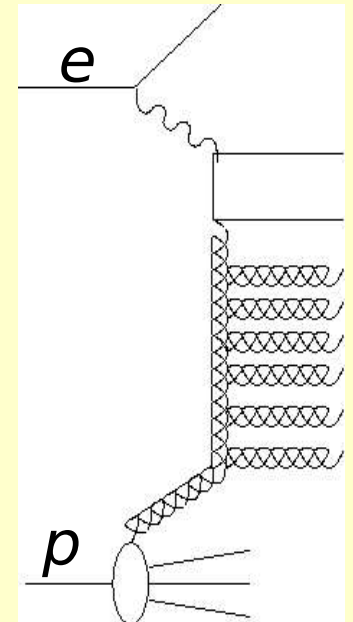


# HERA Today and Tomorrow

- HERA-LHC Workshop
- Status report on HERA and the HERA physics
- Hadronic final states Investigated in Warsaw
  - + Azimuthal Asymmetry in  $ep \rightarrow ehX$
  - Instantons
  - Glueballs
  - Pentaquarks
  - Broken chiral symmetry



# HERA – LHC workshop



**HERA AND THE LHC**  
A workshop on the implications of HERA for LHC physics

March 2004 - January 2005

Parton density functions  
Multijet final states and energy flow  
Heavy quarks  
Diffraction  
Monte Carlo tools

**Startup Meeting**  
March 26-27 2004  
**Midterm Meeting**  
11-13 October 2004  
CERN, Geneva

**Final Meeting**  
January 2005  
DESY, Hamburg

**Organizing Committee:**  
G. Altarelli (CERN), J. Blumlein (DESY),  
M. Boley (DESY), J. Butterworth (DESY),  
A. De Roeck (CERN) (chair), K. Eggert (CERN),  
H. Jung (DESY) (chair), M. Mangano (CERN),  
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G. Terese (ANL)

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J. Sjöstrand (CERN), W. K. Tung (Michigan State),  
A. Wagner (DESY), H. Yasuda (ANL)

[www.desy.de/~heralhc](http://www.desy.de/~heralhc) [heralhc.workshop@cern.ch](mailto:heralhc.workshop@cern.ch)

## Meetings:

CERN October 11-13 2004  
DESY January 17-21 2005

- Parton density functions
- Multijet final states and energy flow
- Heavy quarks production
- Diffraction
- Monte Carlo tools

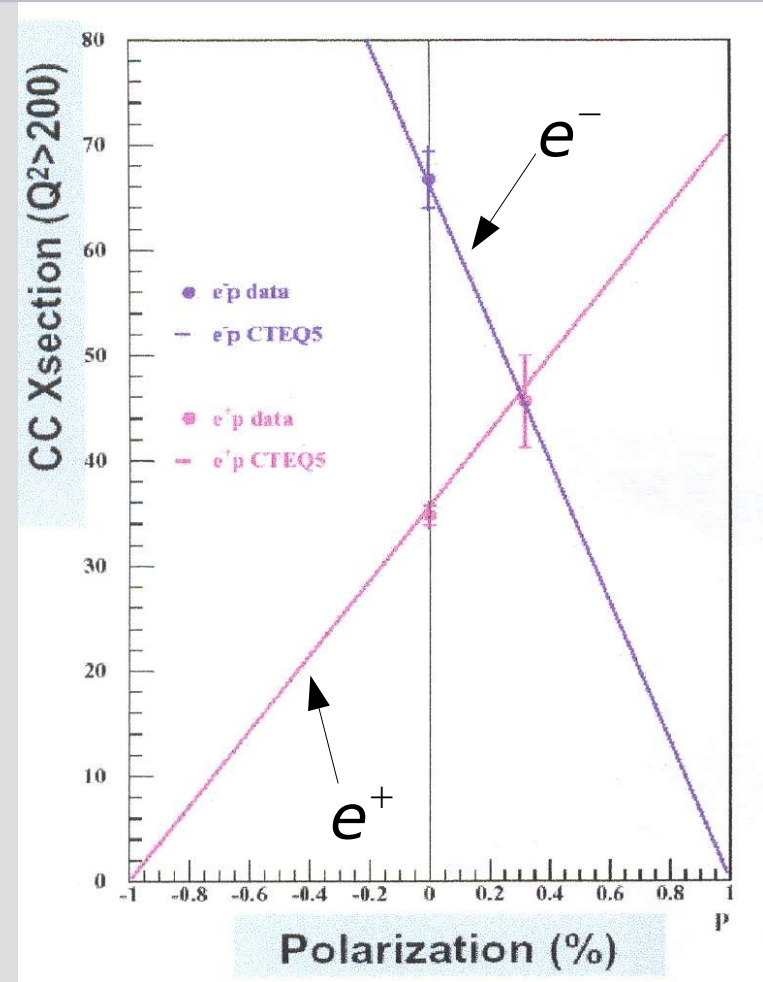
# Status report 2003-2004

- Polarized positron beam  
integrated luminosity  
in 2004  $20 \text{ pb}^{-1}$
- Problems in 2003  
synchrotron radiation  
detectors 15 years old

- Test of standard model

$$\sigma(e_R^- p \rightarrow \nu p) = 0$$

$$\sigma(e_L^+ p \rightarrow \bar{\nu} p) = 0$$

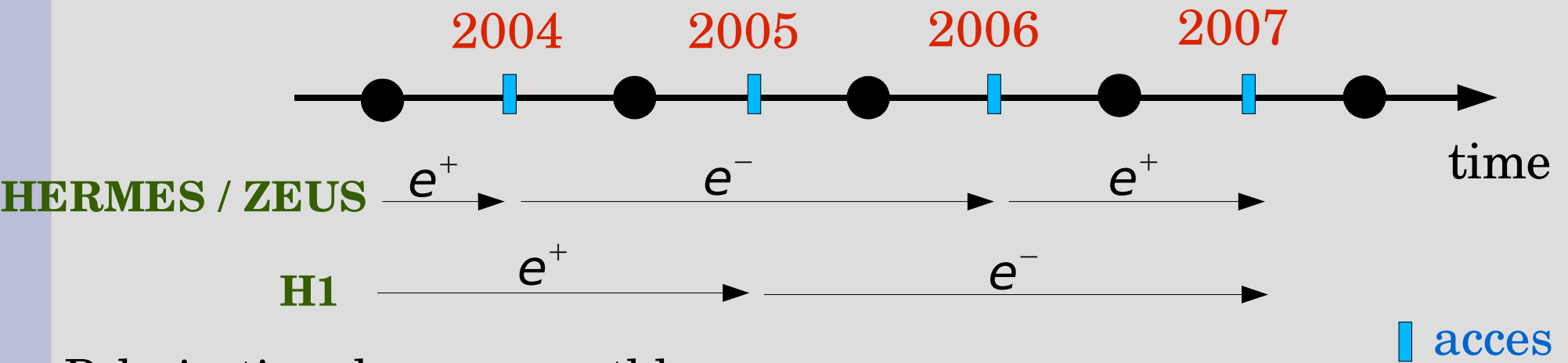


$$\sigma = 38.1 \pm 2.9(\text{stat}) \pm 0.8(\text{syst}) \text{ pb}$$
$$\int L = 6.6 \text{ pb}^{-1}$$



# 2004-2007

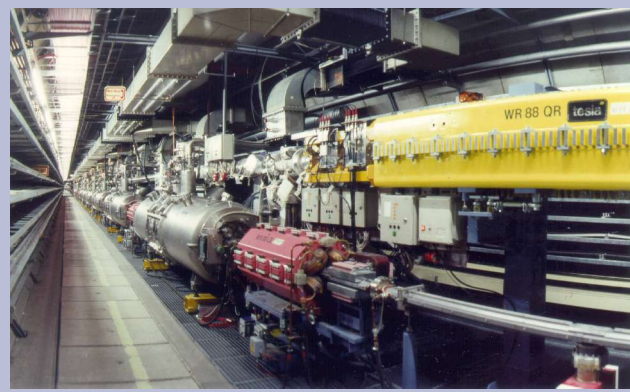
- Wiązka spolaryzowana leptonowa ( $e^+e^-$ ) i protonowa  
 polaryzacja docelowe 50 %  
 aktualne 30-40 %
- Planed time table – decision in April 2004



Polarization changes monthly

## AIMS (H1, ZEUS)

- $F_L$
- beyond SM



## ARGUMENTS

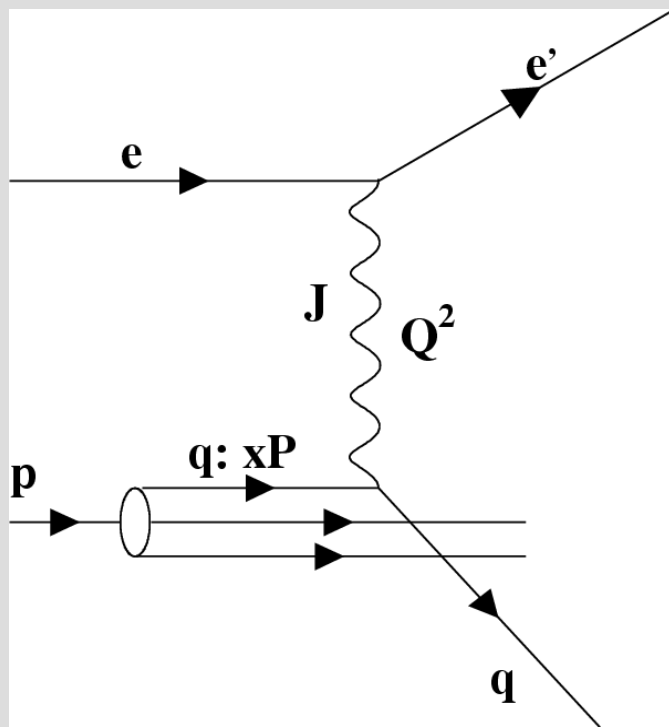
- HERMES  
recoil detector installed summer 2004  
start summer 2005
- H1  
larger luminosity for  $e^+$   
more multilepton events
- ZEUS  
not enough data for  $e^-$   
large sample for  $e^+$

# Azimuthal asymmetry using energy flow method

- Azimuthal angle distribution at  $Q^2 > 100 \text{ GeV}^2$
- Energy flow method and multiplicity method
- Experimental results
- LEPTO predictions for azimuthal asymmetry
  - BGF+QCDC+QPM
    - ★ without PS
    - ★ with PS
    - ★ with hadronization
- NLO DISENT predictions
- Summary



# HERA kinematics

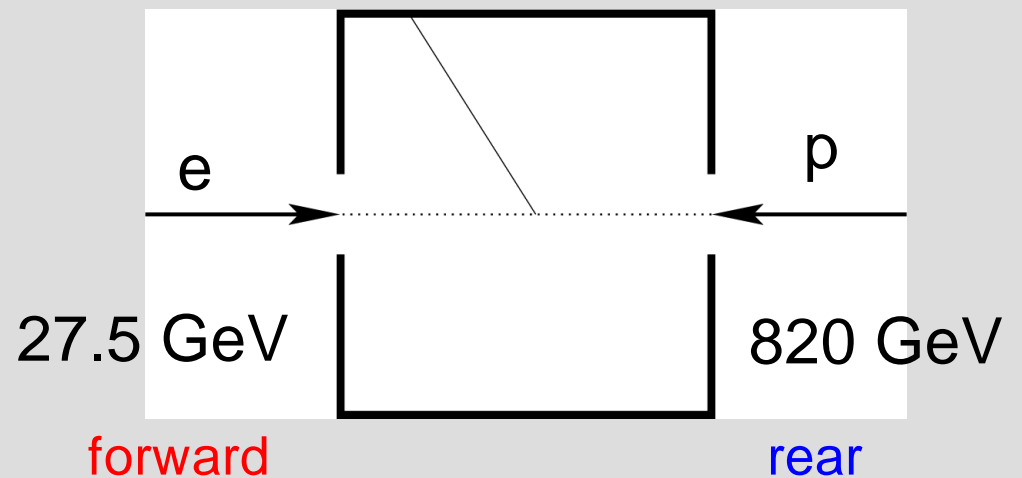


$$J = \left\{ \begin{array}{l} \gamma, Z^0 - NC \\ W - CC \end{array} \right\}$$

$$Q^2 = -q^2 = -(k_{e'} - k_e)^2$$

$$x = Q^2 / 2 P \cdot q$$

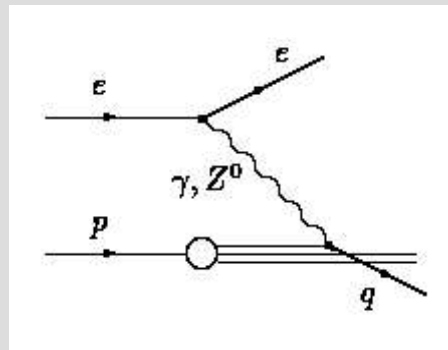
$$y = P \cdot q / P \cdot k_e$$



# Azimuthal angle definition for the $ep \rightarrow ehX$ process

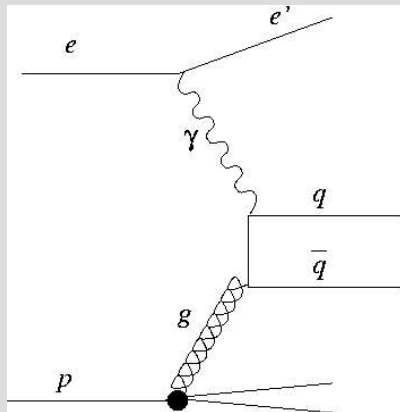
$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = A + B \cos(\phi) + C \cos(2\phi) + D \sin(\phi) + E \sin(2\phi)$$

Zeroth order QCD process  
(Simple DIS, QPM)

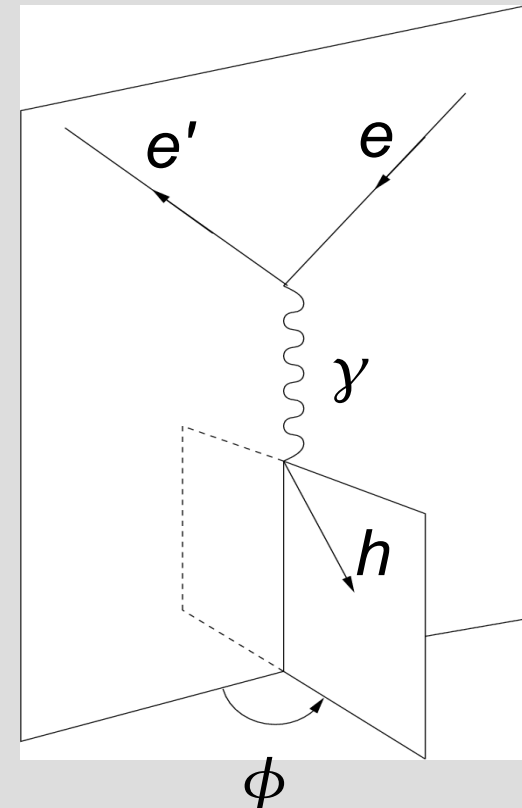
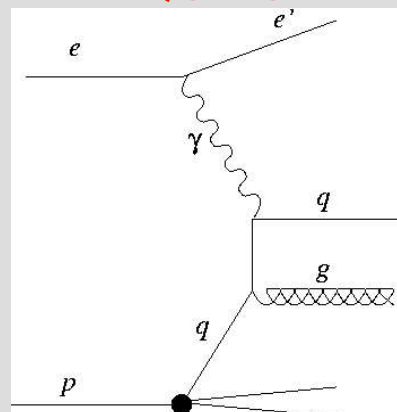


First order QCD processes

BGF



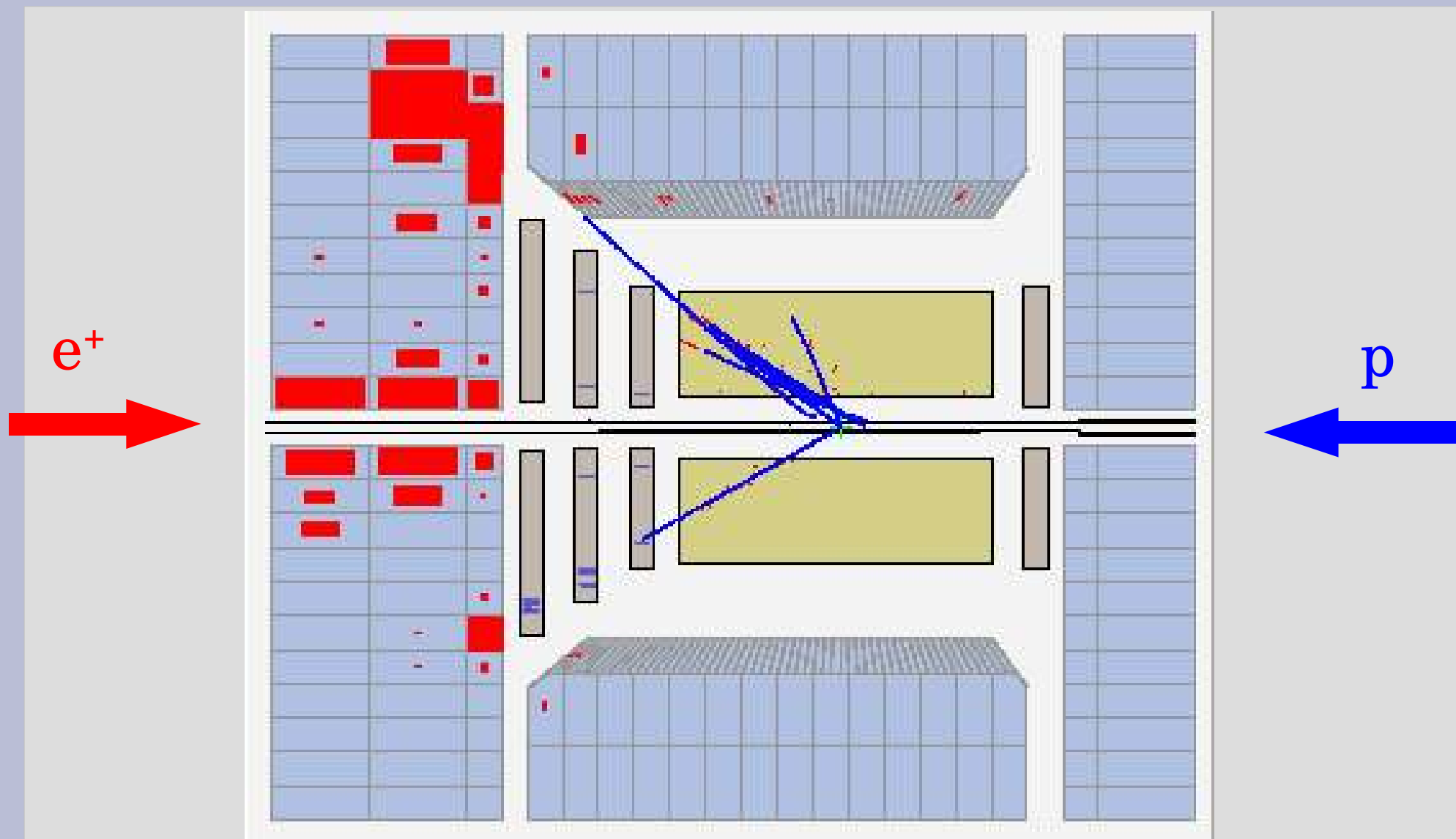
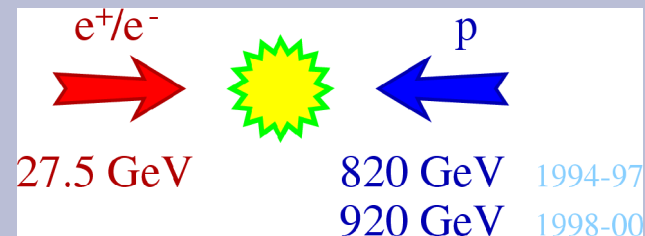
QCDC



