

Non-standard signatures in GMSB

Małgorzata Kazana, *Lisbon*

Alberto de Min, Giovanni Franzoni, *Milano*

Łukasz Gościło, Piotr Traczyk, *Piotr Zalewski*, *Warsaw*

OUTLINE:

- GMSB – introduction
- Why GMSB is interesting from the detection point of view?
- What was done by ATLAS?
- Why is not so simple?
- Discussion LEP–TeVatron–LHC–NLC

Gauge Mediated Supersymmetry Breaking

SUSY is broken in a so called **Secluded Sector** at scale \sqrt{F} and transmitted to **SM particles** via **Messenger Sector**.

Messenger sector contains n flavours of vector GUT representations (eg. SU(5)). Via interaction with **SM singlet scalar field from Secluded sector Messengers** acquire mass M and soft supersymmetry breaking masses squared proportional to $F_M < F$.

This effect feeds down to the MSSM through loop corrections. Gauginos acquire majorana masses at one loop

$$m_j(M) = n \frac{\alpha_j(M)}{4\pi} \Lambda$$

while sfermions acquire SUSY-breaking masses-squared at two loops

$$m_{f_i}^2(M) = 2n \sum_{j=1}^3 C_{ij} \left(\frac{\alpha_j(M)}{4\pi} \right)^2 \Lambda^2$$

where $\Lambda = F_M/M$ sets the effective scale of SUSY breaking in MSSM.

The superpartner mass depends on its SM coupling. Eg. the heaviest scalars are \tilde{q}_L , then $\tilde{q}_R, \tilde{l}_L, \tilde{l}_R$. Gaugino masses are related to sfermion masses via factor \sqrt{n} .

Because squarks of different families are degenerate at relatively low scale M , there is no large FCNC effects.

The splitting of masses is obtained via RGE and depends on $\tan\beta$ the hierarchy M/Λ and a value of the μ parameter (μ^2 is determined by EW symmetry breaking: only sign of μ is free). The most important difference between the GMSB and the mSUGRA models is that in the GMSB scenario gravitino is the LSP

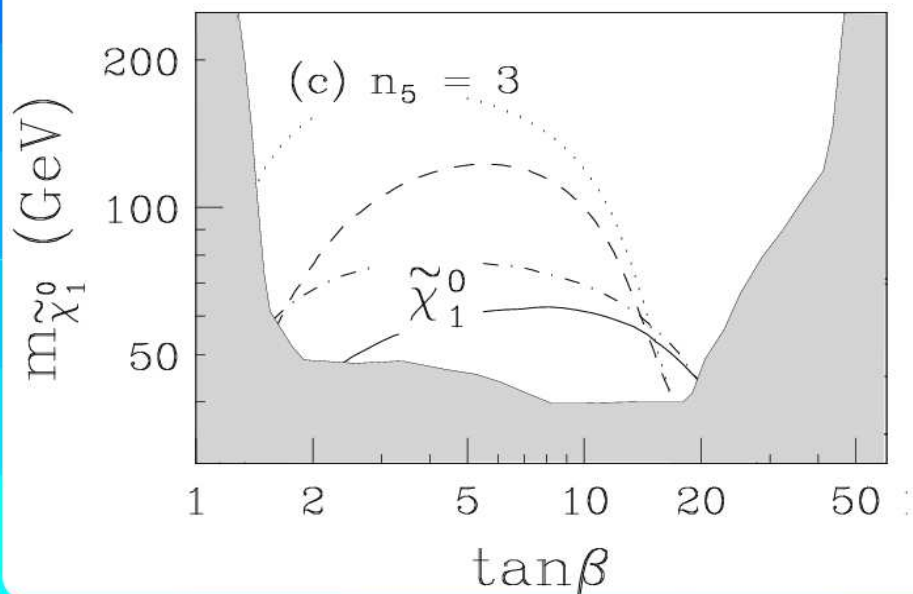
$$m_{\tilde{G}} = \frac{F}{\sqrt{3}M_P}$$

MSSM LSP becomes GMSB NLSP and decays to its SM partner and gravitino with $c\tau$ from microns to hundreds of meters. For $n = 1$ and low $\tan\beta$ neutralino is the NLSP whereas for high $\tan\beta$ or high n stau is the NLSP.

It can be seen below for $n = 3$

NLSP lifetime is proportional to the gravitino mass squared for fixed NLSP mass.

