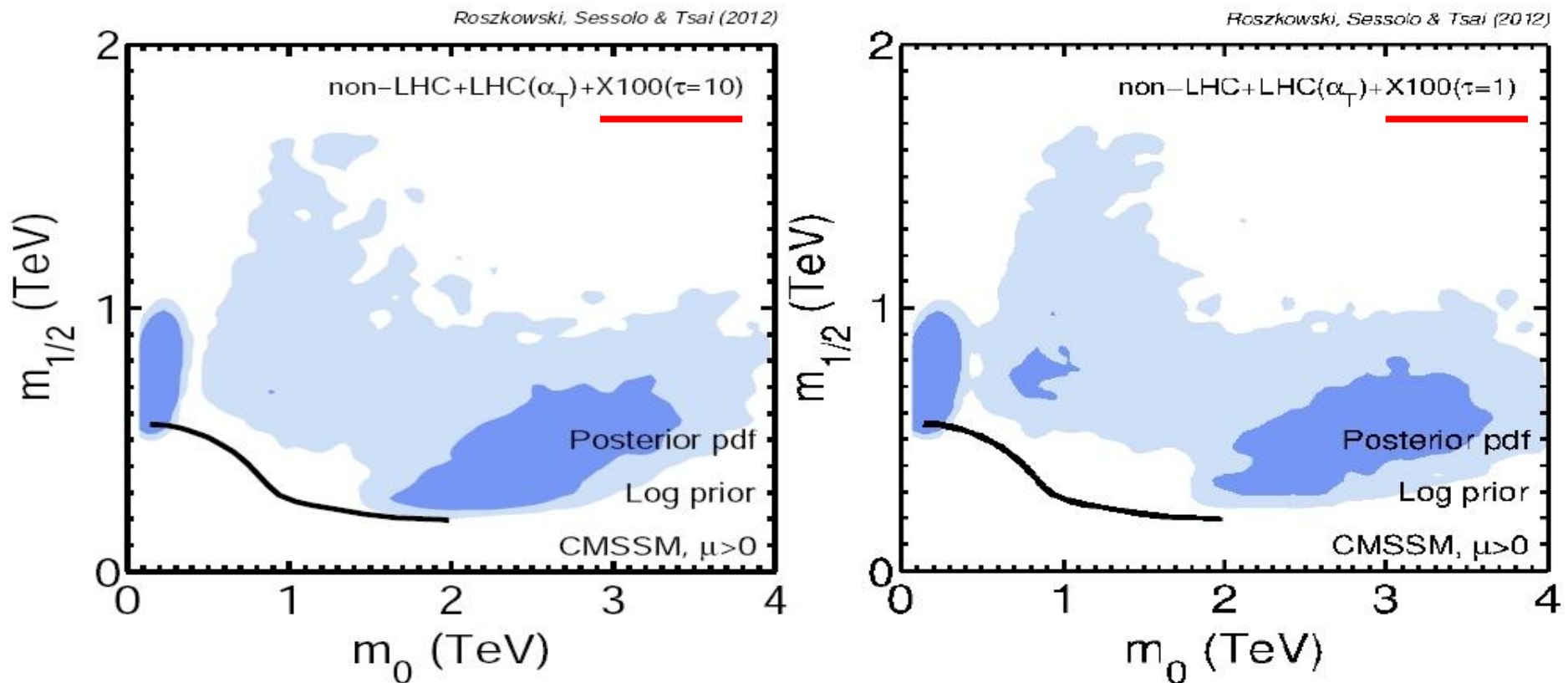
- DM XENON100 limit has small additional effect



LHC limits on CMSSM are stronger

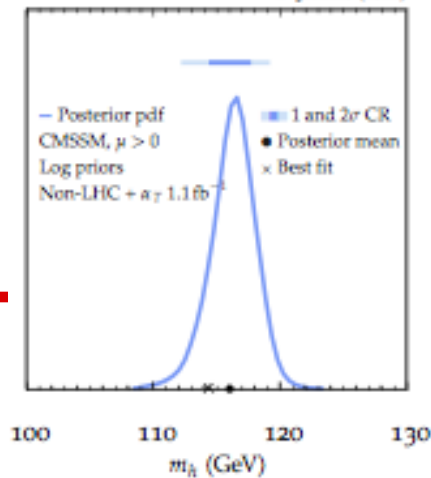
Also: large theoretical uncertainties (\sim factor of 10) in evaluating σ_p

Impact of Xenon100 detection



- The Xenon100 limit shows that DM searches are still marred by large (\sim factor of 10) theoretical uncertainties.
- Even if reduce them to ~ 1 will make little difference.
- Current LHC limits are much stronger.

LHC: Posterior 68% and 95% C.L. ranges of sparticles and Higgs masses



Some examples:

Mass (GeV)	68%	95%	68%	95%
	Non-LHC		Non-LHC + CMS α_T 1.1/fb limit + XENON100	
m_h	(112.3, 116.5)	(110.1, <u>118.4</u>)	(114.4, 117.8)	(112.2, <u>119.4</u>)
m_χ	(56, 291)	(<u>53</u> , 356)	(250, 343)	(<u>128</u> , 390)
$m_{\chi_1^\pm}$	(110, 554)	(104, 676)	(475, 651)	(181, 738)
$m_{\tilde{q}}$	(326, 808)	(<u>254</u> , 1172)	(434, 761)	(<u>398</u> , 1302)
$m_{\tilde{g}}$	(403, 1576)	(<u>384</u> , 1885)	(1380, 1825)	(<u>879</u> , 2043)

Summary

- **With 1.1/fb at LHC: improved limits on SUSY particle masses.**
- **The CMSSM has become severely constrained but not excluded.**
- Constraints from direct detection of dark matter are currently weaker. One reason: still large theoretical uncertainties.
- **Our method is completely general. It can be applied to other models (SUSY or not).**
- **We have developed a framework based on Bayesian approach to include limits and future signals from the LHC.**
- **Work in progress: updated limits including Razor 5/fb and impact of possible Higgs signal.**

The End.
Thank you for your
attention.