$\gamma^* p$  CMS frame



# DIS NC



# Azimuthal angle distribution for the $ep \rightarrow ehX$ process

$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = A + B \cos(\phi) + C \cos(2\phi) + D \sin(\phi) + E \sin(2\phi)$$

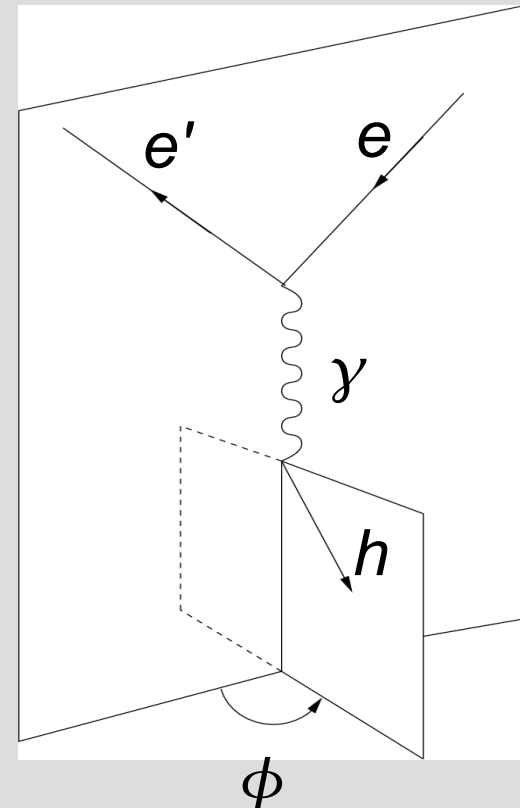
Azimuthal asymmetry comes from:

- ★ Two-body processes (BGF and QCDC)
- ★ Boson polarization
- ★ Longitudinally polarized electron beam
- ★ Parity violating weak interactions
- ★ Final hadron polarization
- ★ Intrinsic parton momentum in the proton

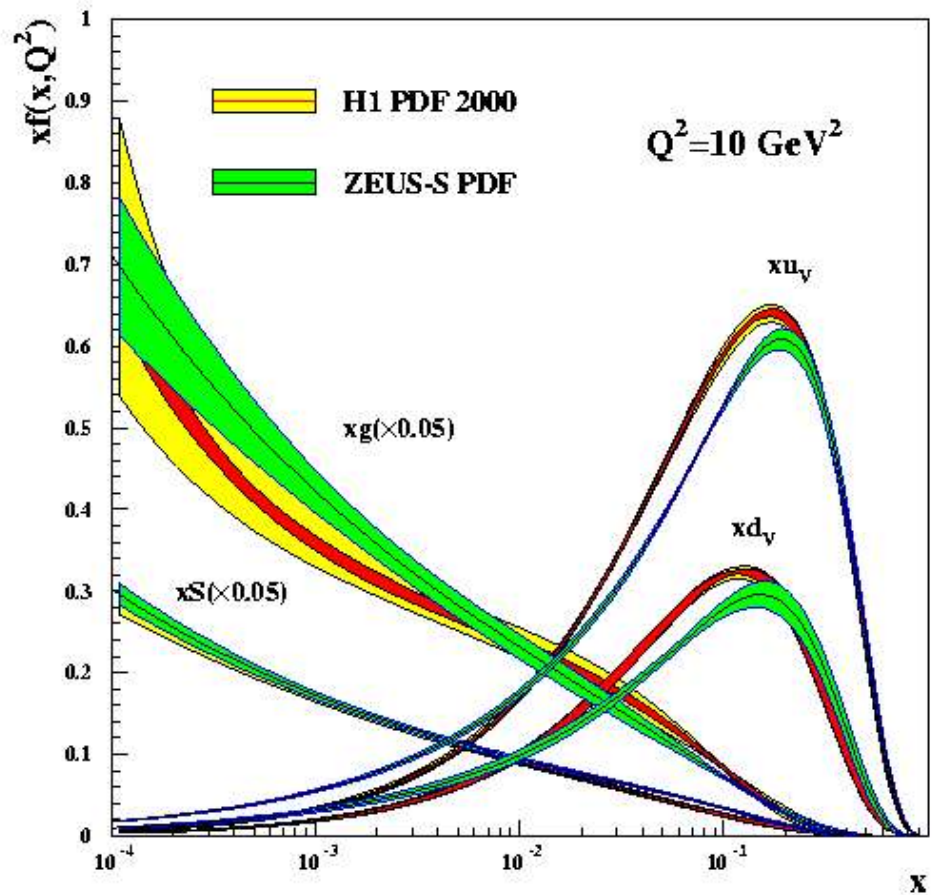
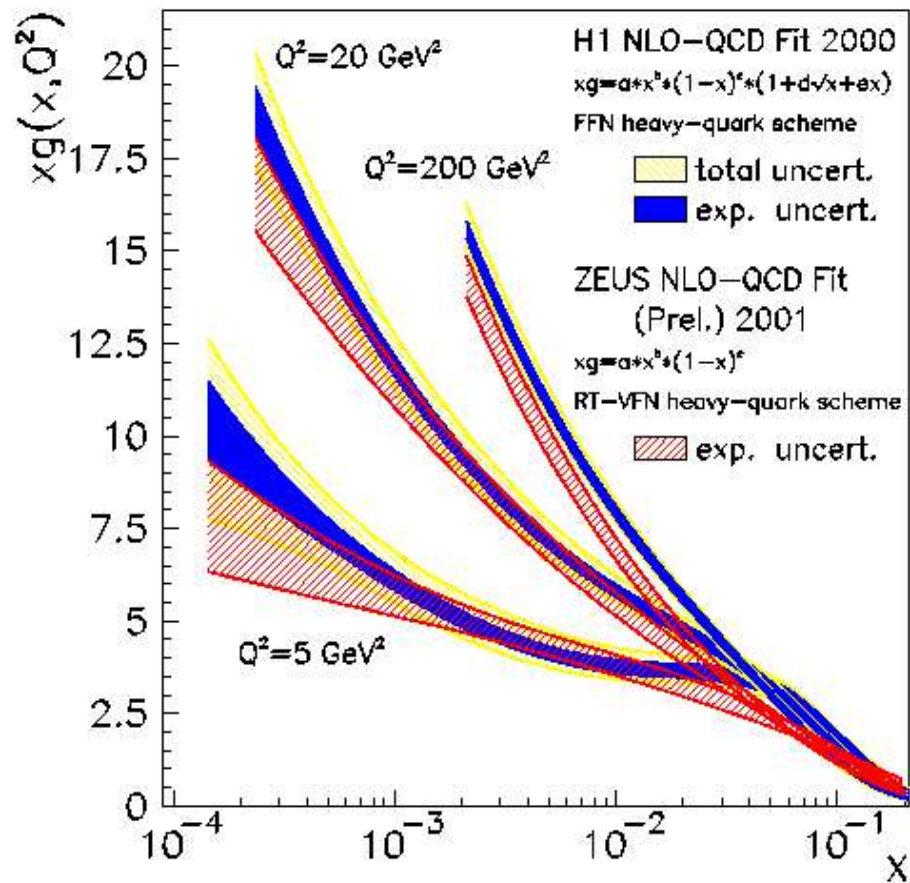
Future:

asymmetry can be measured for

- ◆ Longitudinally polarized lepton beam
- ◆ CC events



# Parton densities



good agreement between quark densities from H1 and ZEUS  
 some differences between gluon densities from H1 and ZEUS

$\phi$  distribution sensitive to gluon content at  $0.01 < x < 0.1$

# Boson polarization

- See J.G.Könner, E.Mirkes, G.Schuler, INT.J.Mod.Phys. A4 (1989) 1781
- Bronkorb, Mirkes, Z.Phys. C66 (1995) 141

$$\frac{d\sigma}{d\phi} = d\sigma_T + d\sigma_L + d\tau_{LT} \cdot \cos(\phi) + d\tau_{TT} \cdot \cos(2\phi)$$

Modify  $p_T$  distribution



## T-odd effects

$$\dots + D \sin(\phi) + E \sin(2\phi)$$

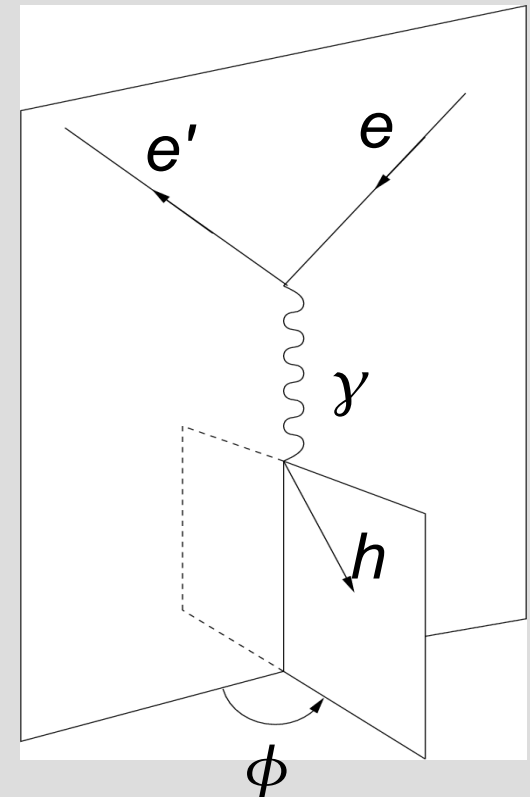
K.Hagiwara, K.Hikasa, N.Kai,  
Phys.Rev. D27 (1983) 84

# Azimuthal angle distribution for the $ep \rightarrow ehX$ process

$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = A + B \cos(\phi) + C \cos(2\phi) + D \sin(\phi) + E \sin(2\phi)$$

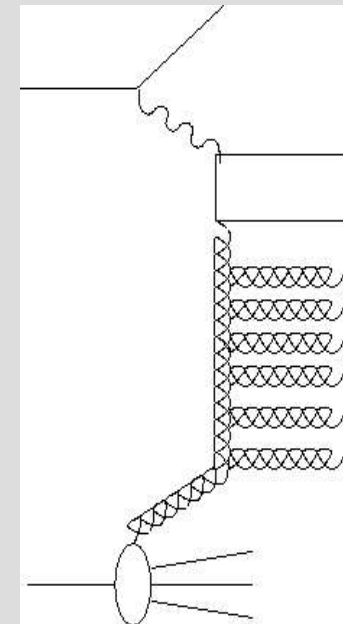
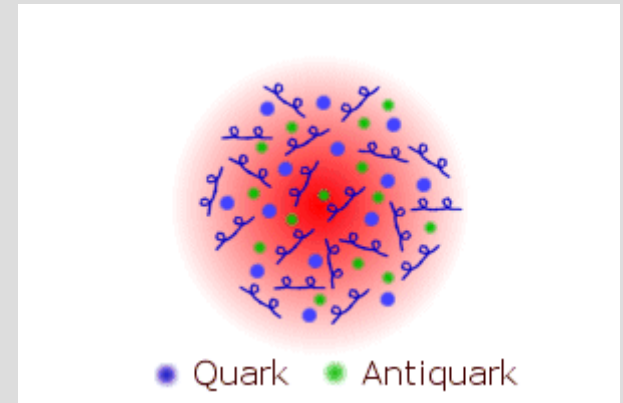
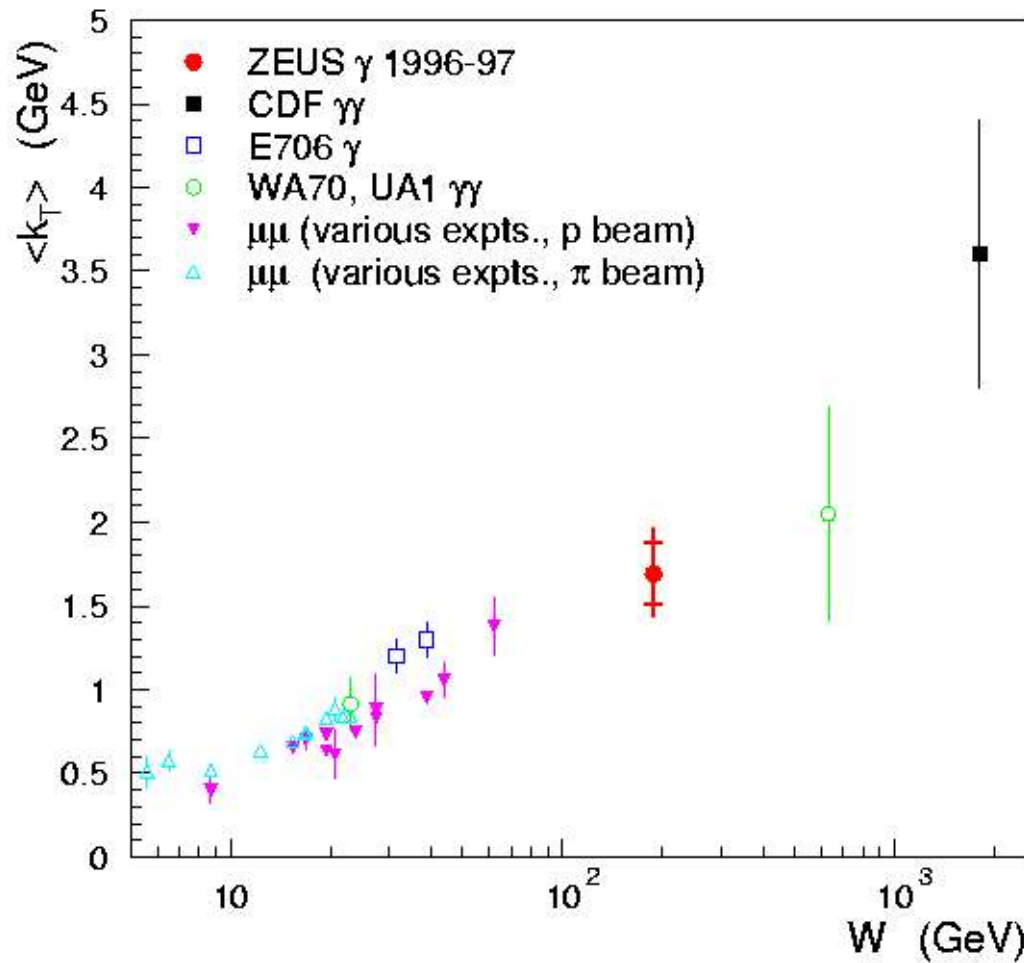
Azimuthal asymmetry comes from:

- ★ Two-body processes (BGF and QCDC)
- ★ Boson polarization
- ★ **Longitudinally polarized electron beam**
- ★ **Parity violating weak interactions**
- ★ **Final hadron polarization**
- ★ **Intrinsic parton momentum in the proton**



# $k_T$ distribution

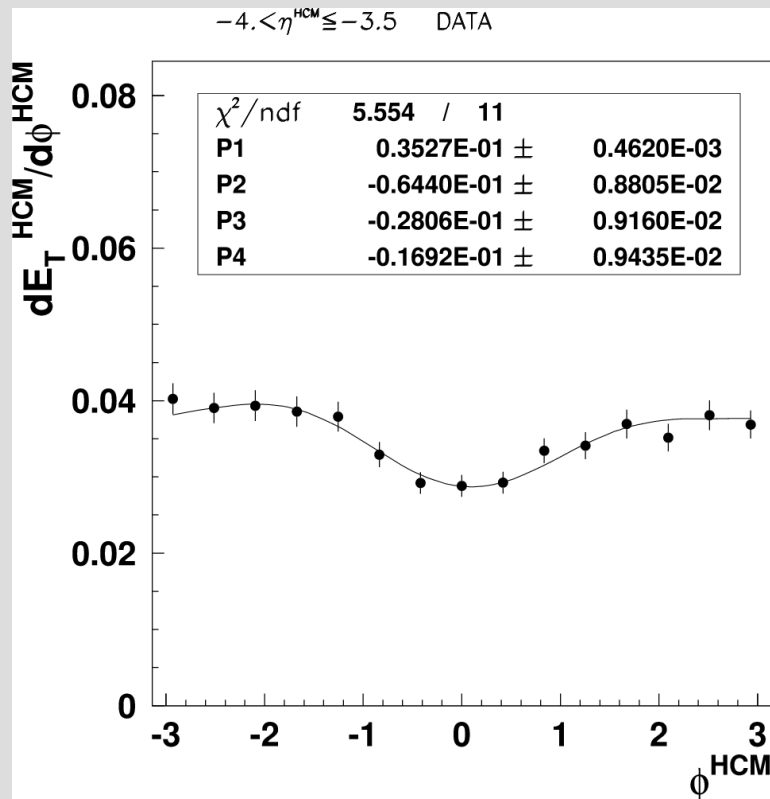
Intrinsic parton momentum



ZEUS:  $ep \rightarrow e \gamma_{jet} X$

# Experimental methods

$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = P_1 + P_2 \cos(\phi) + P_3 \cos(2\phi) + P_4 \sin(\phi)$$