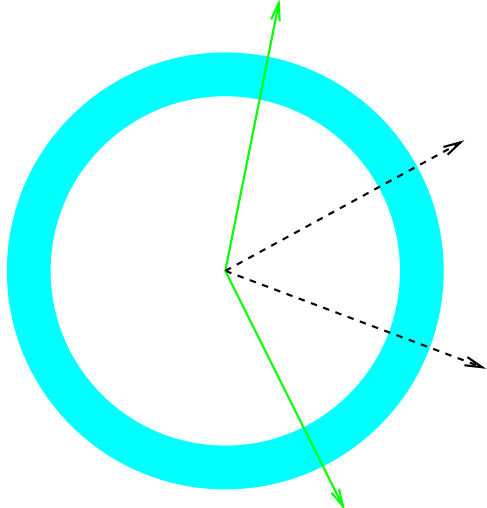
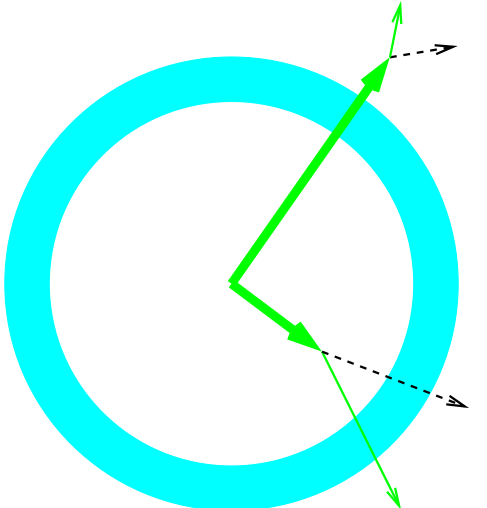
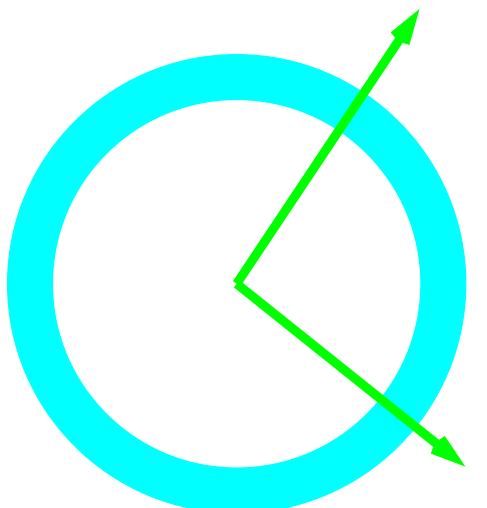
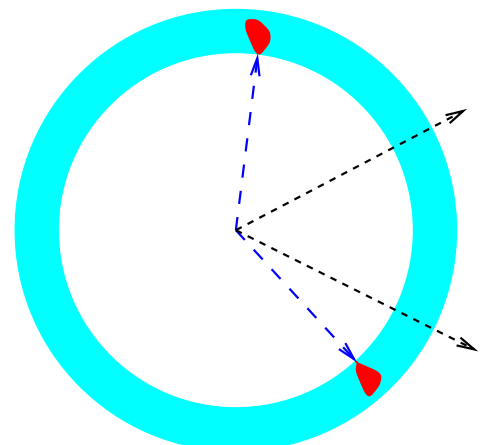
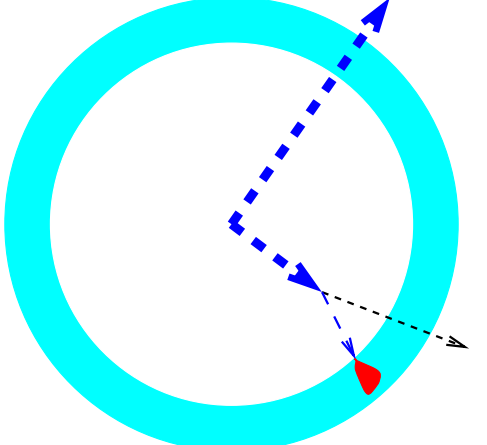
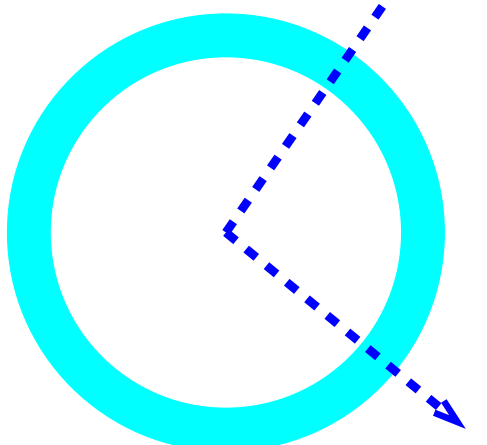
$$c\tau \approx (1.3\text{m}) \left(\frac{100\text{GeV}/c^2}{m(\text{NLSP})} \right)^5 \left(\frac{\sqrt{F}}{1000\text{TeV}/c^2} \right)^4$$

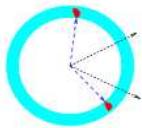


($\text{sgn}\mu, M/\Lambda$): (+, 2) solid, (+, 10^4) dashed, (-, 2) dot-dashed and (-, 10^4) dotted; lines are drawn for equal masses of stau and neutralino.

(The plot is taken from hep-ph/9609444)

Why GMSB?

NLSP	$CT \simeq 0$	$CT \simeq \text{det. size}$	$CT \gg \text{det. size}$
$\tilde{\tau}_1 \rightarrow \tilde{G}\tau$			
$\tilde{N}_1 \rightarrow \tilde{G}\gamma$			



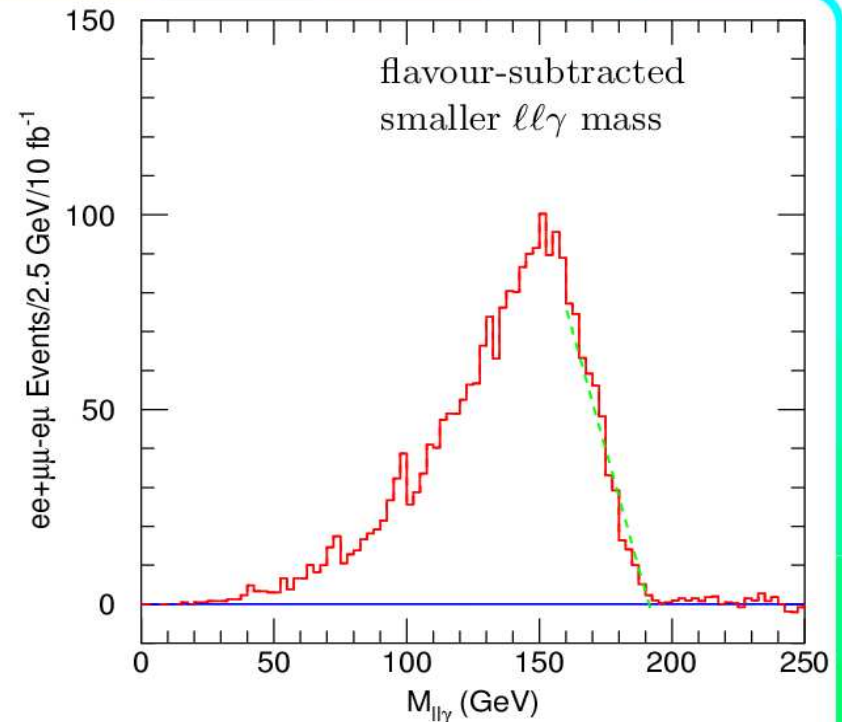
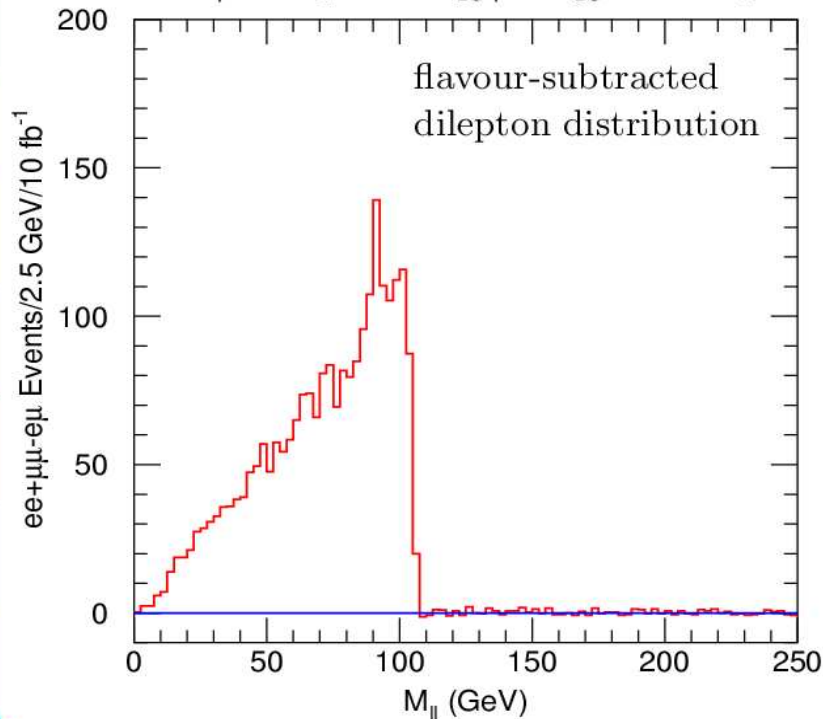
ATLAS point G1a

NLSP: \tilde{N}_1 , 119 GeV; short-lived; 7.6 pb

At this point: $\tilde{N}_1 \rightarrow \tilde{G} \gamma$ and
 $\tilde{N}_2 \rightarrow \tilde{\ell}^\pm \tilde{\ell}^\mp \rightarrow \tilde{N}_1 \ell^+ \ell^- \rightarrow \tilde{G} \ell^+ \ell^- \gamma$

In the flavour-subtracted distribution

$e^+e^- + \mu^+\mu^- - e^\pm\mu^\mp$ one can very sharp endpoint
 $\sqrt{M_{\tilde{N}_2}^2 - M_{\tilde{\ell}_R}^2} \sqrt{M_{\tilde{\ell}_R}^2 - M_{\tilde{N}_1}^2} / M_{\tilde{\ell}_R}$.