- ◆ Fitted function
- ◆ Moments of distributions of trigonometrical functions
  - ☆ Mean
  - ☆ RMS

$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = A + B \cos(\phi) + C \cos(2\phi) + D \sin(\phi) + E \sin(2\phi)$$

## Moments:



$$\langle \cos(n\phi) \rangle = \frac{\int d\sigma \cos(n\phi)}{\int d\sigma}$$

## Means:

$$\langle \cos(\phi) \rangle = \frac{B}{2A}$$

$$\langle \cos(2\phi) \rangle = \frac{C}{2A}$$

$$\langle \sin(\phi) \rangle = \frac{D}{2A}$$

$$\langle \sin(2\phi) \rangle = \frac{E}{2A}$$

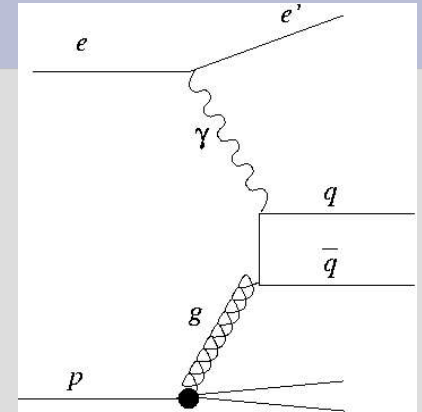
## RMS:

e.g.

$$RMS_{\langle \cos(2\phi) \rangle} = \sqrt{\frac{1}{2} + \frac{C}{2A}}$$



$$\frac{d\sigma^{ep \rightarrow ehX}}{dx dQ^2 d\phi} = A + B \cos(\phi) + C \cos(2\phi) + D \sin(\phi) + E \sin(2\phi)$$



$$\langle \cos(\phi) \rangle = \frac{B}{2A}$$

$$\frac{d\sigma}{dx dy dz dp_T^2} = \sum_i \text{str.func.} \cdot \frac{d\sigma_{\gamma i \rightarrow j}}{dx_h dy dz h dp_T^2} \cdot D_j(\xi', Q^2)$$

$D_j = PS + \text{hadronization}$

$$x_h = \frac{x}{\xi} = \frac{Q^2}{2p_i q}$$

$$z_h = \frac{z}{\xi'} = \frac{p_i p_j}{p_i q}$$

$$\frac{\alpha^2 Q_q^2}{16 \pi Q^2} y \delta(\dots) L_{\mu\nu} M_{ij}^{\mu\nu}$$

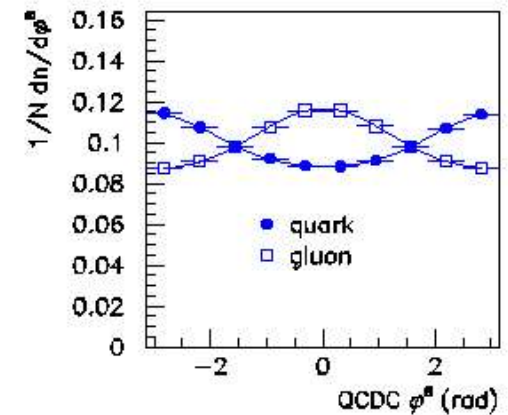


# For pure photon exchange the currents can be written:

**QCDC**

$$L_{\mu\nu}M_{qq}^{\mu\nu} = \frac{64\pi}{3}\alpha_s Q^2 \frac{1}{y^2} \left[ (1 + (1-y)^2) \left( \frac{x_h^2 + z_h^2}{(1-x_h)(1-z_h)} + 2(x_h z_h + 1) \right) \right. \\ \left. - 4(2-y)\sqrt{1-y} \sqrt{\frac{x_h z_h}{(1-x_h)(1-z_h)}} (x_h z_h + (1-x_h)(1-z_h)) \cos \phi \right. \\ \left. + 4x_h(1-y)z_h \cos 2\phi \right]$$

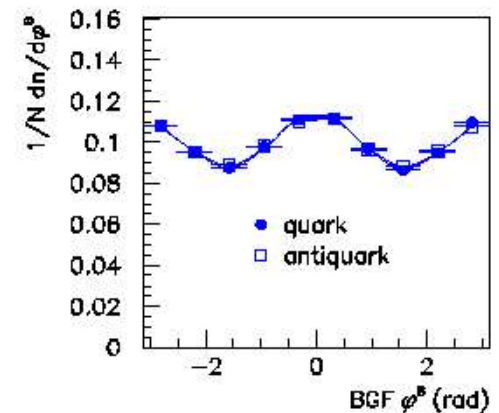
$e^\eta$



$\eta$  pseudorapidity

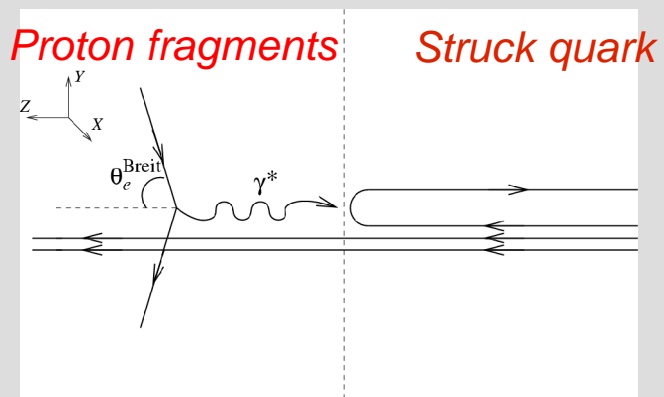
**BGF**

$$L_{\mu\nu}M_{qq}^{\mu\nu} = 8\pi\alpha_s Q^2 \frac{1}{y^2} \left[ (1 + (1-y)^2) (x_h^2 + (1-x_h)^2) \frac{z_h^2 + (1-z_h)^2}{z_h(1-z_h)} \right. \\ \left. - 4(2-y)\sqrt{1-y} \sqrt{\frac{x_h(1-x_h)}{z_h(1-z_h)}} (1 - 2x_h)(1 - 2z_h) \cos \phi \right. \\ \left. + 8x_h(1-x_h)(1-y) \cos 2\phi \right]$$



# Previous Measurement

# The HCM / Breit frame

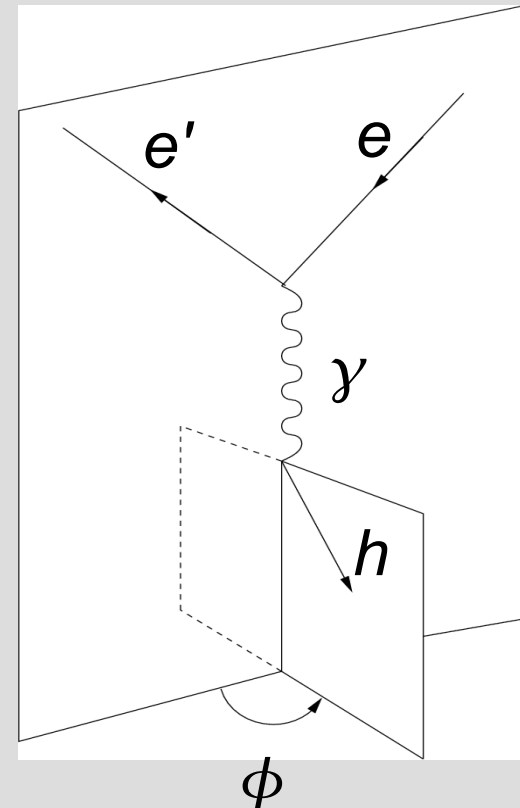


$$\mathbf{E}_T^{\text{BR}} = \mathbf{E}_T^{\text{HCM}}$$

$$\phi^{\text{BR}} = \phi^{\text{HCM}}$$

For massless partons or particles

$$\eta^{\text{BR}} \approx \eta^{\text{HCM}} + 0.5 \cdot \ln\left(\frac{1-x}{x}\right)$$



# Distribution of the azimuthal angle – ZEUS paper 2002

Breit frame

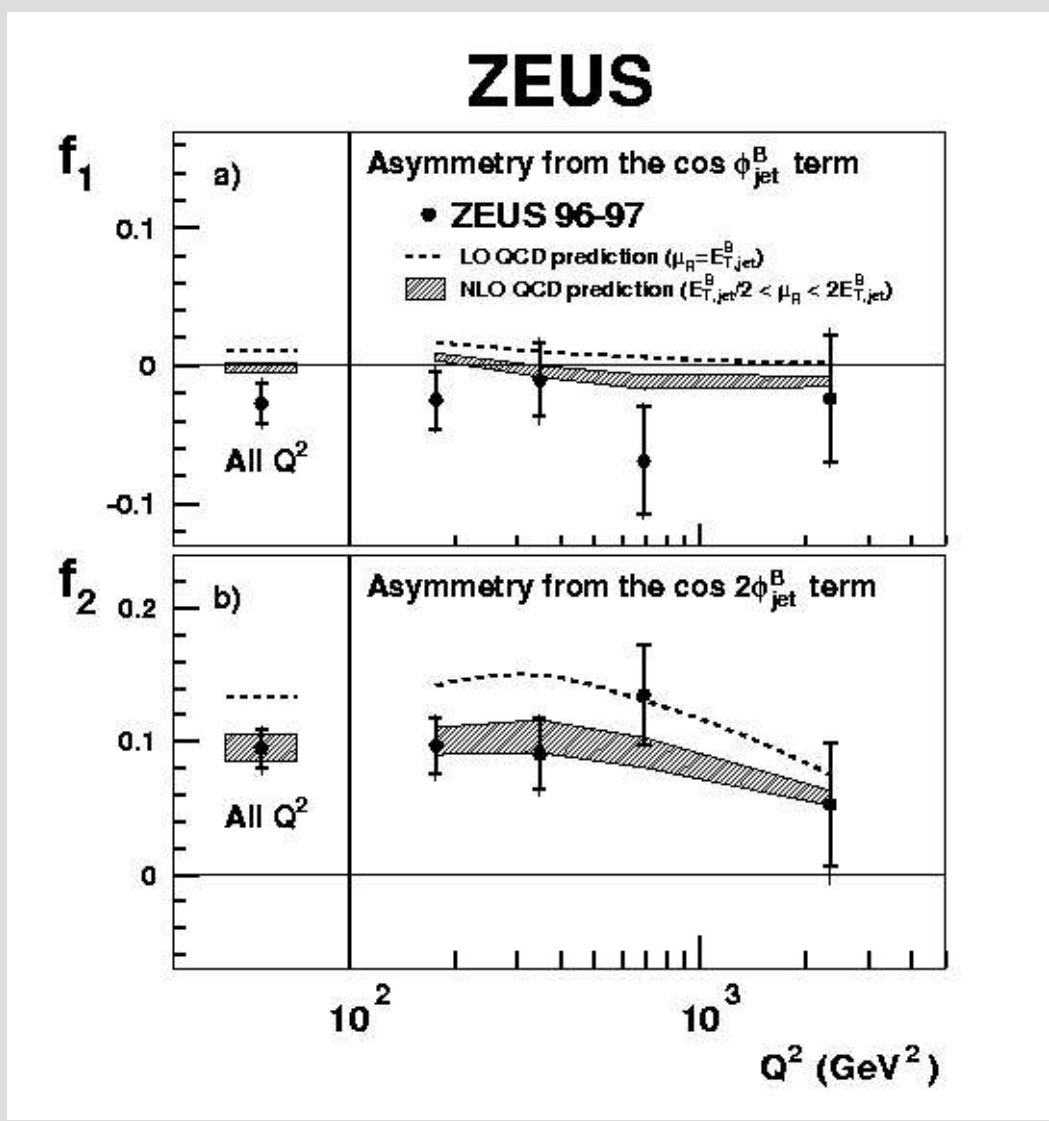
## Jet analysis

Fitted to experimental data

$$\frac{1}{\sigma} \frac{d\sigma}{d|\phi_{jet}^B|} = \frac{1}{\pi} [1 + f_1 \cos(\phi_{jet}^B) + f_2 \cos(2\phi_{jet}^B)]$$

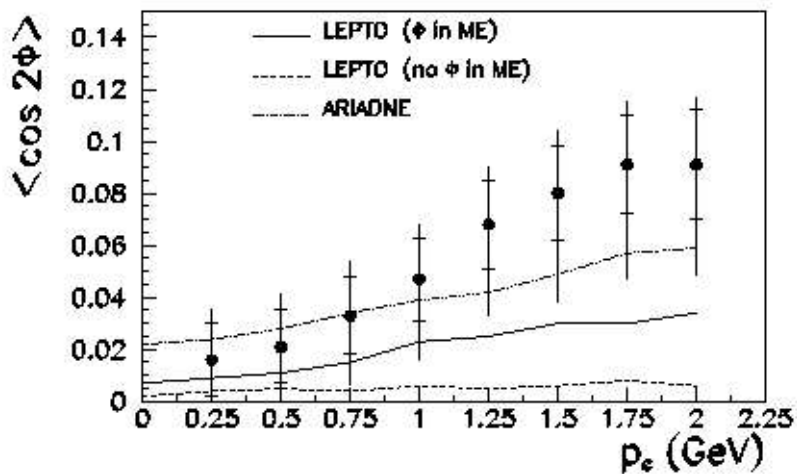
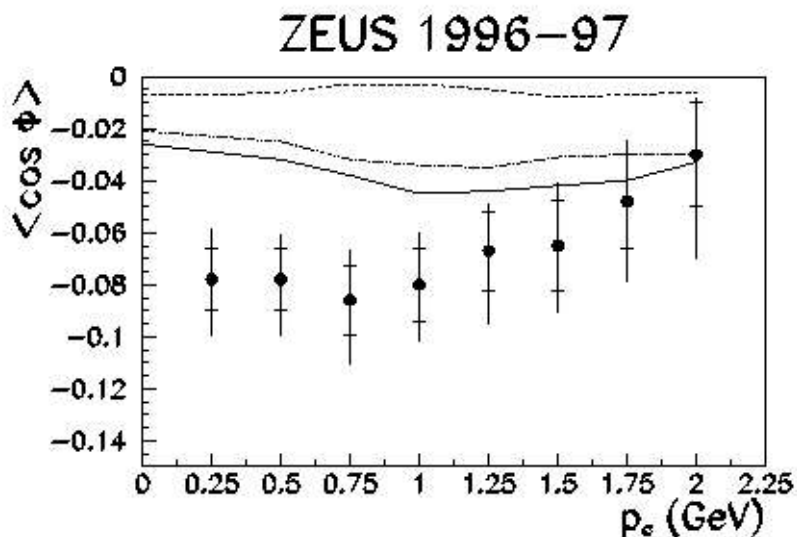
↑
↑  
 Small                      Large

See Oscar Gonzalez thesis and  
 DESY-02-171, Phys.Lett. B551 (2003) 226-240



# Distribution of the azimuthal angle – ZEUS paper 2000

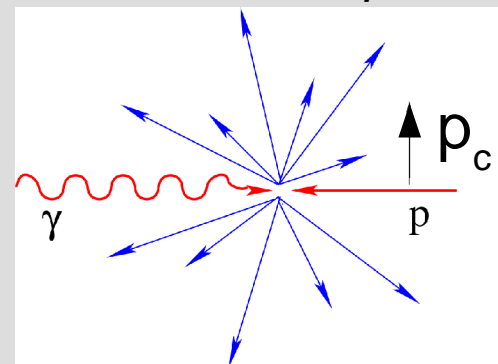
HCM frame  
z, p<sub>T</sub> method



**Multiplicity method**  
**Charged hadrons**

For hadrons with  $p_T > p_c$

$$0.2 < z_h = \frac{P \cdot p_h}{P \cdot q} < 1$$



See Eduardo Rodriguez thesis and  
DESY-00-040, Phys.Lett. B481 (2000) 199-212

# Why pseudorapidity?

$$E_T^{\text{BR}} = E_T^{\text{HCM}} = E_T^*$$

$$\phi^{\text{BR}} = \phi^{\text{HCM}}$$

For massless partons or particles

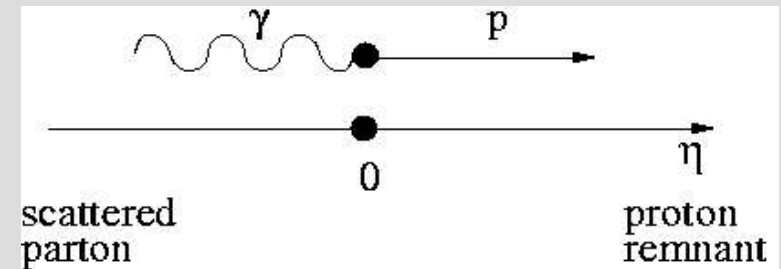
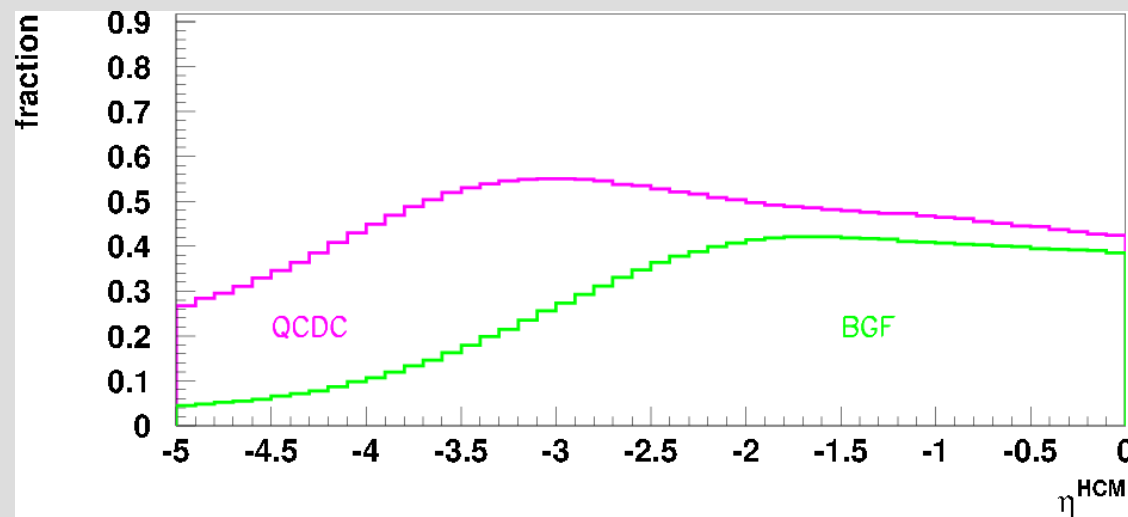
$$\eta^{\text{BR}} \approx \eta^{\text{HCM}} + 0.5 \cdot \ln\left(\frac{1-x}{x}\right)$$

$$z_h = \frac{E_T^*}{Q} \cdot e^{-\eta^{\text{BR}}}$$

$$\frac{E_T^*}{Q} = \sqrt{\frac{z_h(1-x_h)(1-z_h)}{x_h}}$$

$$\sqrt{\frac{x_h z_h}{(1-x_h)(1-z_h)}} = e^{\eta^{\text{BR}}}$$

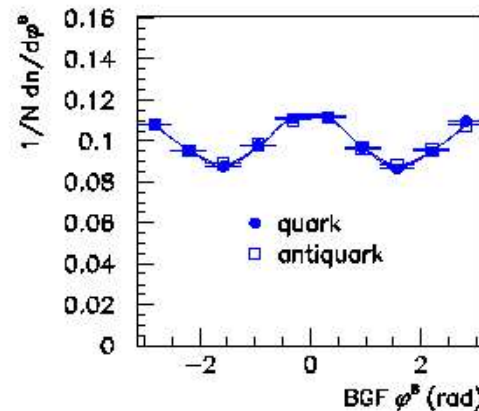
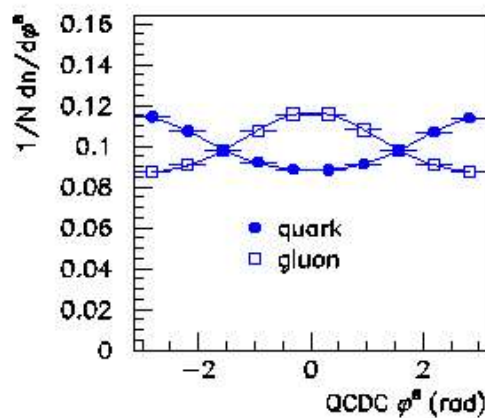
# Look into pseudorapidity



If integrated over

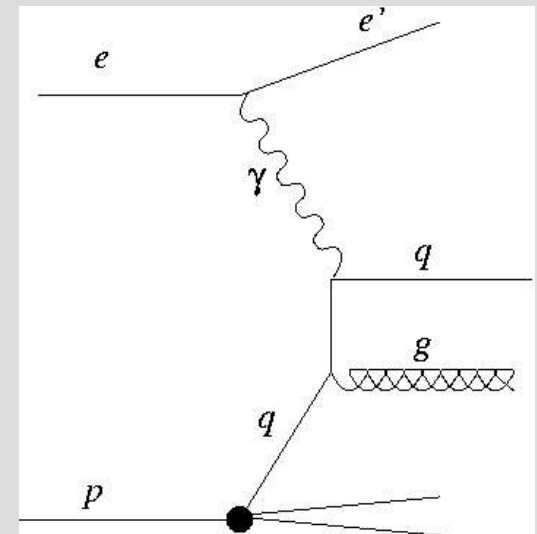
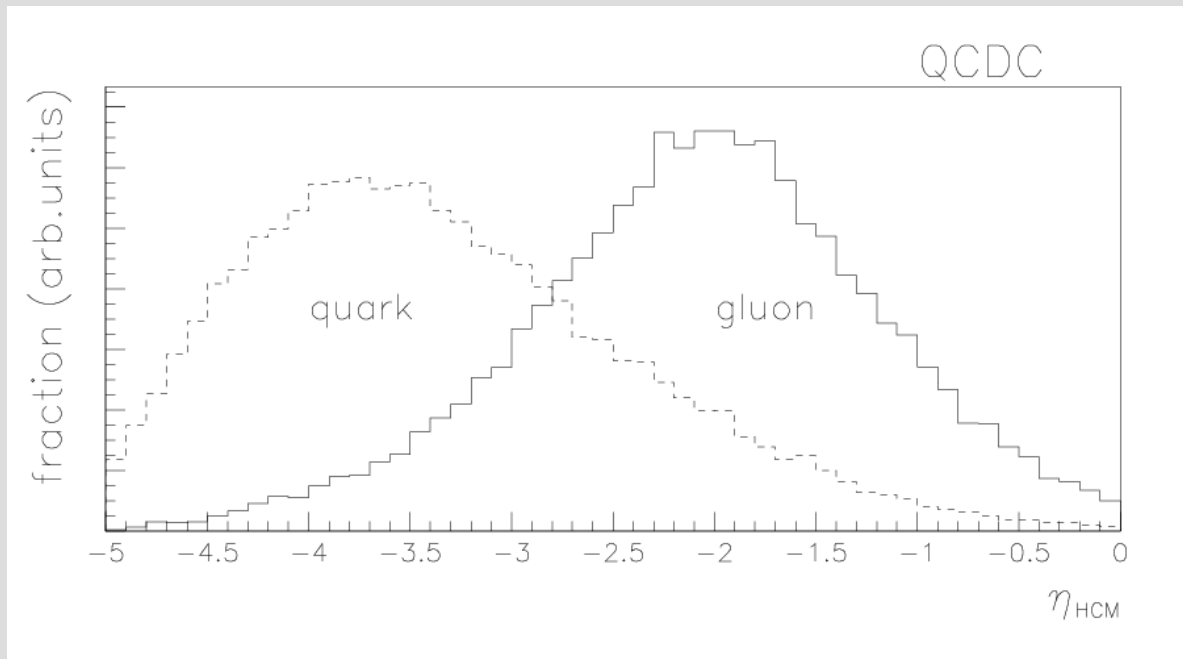
the dominant contribution:  
from QCDC  $\gamma^* q \rightarrow qg$   
to  $\cos(\phi)$

from BGF  $\gamma^* g \rightarrow q\bar{q}$   
to  $\cos(2\phi)$





# QCD Compton



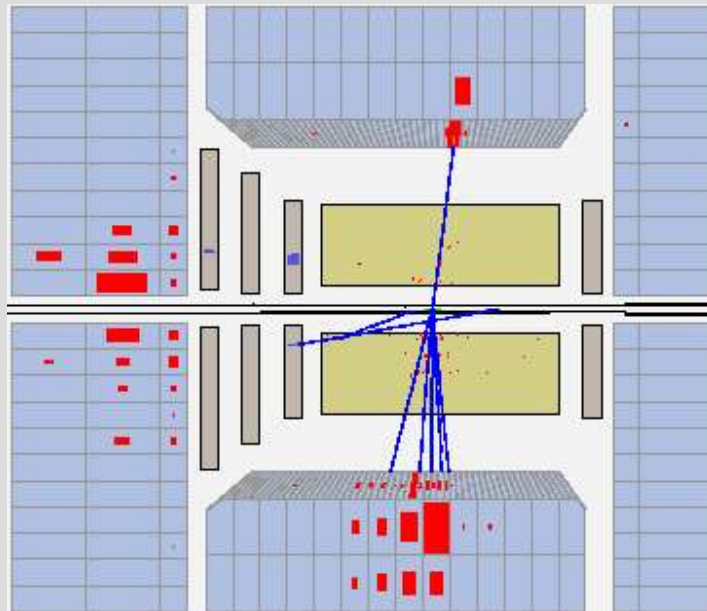
## **Pseudorapidity because:**

- separation of BGF and QCDC
- separation of hadrons from  $q$  and  $g$

## **Important because:**

- BGF  $\rightarrow$  gluonic structure function of proton

# Why energy flow method instead of multiplicity method?



Tracking detector

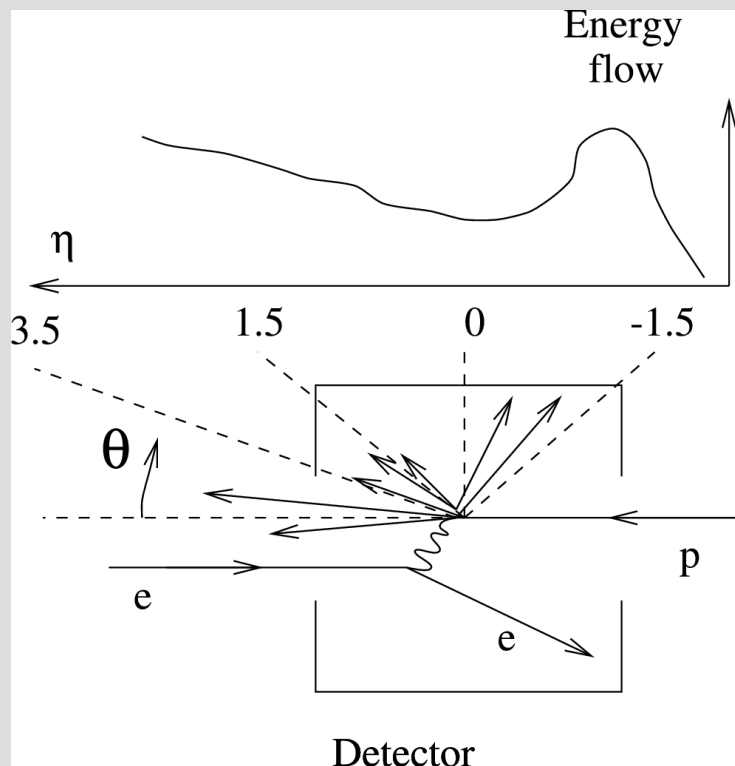
$$-1.75 < \eta^{\text{LAB}} < 1.75$$

Calorimeter

$$-3 < \eta^{\text{LAB}} < 4$$

# Energy flow method in the laboratory frame

ZEUS UFO



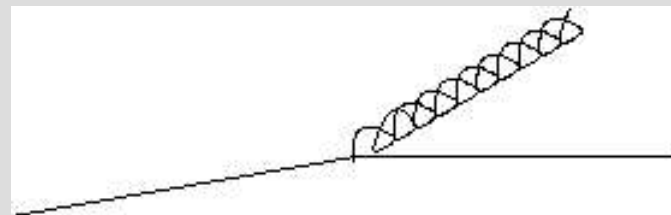
Energy flow objects EFO  
EFO used as pseudohadrons

pQCD

infrared and collinear  
singularities out

Peccei, Rückl (1978)

Each particle direction is  
weighet with its energy



# Comments on the energy flow method and pseudorapidity

## Advantages

- ▶ charged and neutral hadrons included
- ▶ hard partons ( $E_T^*$  larger) provides a larger contribution
- ▶ hadrons nearby in the HCM frame  $\rightarrow$  nearby in LAB
- ▶ sensitive to parton fragmentation  $\rightarrow$  no dependence on jet algorithms
- ▶ multiplicity method with charged hadrons  $\rightarrow$  sensitive to hadronization
- ▶ calorimeter energy scale is cancelled for i.e.  $\langle \cos(\phi) \rangle$

## Disadvantages

- ▶ no hadrons but clusters of energy  
the quantities like  $z = Pp_h / Pq$  are not well measured

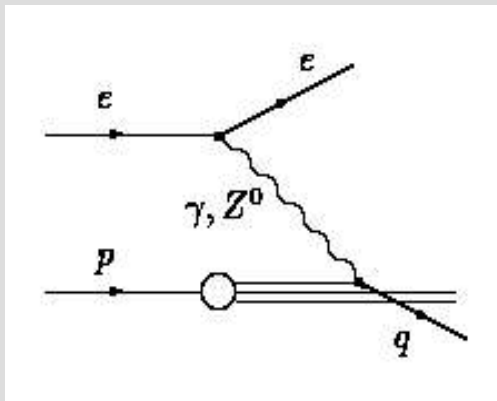
# Global selection criteria

$$E_e^{\text{LAB}} > 10 \text{ GeV}$$

$$100 < Q^2 < 8000 \text{ GeV}^2$$

$$0.2 < y < 0.8$$

$$0.01 < x < 0.1$$



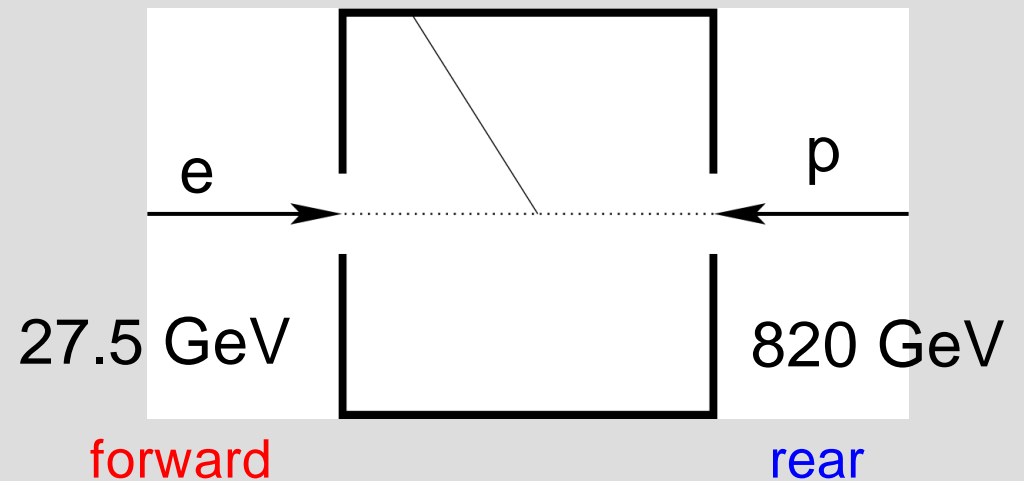
$$\theta_{\text{particle}}^{\text{LAB}} > 8^\circ \text{ (First ring)}$$

$$p_{\text{T}}^{\text{LAB}} > 150 \text{ MeV}$$

$$Q^2 = -q^2 = -(k_{e'} - k_e)^2$$

$$x = Q^2 / 2 P \cdot q$$

$$y = P \cdot q / P \cdot k_e$$





# Monte Carlo Models

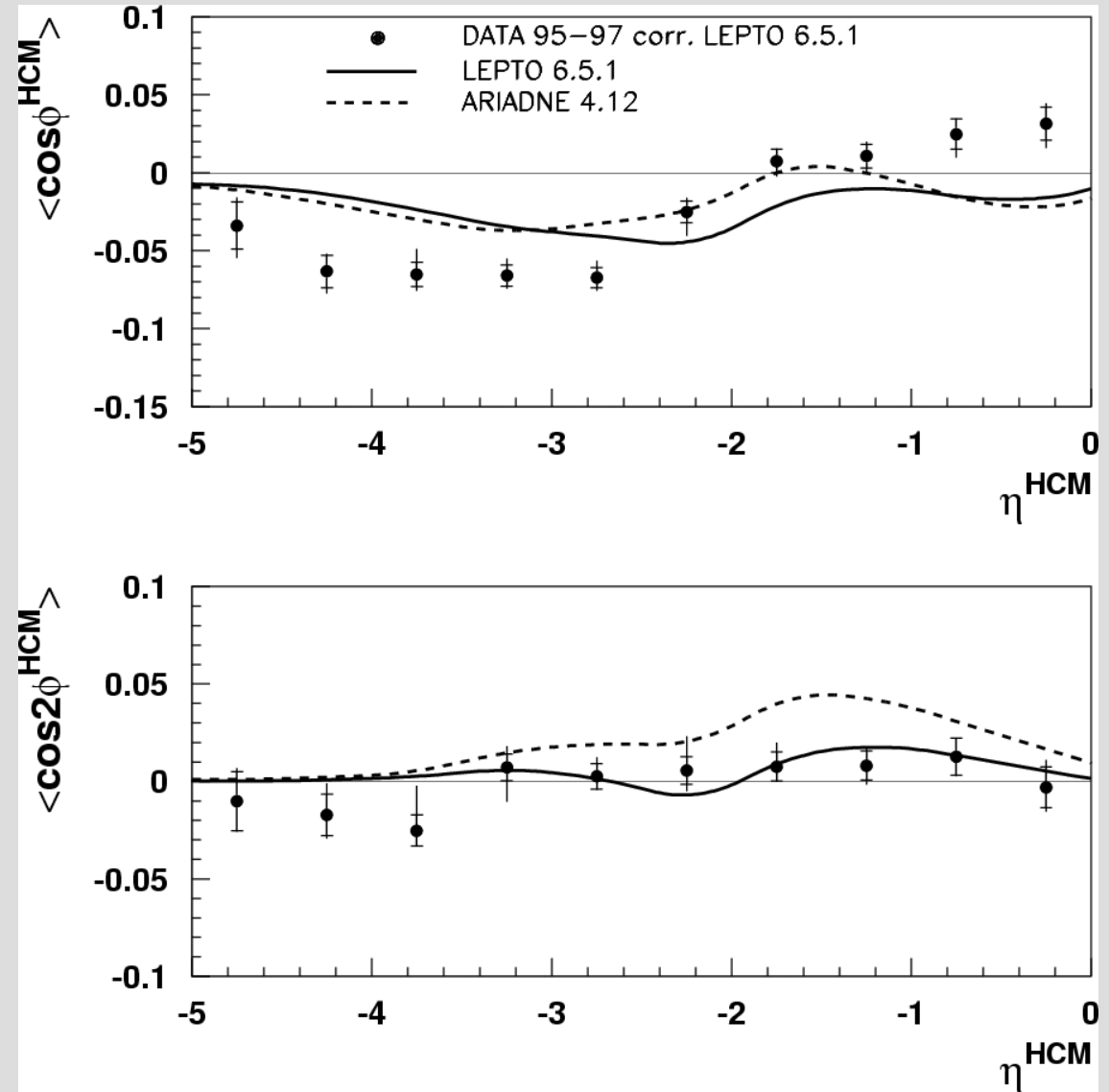
LEPTO 6.5.1 – matrix element and parton shower

ARIADNE 4.12 – colour dipol model (LO)