Taking into account the whole 4-body decay chain one can measure 3 more endpoints

$$M_{\ell\ell\gamma} = \sqrt{M_{\tilde{N}_2}^2 - M_{\tilde{N}_1}^2} \quad (\text{linear} - \text{figure})$$

$$M_{\ell\gamma}^{(1)} = \sqrt{M_{\tilde{\ell}_R}^2 - M_{\tilde{N}_1}^2} \quad (\text{edge})$$

$$M_{\ell\gamma}^{(2)} = \sqrt{M_{\tilde{N}_2}^2 - M_{\tilde{\ell}_R}^2} \quad (\text{linear})$$

Which is sufficient to determine $\tilde{N}_2, \tilde{\ell}_R$ and \tilde{N}_1 masses

Fitting GMSB parameters by ATLAS

Point **G1A** 7.6 pb, short-lived \tilde{N}_1

quantity	low-lumi	ultimate	
$M_h(GeV)$	109.47 ± 3.0	109.27 ± 0.2	
$M_{\ell\ell}^{max}(GeV)$	105.1 ± 0.1	same	
$M_{\ell\ell\gamma}^{max}(GeV)$	189.7 ± 0.3	189.7 ± 0.3	
$M_{\ell\gamma}^{(1)}(GeV)$	112.7 ± 0.15	112.7 ± 0.1	
$M_{\ell\gamma}^{(2)}(GeV)$	152.6 ± 0.3	152.6 ± 0.2	
parameter	low-lumi	high-lumi	ultimate
$\Lambda(TeV)$	90 ± 1.7	90 ± 0.89	same
$M(TeV)$	500 ± 170	500 ± 110	same
n	1.0 ± 0.014	1.0 ± 0.011	same
$\tan\beta$	5.0 ± 1.3	5.0 ± 0.4	5.0 ± 0.14

Point **G2A** 23 pb, short-lived $\tilde{\tau}_1$

quantity	low-lumi	ultimate	
$M_h(GeV)$	106.6 ± 3.0	106.6 ± 0.2	
$M_{\ell\ell}^{max(1)}(GeV)$	52.21 ± 0.05	same	
$M_{\ell\ell}^{max(2)}(GeV)$	175.94 ± 0.18	same	
$M_{\ell\ell q}^{max}(GeV)$	640 ± 7	same	
$M_{\ell q}^{max}/M_{\ell\ell q}^{max}$	0.450 ± 0.004	same	
parameter	low-lumi	high-lumi	ultimate
$\Lambda(TeV)$	30 ± 0.54	same	same
$M(TeV)$	250 ± 60	same	same
n	3.0 ± 0.05	same	same
$\tan\beta$	5.0 ± 1.0	5.0 ± 1.0	5.0 ± 0.06

Point **G1B** 7.6 pb, long-lived \tilde{N}_1

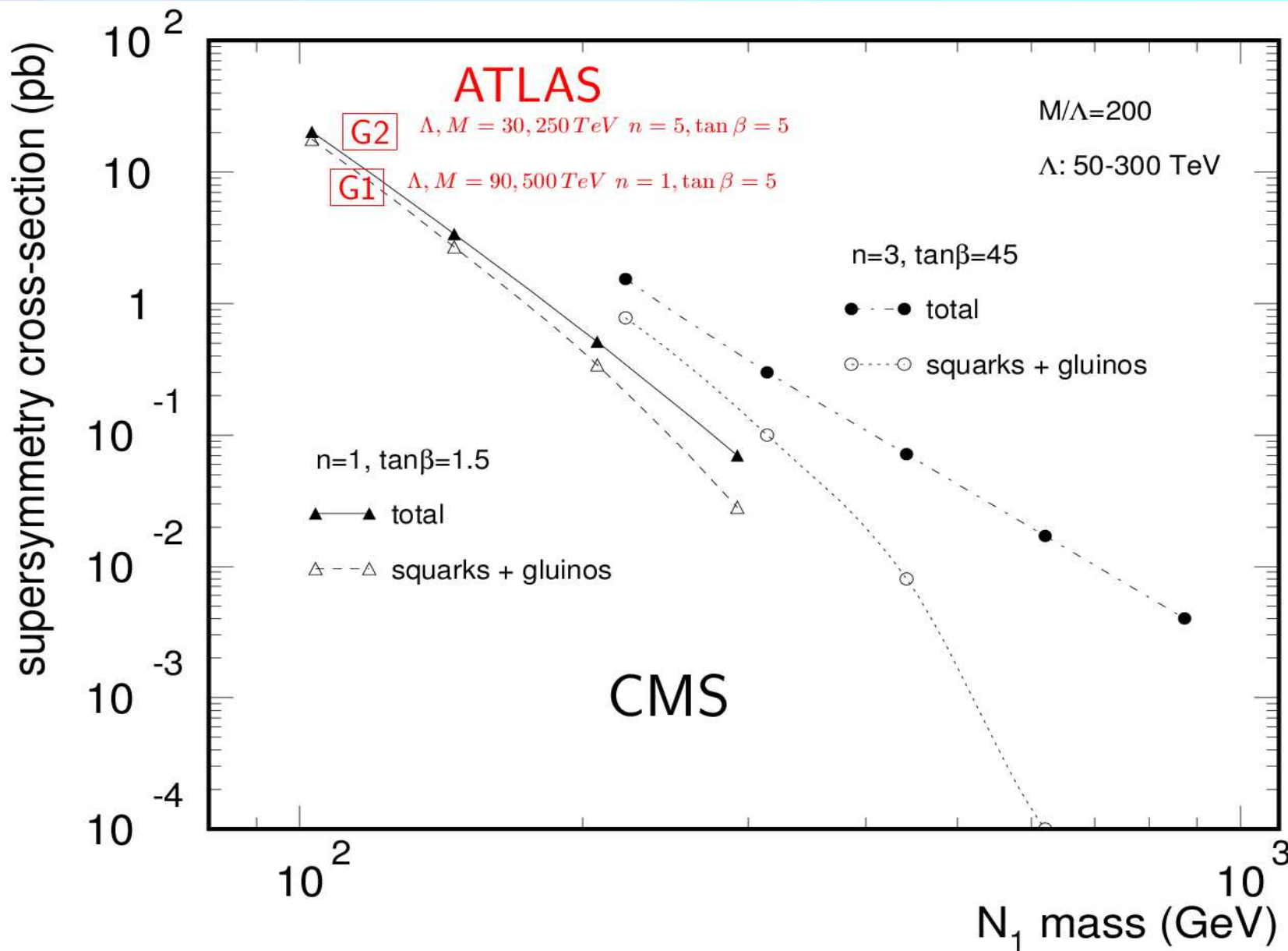
quantity	low-lumi	ultimate	
$M(\tilde{g}) - M(\tilde{N}_2)(GeV)$	189.7 ± 0.3	189.7 ± 0.3	
$M(\tilde{q}_L)(GeV)$	112.7 ± 0.15	112.7 ± 0.1	
$(M_h \& M_{\ell\ell}^{max}$ as in point G1A)			
For 30/fb the fit gives (for $\mu > 0$):			
$n \times \Lambda = (90 \pm 0.88) TeV, \Lambda = (90 \pm 11.5) TeV$			
$M < 7 \times 10^5 TeV, \tan\beta = 5.0_{-1.8}^{+2.7}$			