DISENT – NLO colour dipol model

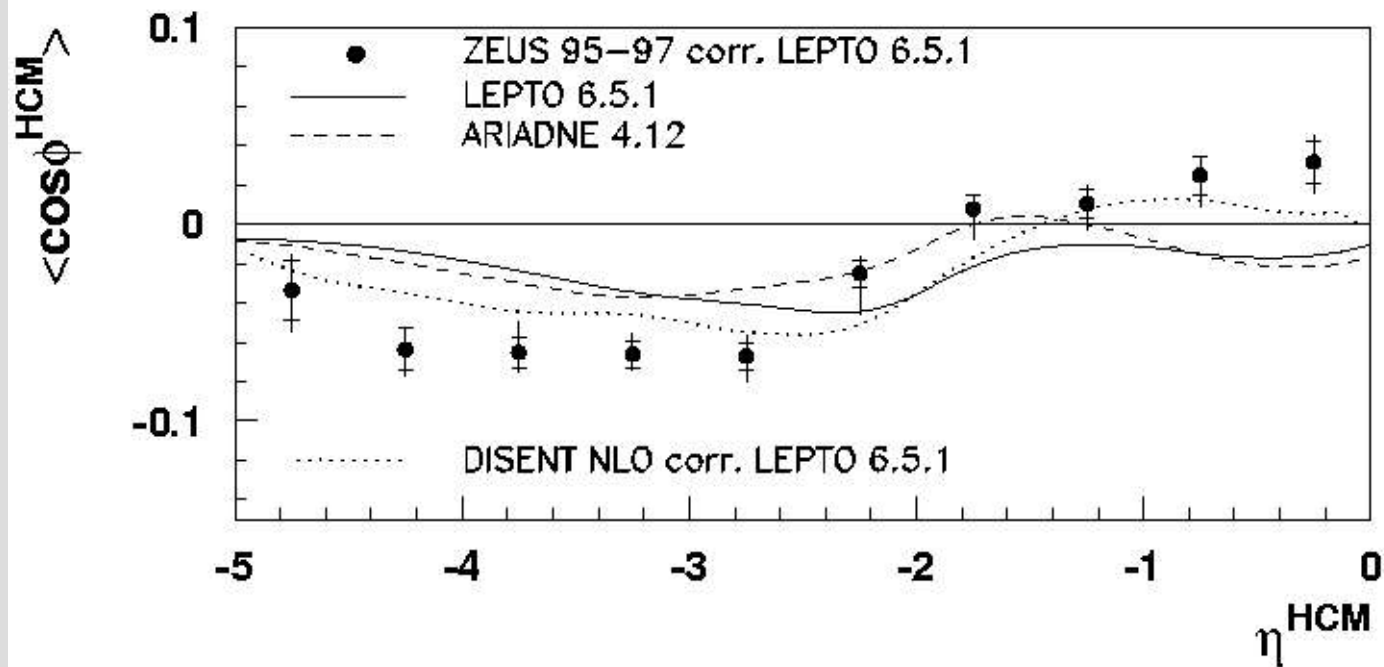
# DATA versus LO models

*Energy flow method*





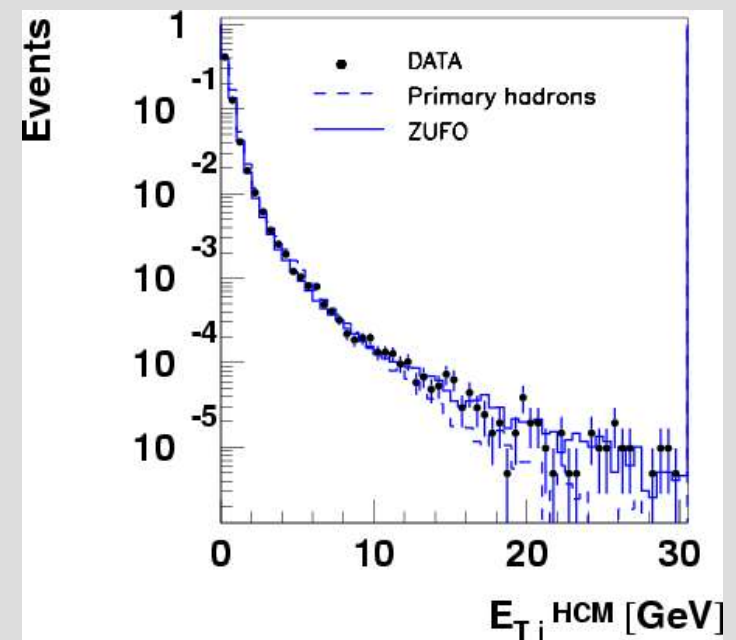
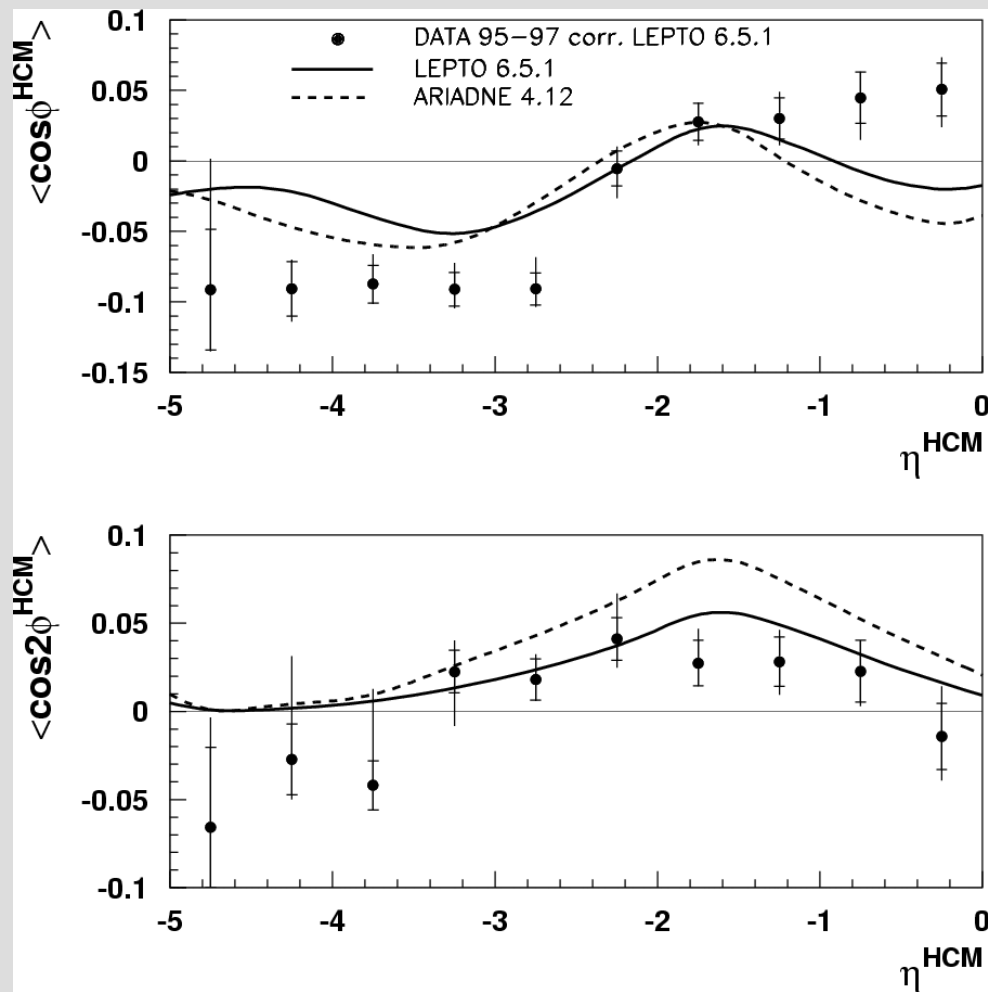
# Azimuthal asymmetry energy flow method



# Compared ZEUS DATA 95-97 with predictions

*Energy flow method*

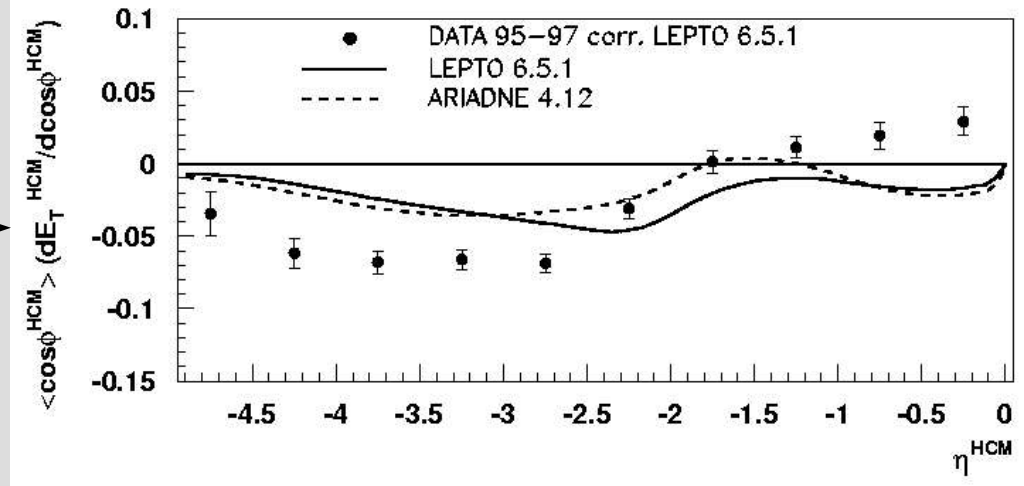
$E_T^{\text{HCM}} > 1 \text{ GeV}$



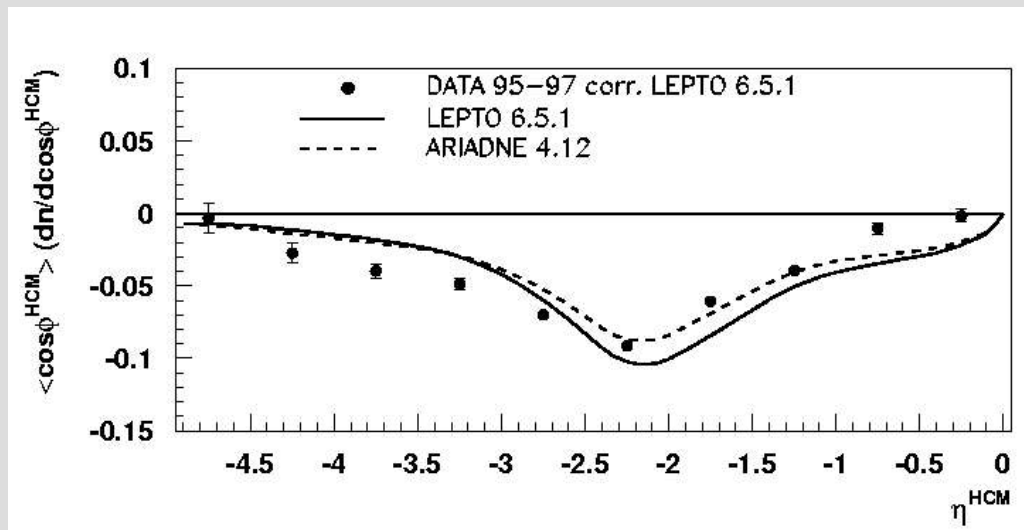


# Energy flow and multiplicity

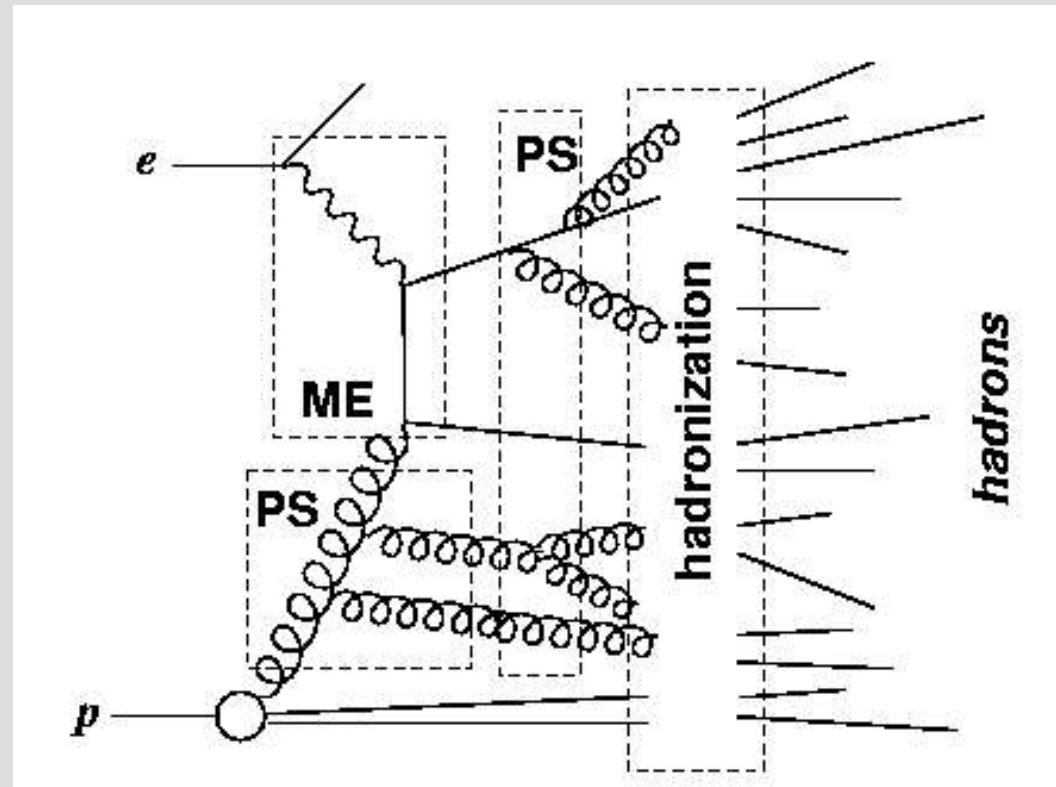
Energy flow method



Multiplicity method



# Sequence of subprocesses in DIS

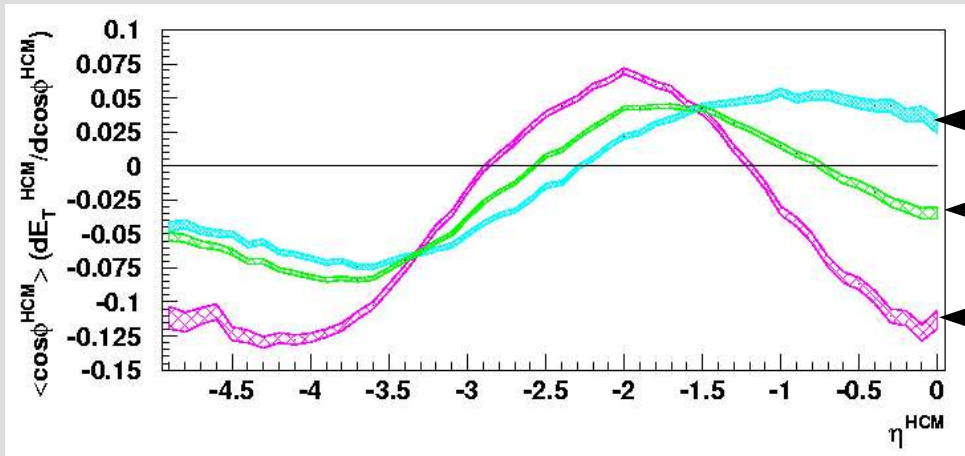


Contribution to asymmetries from: **BGF** + **QCDC**

# ME prediction:

# parton level without PS

## Energy flow method



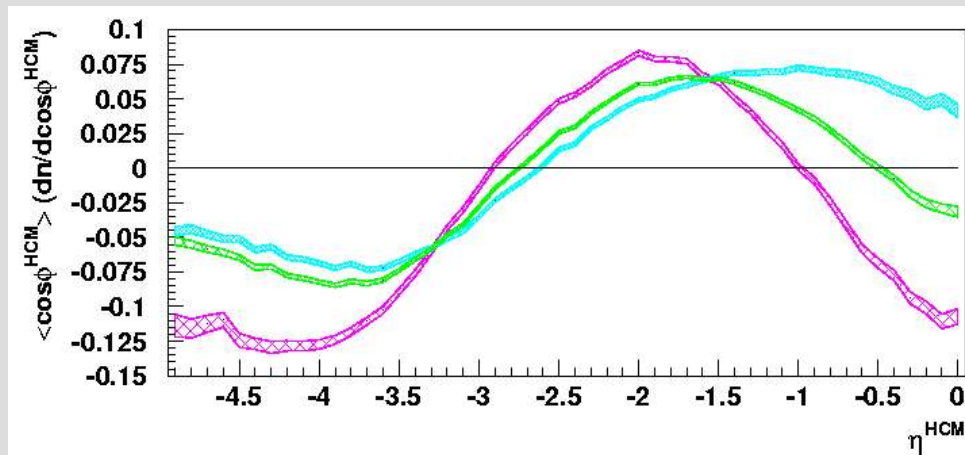
LEPTO 6.5.1

QCDC

BGF + QCDC + QPM

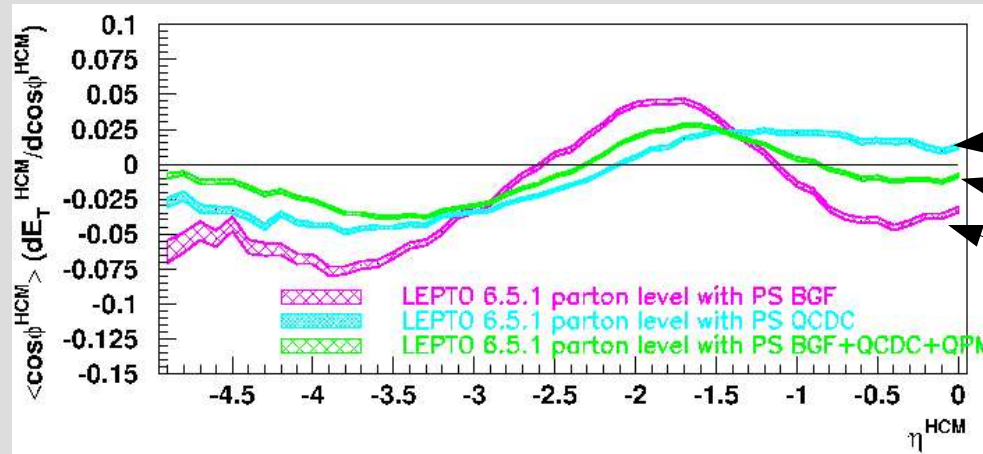
BGF

## Multiplicity method



# ME + PS prediction: parton level with PS

## Energy flow method



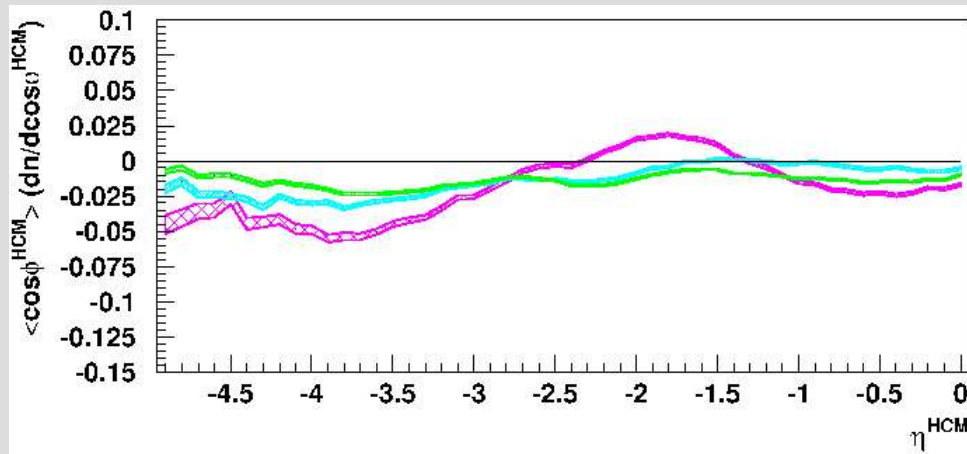
LEPTO 6.5.1

QCDC

BGF + QCDC + QPM

BGF

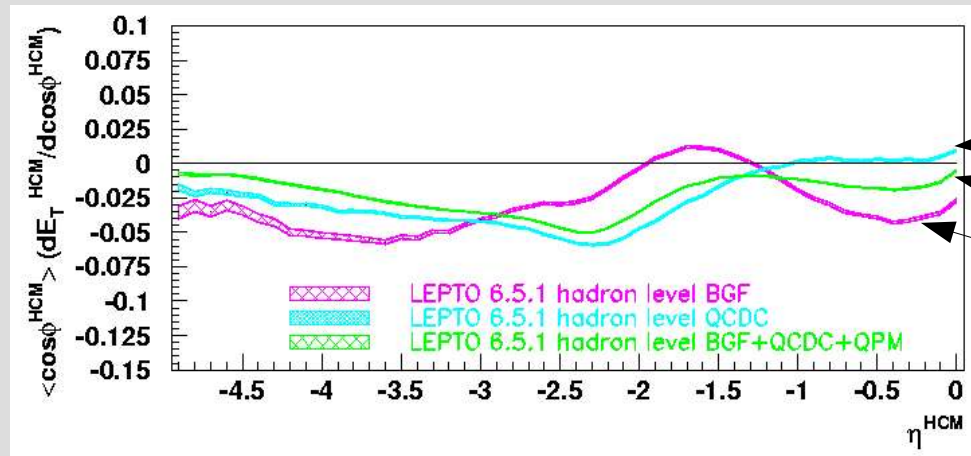
## Multiplicity method



# Prediction:

# hadron level + detection

## Energy flow method



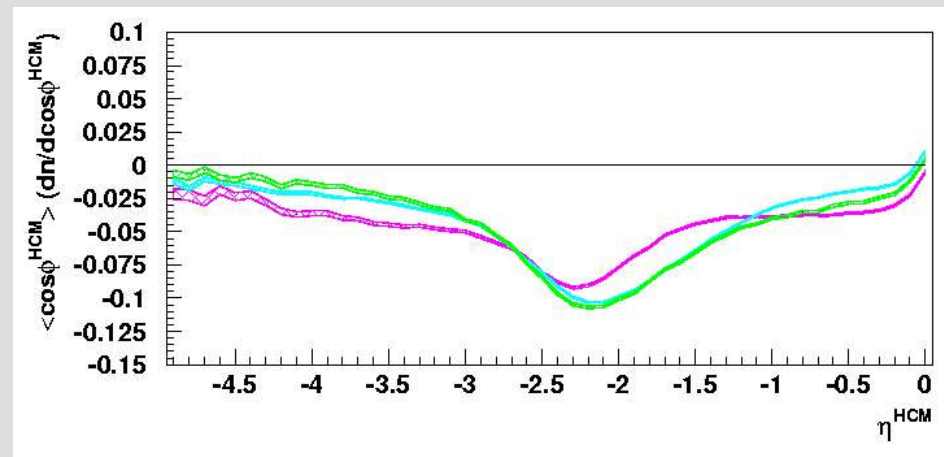
LEPTO 6.5.1

QCDC

BGF + QCDC + QPM

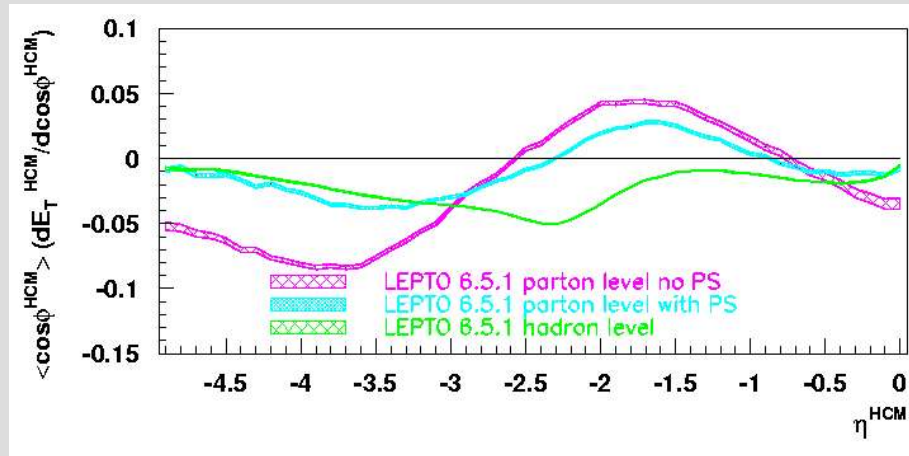
BGF

## Multiplicity method



# Prediction: parton level vs. hadron level

## Energy flow method



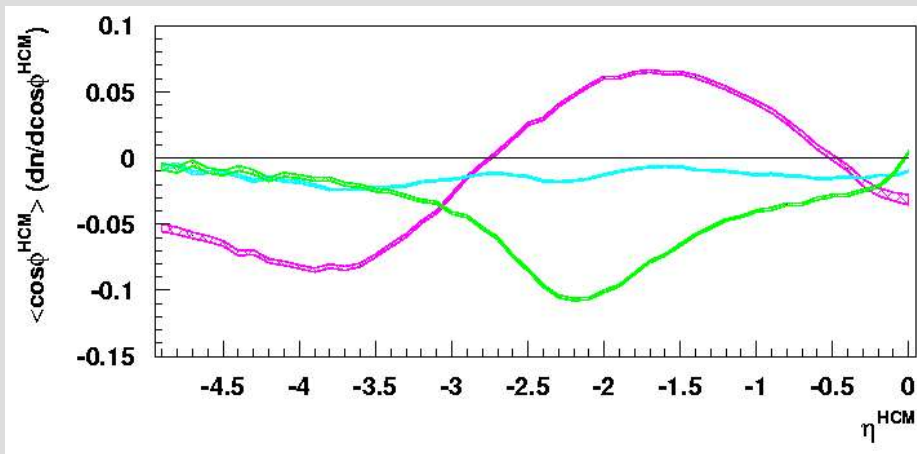
LEPTO 6.5.1

Parton level no PS

Parton level with PS

Hadron level

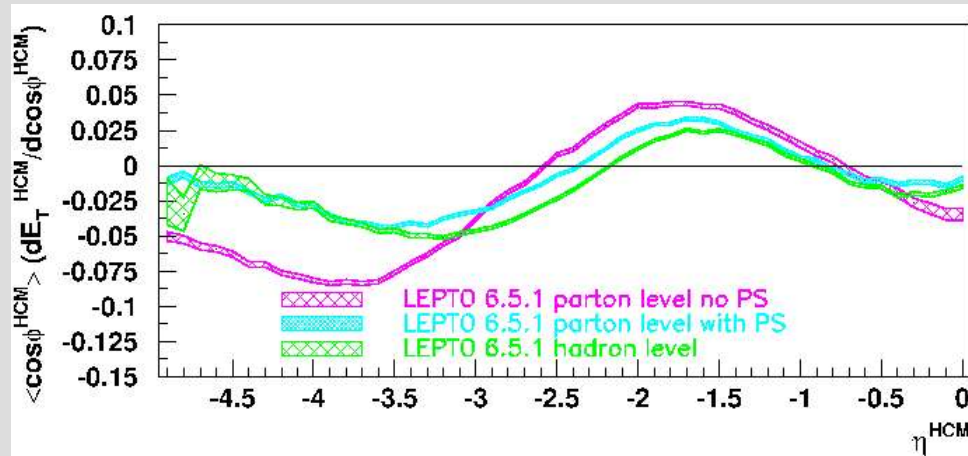
## Multiplicity method





# Prediction: parton level vs. hadron level

## Energy flow method



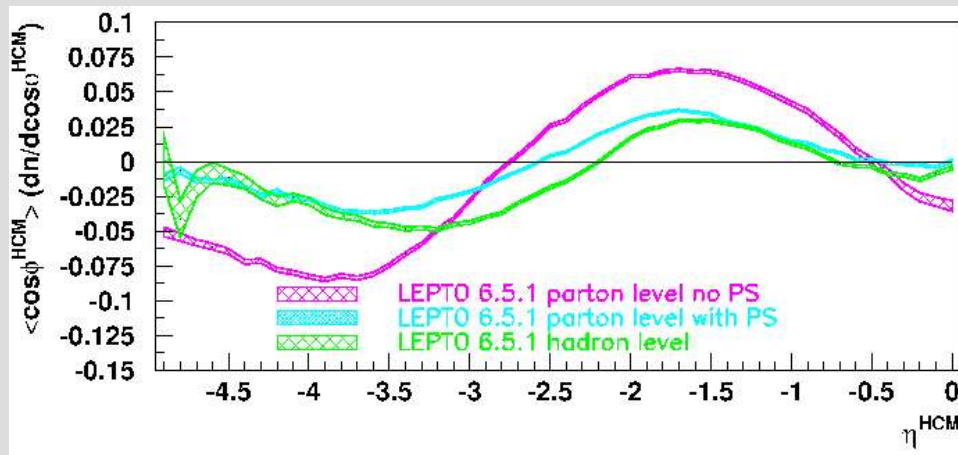
LEPTO 6.5.1

Parton level no PS

Parton level with PS

Hadron level

## Multiplicity method



$$E_T^{HCM} > 1 \text{ GeV}$$

# Conclusions (part I)

- The multiplicity method is very sensitive to detection inefficiencies
- Therefore the obtained azimuthal asymmetries are dominated by hadronization effects
- Azimuthal asymmetries for  $0.01 < x < 0.1$  with energy flow the dominant contribution depends on the  $\eta$  interval
  - $-5. < \eta^{\text{HCM}} \leq -3.$  is dominated by QCDC and NLO
  - $-3. < \eta^{\text{HCM}} \leq -1.$  is sensitive to BGF content
  - $-1. < \eta^{\text{HCM}} \leq -0.$  by both QCDC + BGF and NLO



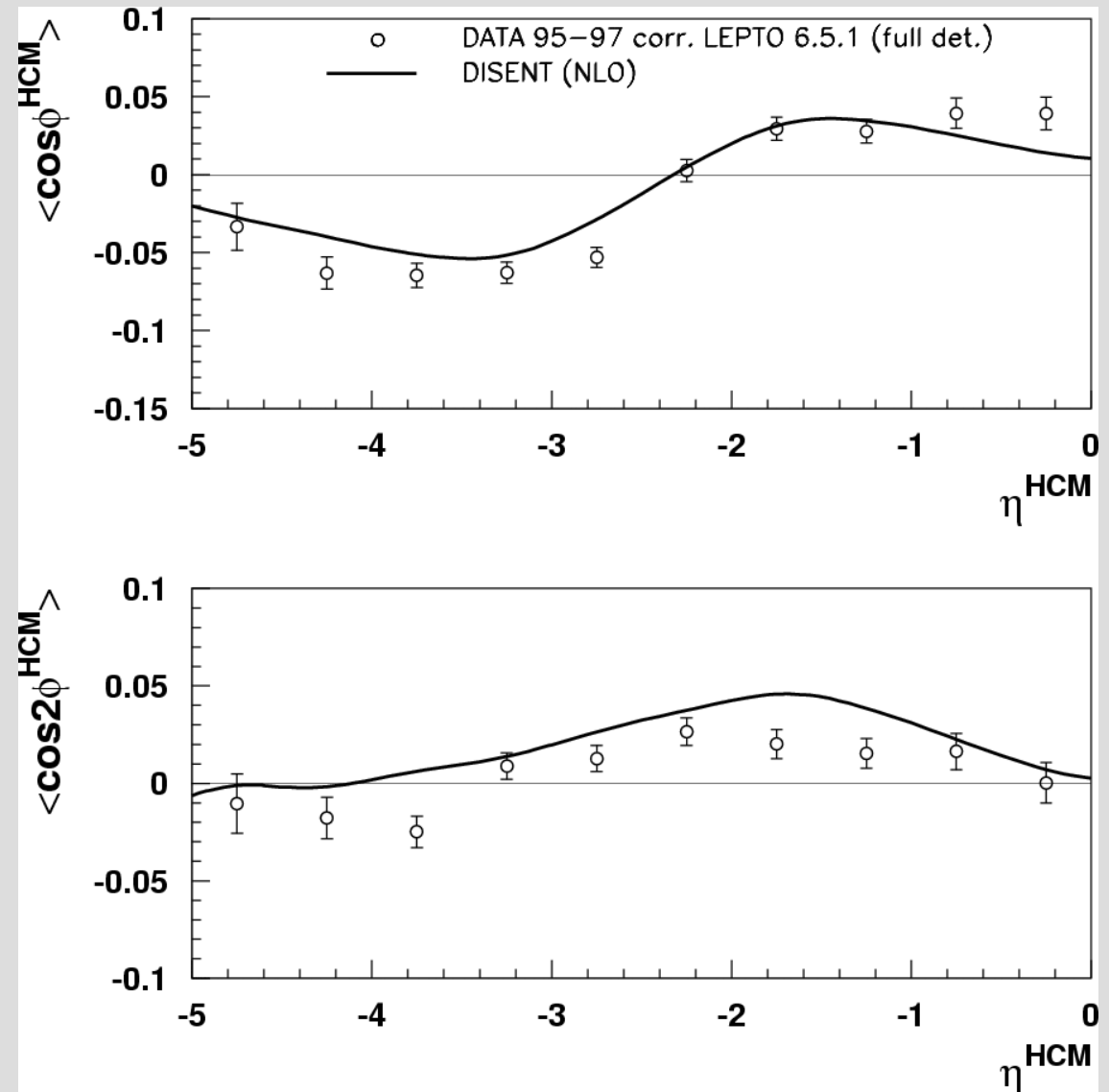
# DATA versus NLO

## *Energy flow method*

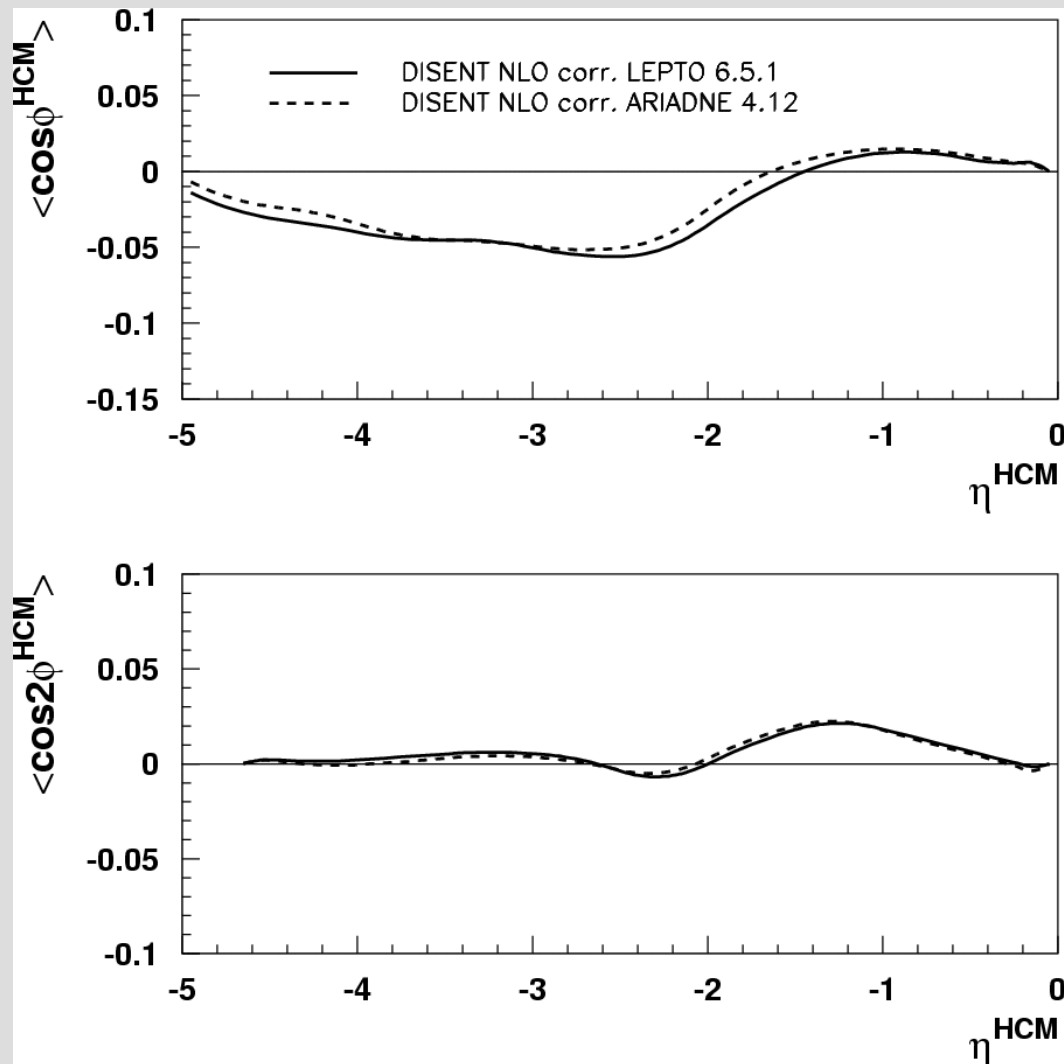
DATA corrected  
for losses  
in the HERA beampipe

(Against rules!)