Point **G2B** 23 pb, long-lived $\tilde{\tau}_1$

Quasi-stable sleptons:

many precise mass measurements

parameter	low-lumi	high-lumi	ultimate
$\Lambda(TeV)$	30 ± 0.25	same	same
$M(TeV)$	250 ± 32	same	same
n	3.0 ± 0.02	same	same
$\tan\beta$	5.0 ± 0.3	5.0 ± 0.3	5.0 ± 0.03



SUSY GMSB event: cascade decays to

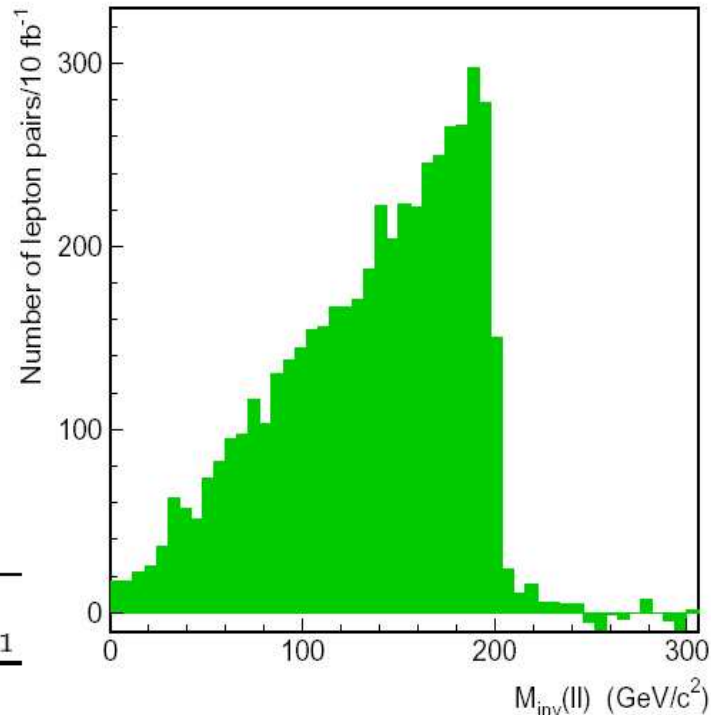
Next to the Lightest Supersymmetry Particle

neutralino \tilde{N}_1

$$\begin{aligned} \tilde{N}_2 &\rightarrow \tilde{\ell}_R^\pm + \ell^\mp \\ &\hookrightarrow \tilde{N}_1 + \ell^\pm \end{aligned}$$

$$M_{inv}(l\bar{l}) = \frac{\sqrt{m_{\tilde{N}_2}^2 - m_{\tilde{\ell}_R}^2} \sqrt{m_{\tilde{\ell}_R}^2 - m_{\tilde{N}_1}^2}}{m_{\tilde{\ell}_R}}$$