Stat. uncertainties only



# Corrections for DISENT NLO

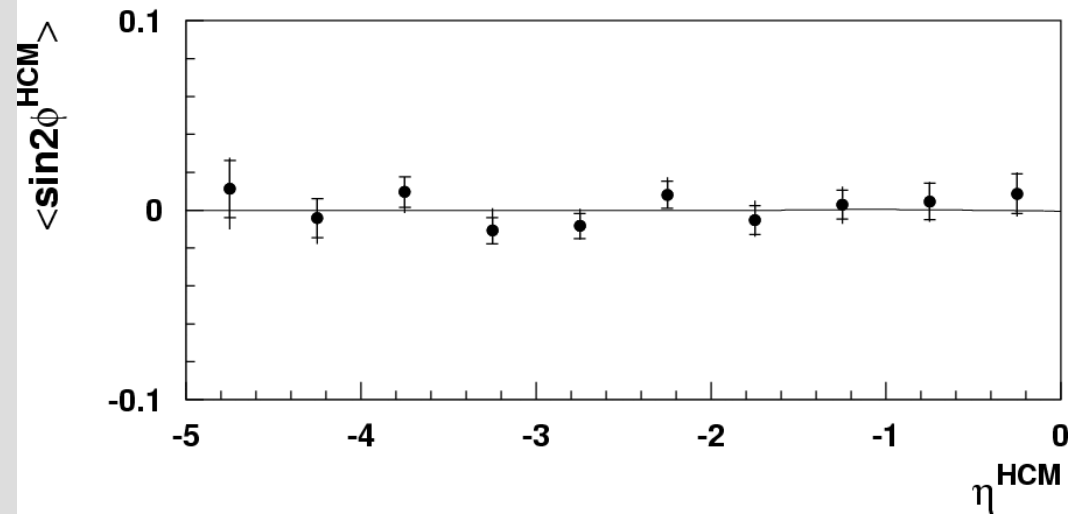
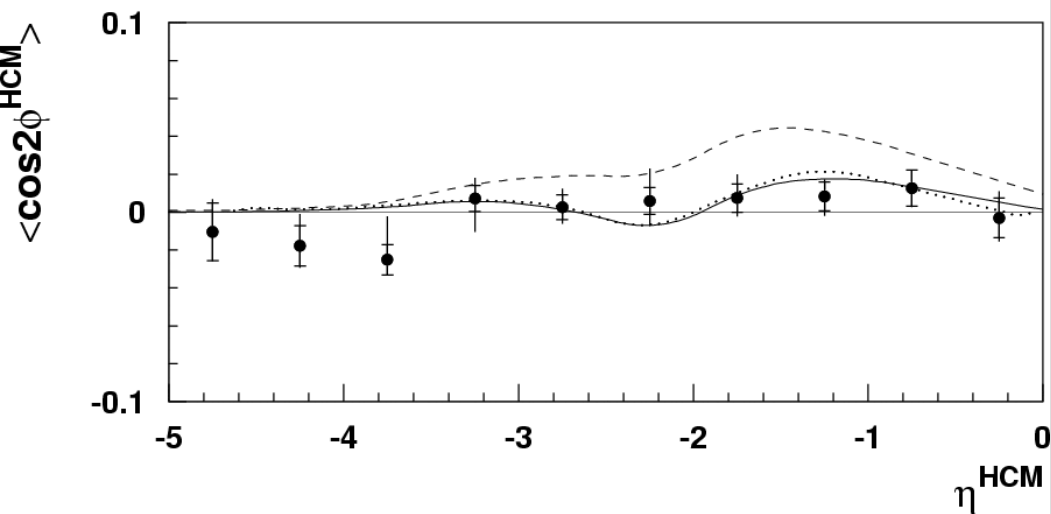
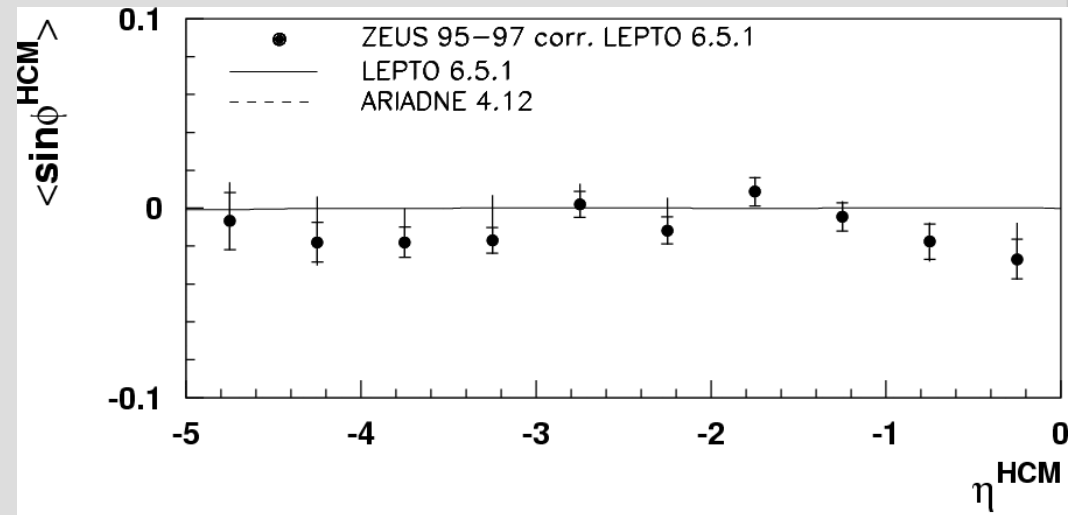
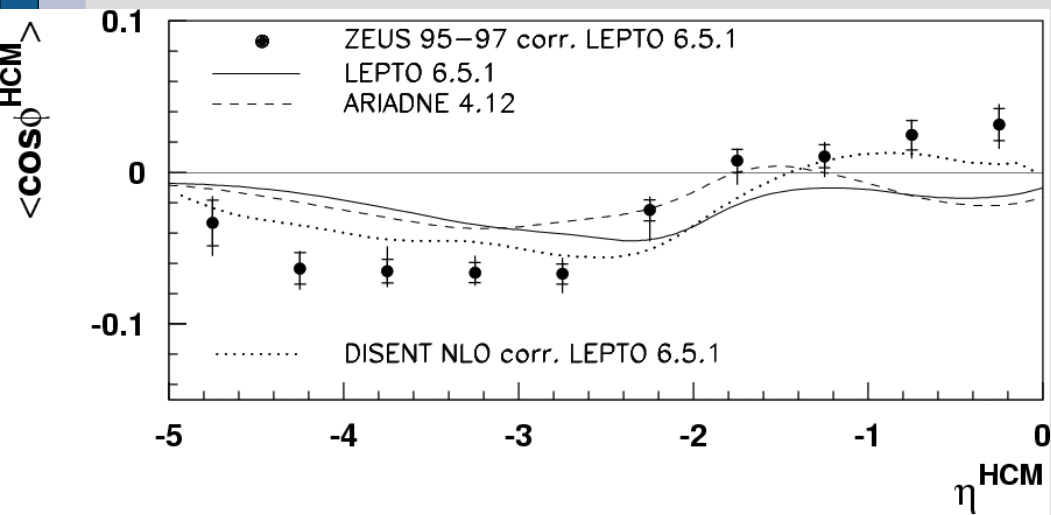


No model dependence

Uncertainties due to factorization and renormalization scales were calculated by changing  $\mu_{F,R} = Q$  to  $0.5 \cdot Q$  and  $2 \cdot Q$

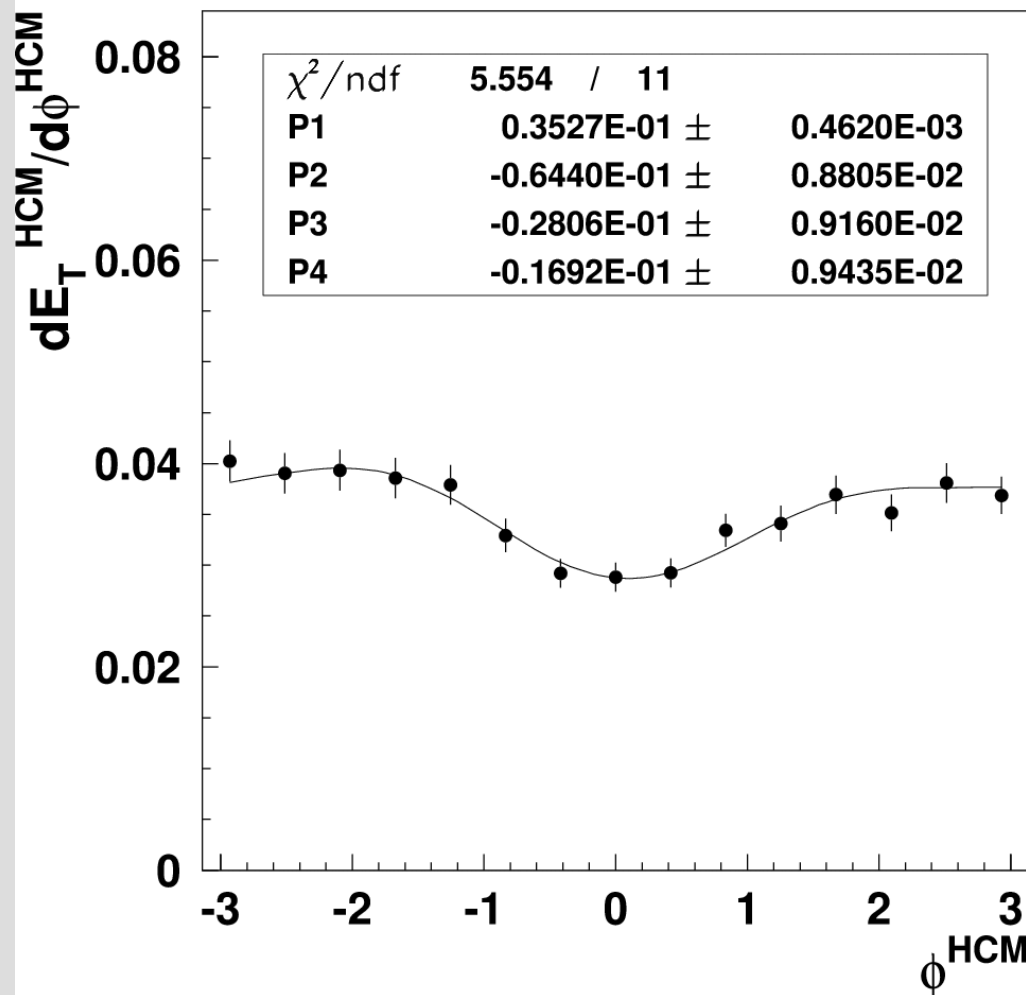
The asymmetries fluctuated within a band of  $\pm 0.005$

# Azimuthal asymmetry energy flow method



# Azimuthal angle distribution energy flow method

$$-4. < \eta^{\text{HCM}} \leq -3.5$$



DATA 95-97 corr. LEPTO 6.5.1

$$\frac{d\sigma}{d\phi} = P_1 + P_2 \cos(\phi) + P_3 \cos(2\phi) + P_4 \sin(\phi)$$

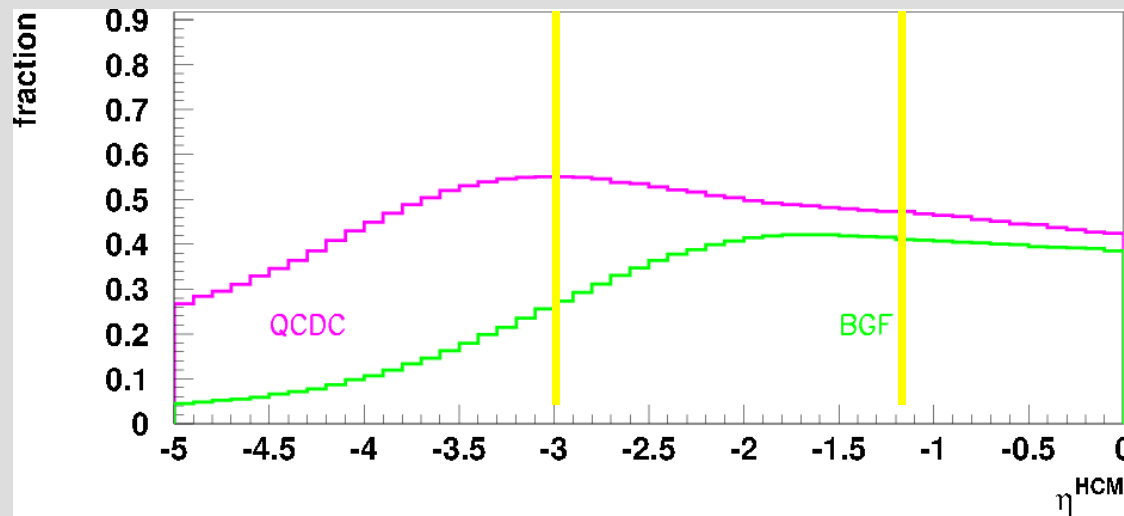
The second independent analysis  
Fitted values  $P_2, P_3, P_4$   
are consistent with mean ones

$$\langle \cos(\phi^{\text{HCM}}) \rangle, \langle \cos(2\phi^{\text{HCM}}) \rangle, \langle \sin(\phi^{\text{HCM}}) \rangle$$



# BGF versus QCDC for energy flow method

## LEPTO 6.5.1



charged hadrons  
(multiplicity method)  
Eduardo Rodrigues  
DESY-00-040  
Phys.Lett. B481 (2000) 199-212

jet method  
Oscar Gonzalez  
DESY-02-171  
Phys.Lett. B551 (2003) 226-240



jet method



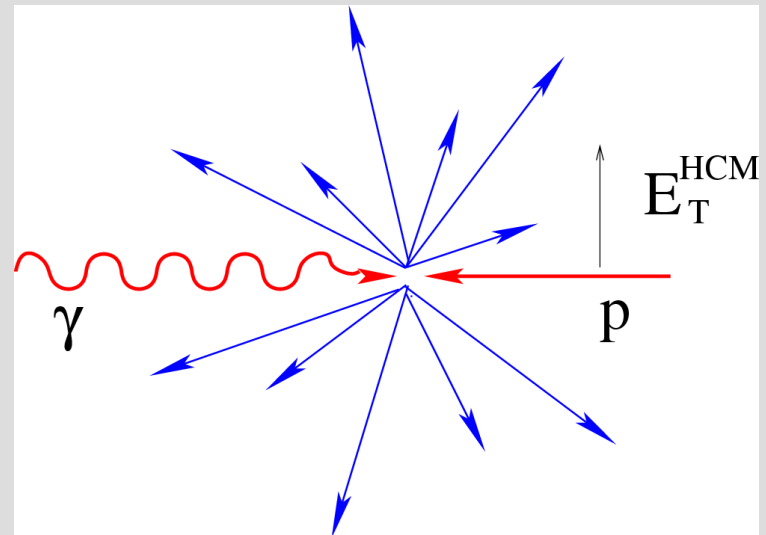
charged hadrons  
(multiplicity method)

# Azimuthal asymmetry as $f(E_T^{\text{HCM}}(\text{min}))$

$$-5. < \eta^{\text{HCM}} \leq -2.5$$

$$-2.5 < \eta^{\text{HCM}} \leq -1.$$

$$-1. < \eta^{\text{HCM}} \leq 0.$$



$E_T^{\text{HCM}}(\text{min})$  large  $\rightarrow$  pQCD

Theorists claim that cut on  $E_T^{\text{HCM}} \geq 2 \text{ GeV}$   
is sufficient to removed nonperturbative effects