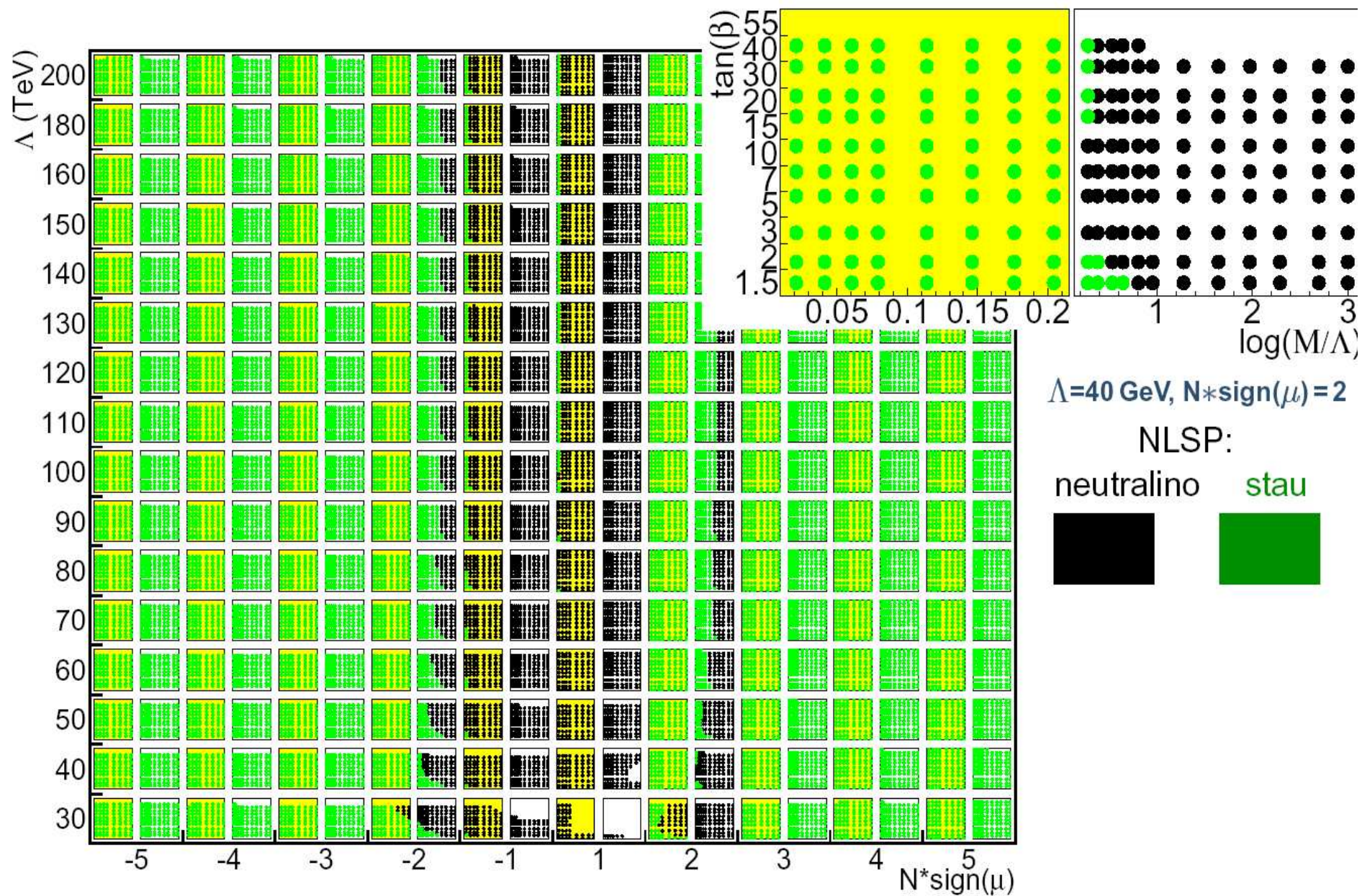
In $M_{inv}(l\bar{l})$ of $\mu^+\mu^- + e^+e^- - \mu^+e^- - e^+\mu^-$ sleptons: $\tilde{\tau}_1$ or $\tilde{e}_R \tilde{\mu}_R$