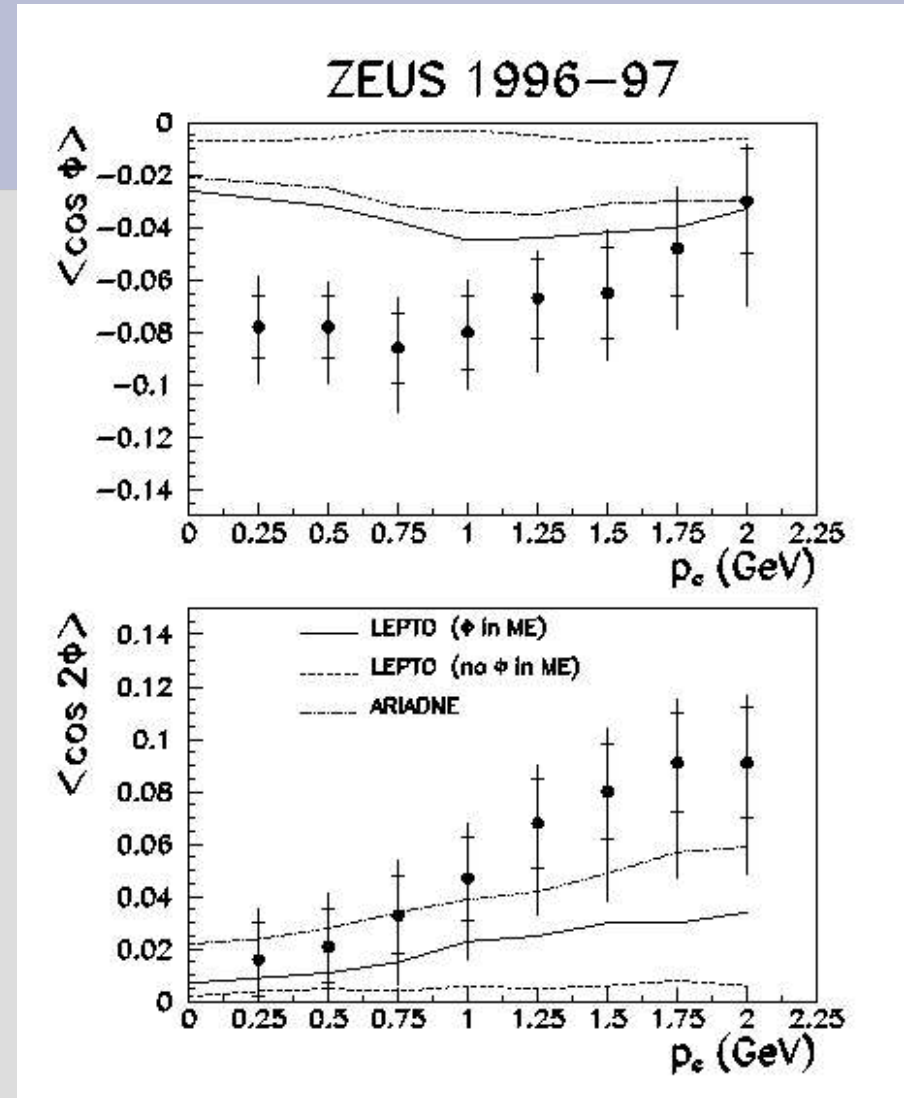
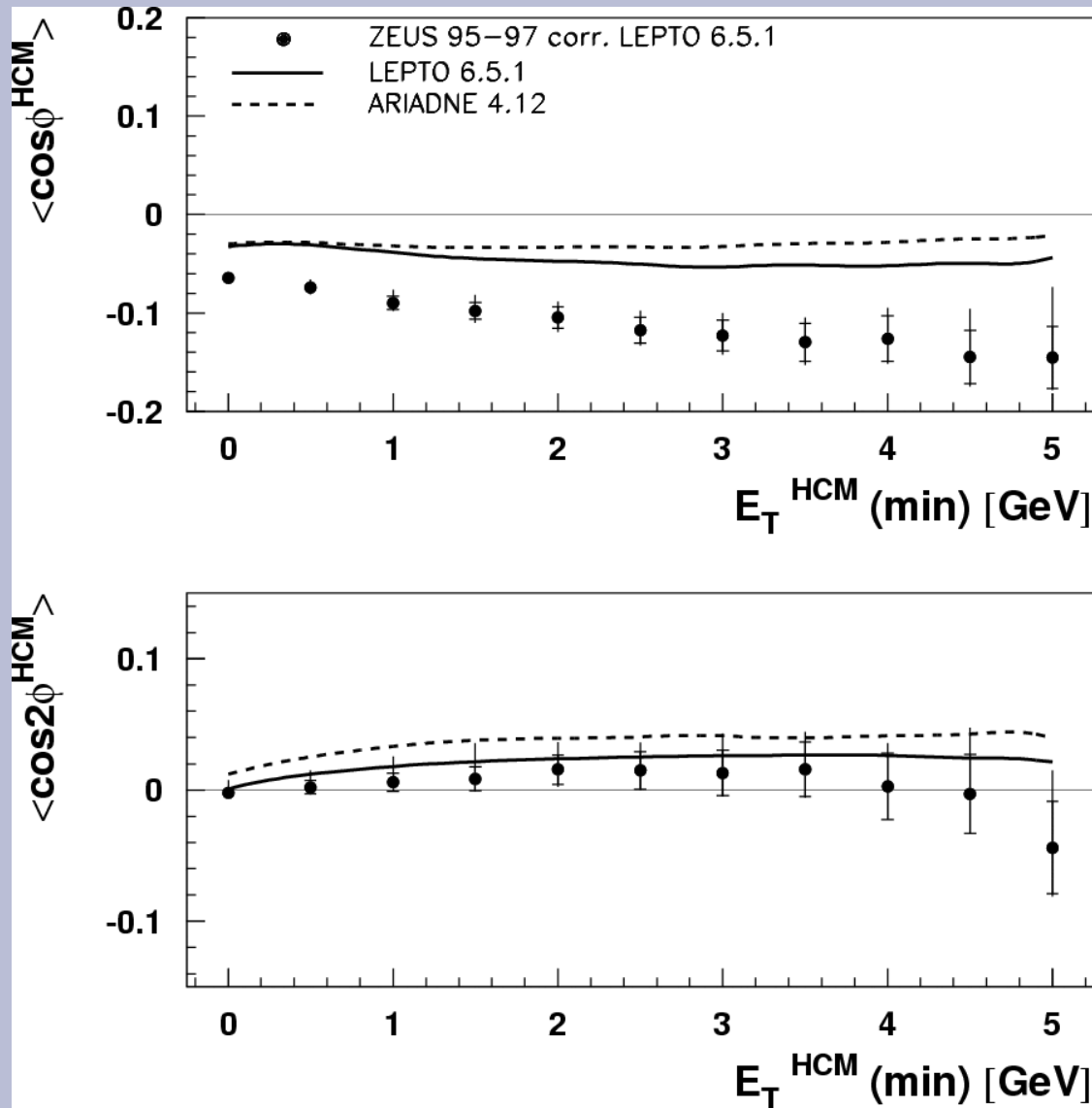
Method  $p_T$  - cut:

Chay et al. Phys.Rev. D45 (1992) 46





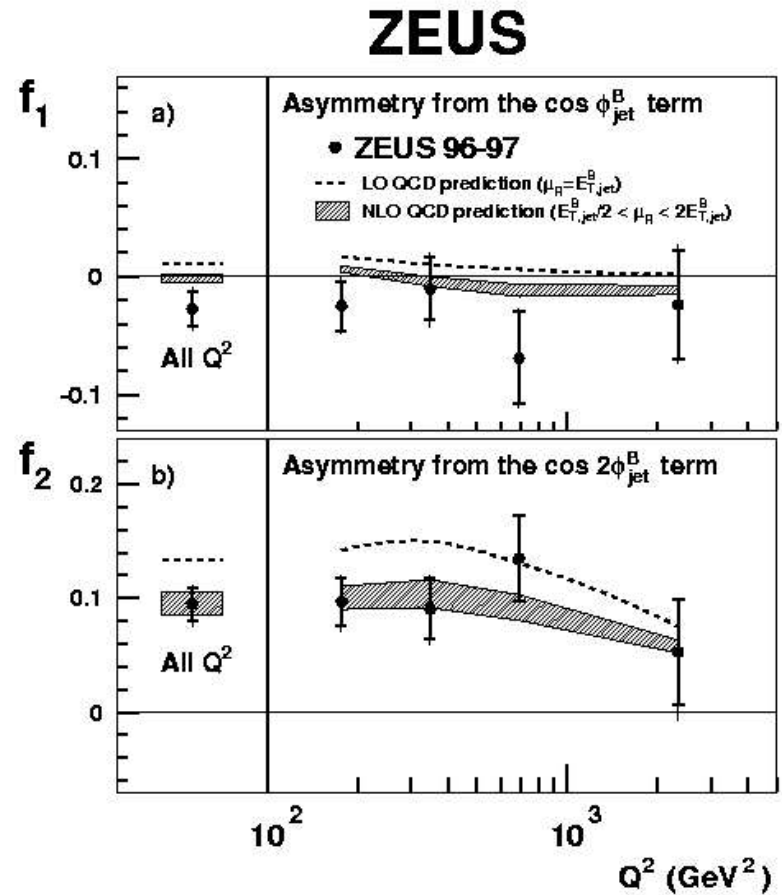
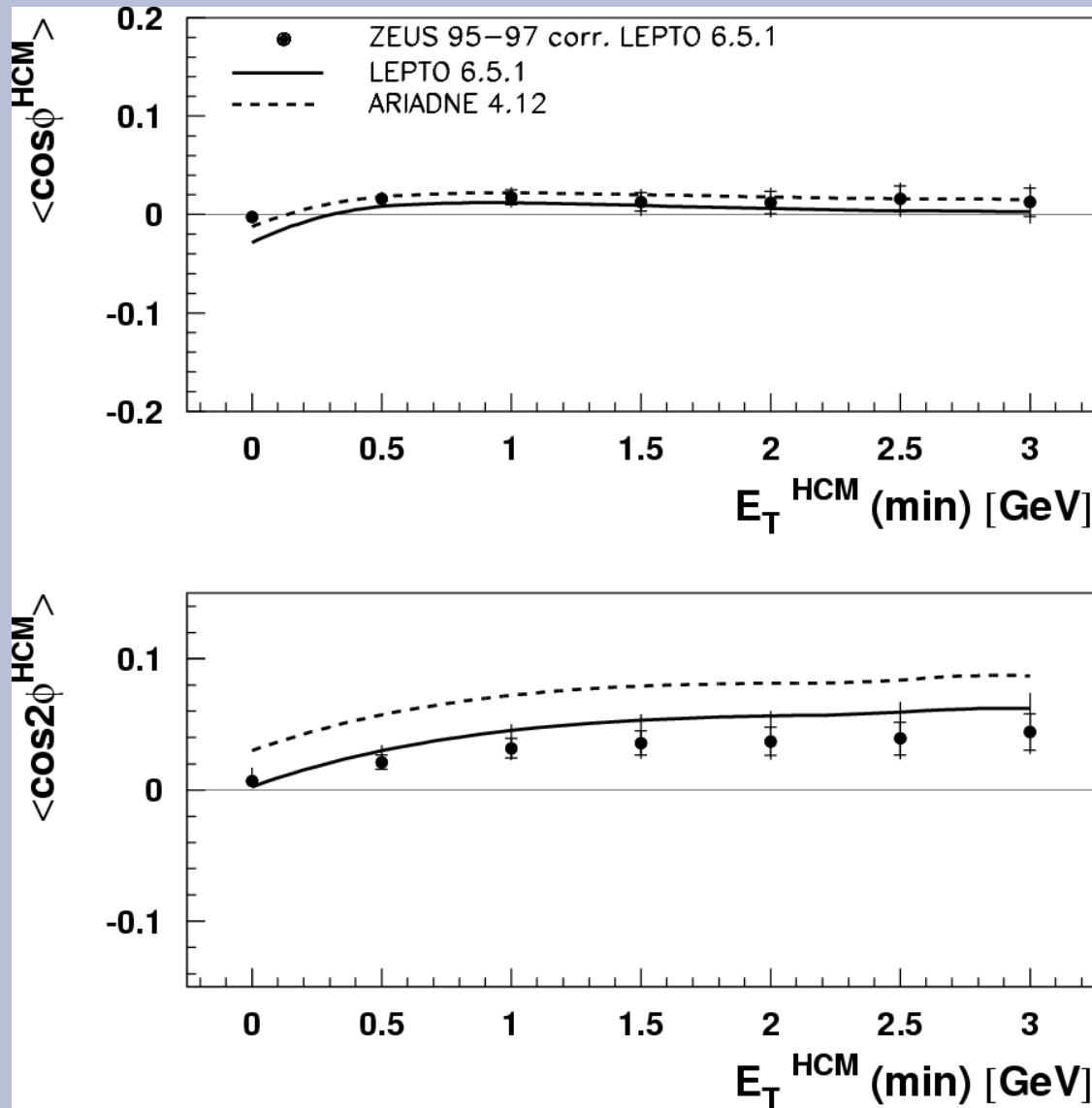
$$-5. < \eta^{\text{HCM}} \leq -2.5$$



$$0.2 < z_h = \frac{P \cdot p_h}{P \cdot q} < 1$$

$$|\eta^{\text{LAB}}| < 1.75$$

$$-2.5 < \eta^{\text{HCM}} \leq -1.$$

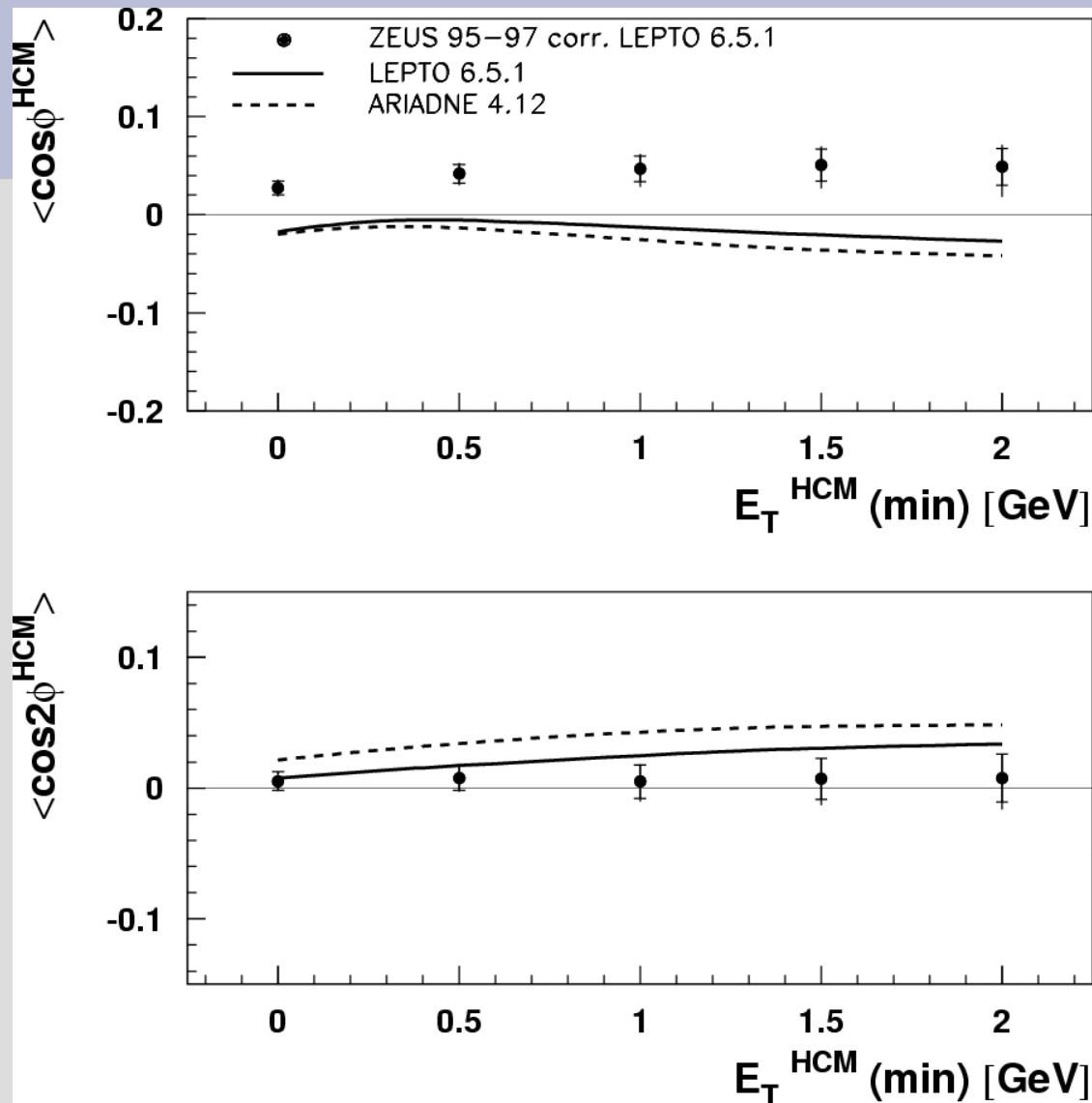


$$\frac{1}{\sigma} \frac{d\sigma}{d|\phi_{\text{jet}}^B|} = \frac{1}{\pi} [1 + f_1 \cos(\phi_{\text{jet}}^B) + f_2 \cos(2\phi_{\text{jet}}^B)]$$

$$f_1 \Leftrightarrow \langle \cos(\phi^{\text{HCM}}) \rangle; \quad f_2 \Leftrightarrow \langle \cos(2\phi^{\text{HCM}}) \rangle$$

$$-1. < \eta^{\text{HCM}} \leq 0.$$

First measurements  
in this region  
of phase space



NLO prediction expects:

- the positive values for  $\langle \cos(\phi^{\text{HCM}}) \rangle$
- the smaller positive ones for  $\langle \cos(2\phi^{\text{HCM}}) \rangle$

# Summary and conclusions

A novel approach to azim. asym. is proposed which provides precise measurements and small systematic errors in the wider interval of phase space

## The method permits to:

- include charged and neutral hadrons
- enhance contributions of hard partons by weighting with energy, i.e. energy flow
- investigate contributions of BGF w.r.t QCDC
- compare these results with the previous ZEUS measurements

## Disadvantage:

- cuts on individual hadrons – unreliable



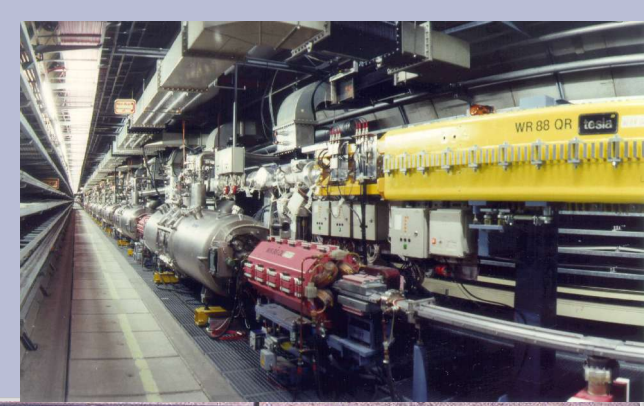
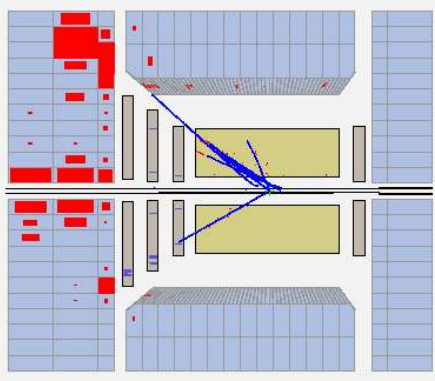
# Summary and conclusions

## The main results are:

- the NLO effects give not negligible contribution
- they provide better agreement with experimental data
- some small discrepancies are visible which cannot be explained by experimental errors



# New look at HERA



HERA jest polską marką odzieżową powszechnie rozpoznawalną i cenioną. Dzięki wysoko wykwalifikowanej kadrze oraz nowoczesnym liniom produkcyjnym, HERA oferuje produkty na najwyższym światowym poziomie. Wieloletnie doświadczenie poparte obserwacją aktualnych światowych trendów mody, pozwoliły Pani Halinie Widlińskiej na tworzenie kolekcji spójnych wzorniczo i kolorystycznie z obowiązującymi liniami największych domów mody. Współpraca z wiodącymi producentami tkanin i dodatków nadaje wyrobom oryginalny i lekki charakter. Kolekcje łączą w sobie zarówno elementy awangardowe, tworzone z nowatorskich oryginalnych materiałów jak i klasyczne całodzienne kostiumy dla kobiet aktywnych zawodowo. Dwa razy do roku HERA wprowadza na rynek nowe kolekcje obejmujące sezony: jesienno-zimowy oraz wiosenno-letni. Bardzo krótkie serie poszczególnych modeli połączone z szeroką gamą wzorniczą pozwalają klientkom firmy na kreowanie własnego, niepowtarzalnego wizerunku, zgodnego ze światowymi trendami.

Myślą przewodnią tworzonych przez firmę HERA kolekcji jest nadanie klientkom charakteru nowoczesności, lekkości, oryginalności a przede wszystkim poczucia wiecznej młodości.