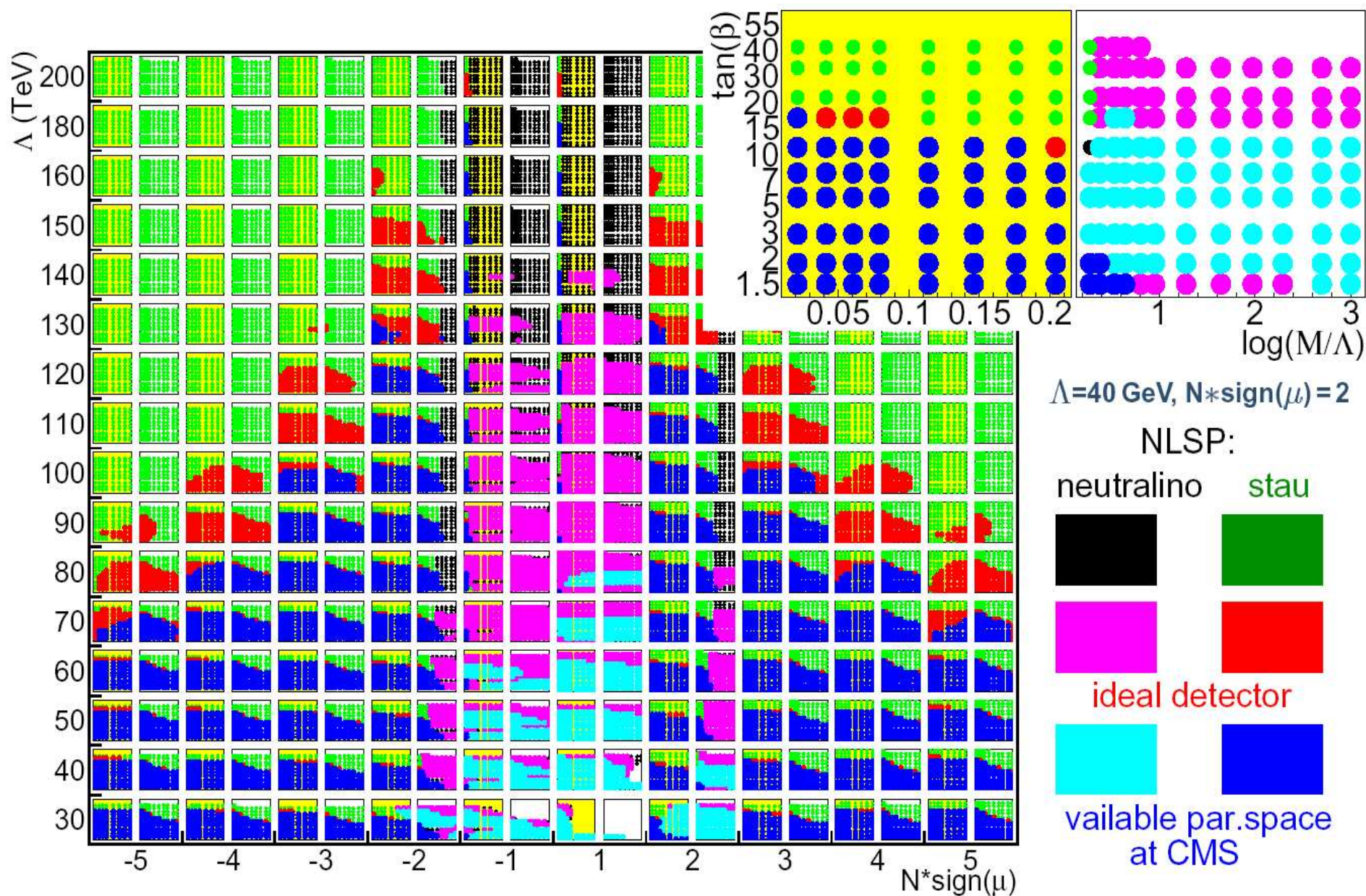
$$\begin{aligned} \tilde{N}_1 &\rightarrow \tilde{\ell}_R^\pm + \ell^\mp \\ &\hookrightarrow \tilde{G} + \ell^\pm \end{aligned}$$

$$M_{inv}(l\bar{l}) = \sqrt{m_{\tilde{N}_1}^2 - m_{\tilde{\ell}_R}^2}$$

Scan over parameters space



Scan over parameters space



Models: Case Studies

Goals:

- For more complex signals (cascades, production of multiple particles), check if consistent with various models
- Large number of models, phenomenology depends on free parameters → case studies

Projects:

- SUSY
 - mSUGRA (Alexander Belyaev, Tadas Krupovnickas)
 - MSSM (Howie Baer, Georg Weiglein, Csaba Balazs)
 - NMSSM (Sabine Kraml, Cyril Hugonie)
 - Stops (Sabine Kraml, Joe Lykken, Shoji Asai)
 - Disentangling Models (Dirk Zerwas et al. (SFITTER), Bob Kehoe, Gordon Kane)
 - Tools (Michael Spira (PROSPINO), Peter Skands (SLHA))
- Universal Extra Dimensions (Kyoungchul Kong, Konstantin Matchev, Satya Nandi)
- Signal Generators and tools for Extra Dimensions (Albert de Roeck et al.)
- Little Higgs Models with T-Parity (Jay Hubisz)
- Technicolor (Ken Lane)
- Higgsless Models (Andreas Birkedal)

Please add your name to the list!

SUSY Fits in Tevatron and LHC Data (SFITTER)

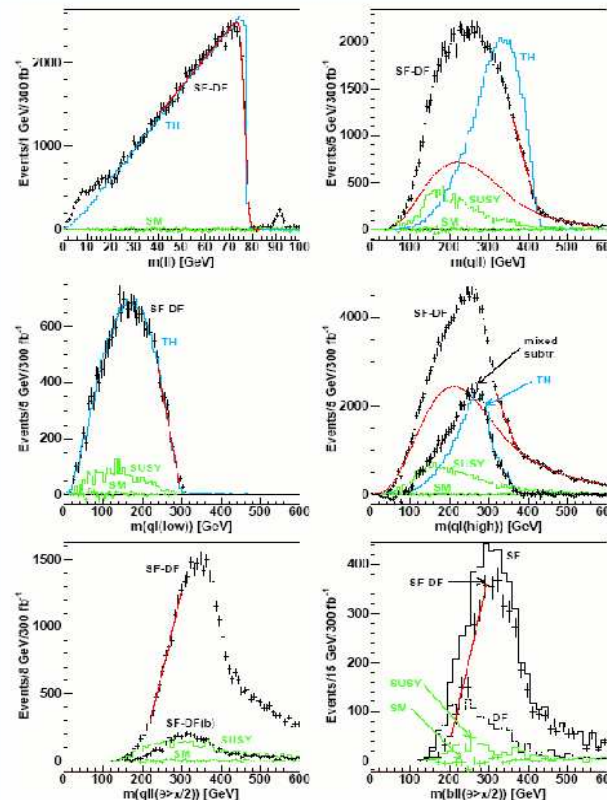
Examples of measurements at LHC

Gjelsten et al: ATLAS-PHYS-2004-007/29

$$\begin{aligned}
 (m_{ll}^2)^{\text{edge}} &= \frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{t}_R}^2)(m_{\tilde{t}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{t}_R}^2} \\
 (m_{qll}^2)^{\text{edge}} &= \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{\chi}_2^0}^2} \\
 (m_{ql}^2)^{\text{edge}_{\text{min}}} &= \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{t}_R}^2)}{m_{\tilde{\chi}_2^0}^2} \\
 (m_{ql}^2)^{\text{edge}_{\text{max}}} &= \frac{(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{t}_R}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{t}_R}^2} \\
 (m_{ql}^2)^{\text{thres}} &= \frac{[(m_{\tilde{q}_L}^2 + m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{t}_R}^2)(m_{\tilde{t}_R}^2 - m_{\tilde{\chi}_1^0}^2) - (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)\sqrt{(m_{\tilde{\chi}_2^0}^2 + m_{\tilde{t}_R}^2)^2(m_{\tilde{t}_R}^2 + m_{\tilde{\chi}_1^0}^2)^2 - 16m_{\tilde{\chi}_2^0}^2 m_{\tilde{t}_R}^4 m_{\tilde{\chi}_1^0}^2} + 2m_{\tilde{t}_R}^2(m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)]}{4m_{\tilde{t}_R}^2 m_{\tilde{\chi}_2^0}^2}
 \end{aligned}$$

Edge	Nominal Value	Fit Value	Syst. Error Energy Scale	Statistical Error
$m(ll)^{\text{edge}}$	77.077	77.024	0.08	0.05
$m(qll)^{\text{edge}}$	431.1	431.2	4.3	2.4
$m(ql)^{\text{edge}_{\text{min}}}$	302.1	300.8	3.0	1.5
$m(ql)^{\text{edge}_{\text{max}}}$	380.3	379.4	3.8	1.8
$m(ql)^{\text{thres}}$	203.0	204.6	2.0	2.8
$m(bl)^{\text{thres}}$	183.1	181.1	1.8	6.3

plus other mass differences and edges...



From edges to masses:
System overconstrained

What is the problem?

CPX scenario:

→ emphasize “possible” large effects:

[*M. Carena, J. Ellis, A. Pilaftsis, C. Wagner '01*]

$$M_{\text{SUSY}} = 500 \text{ GeV}, \quad |A_t| = 1 \text{ TeV},$$

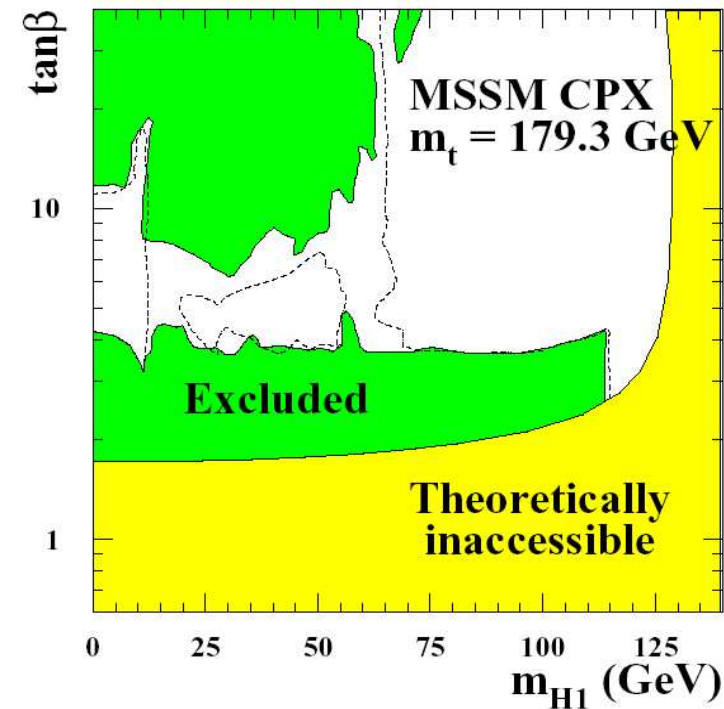
$$A_b = A_\tau = A_t,$$

$$M_2 = 500 \text{ GeV}, \quad |m_{\tilde{g}}| = 1 \text{ TeV},$$

$$\mu = 2 \text{ TeV}$$

$$\Phi = \Phi_{A_{t,b,\tau}} = \Phi_{m_{\tilde{g}}}$$

$$M_{H^\pm}, \tan\beta \text{ varied}$$



LEP search left uncovered holes with low $m_{h_1} \Rightarrow$ difficult for LHC

Q: Can the Tevatron cover these holes?

A'': However, another thing . . .

LEP analysis:

Two codes: **FeynHiggs2.0** and **CPH**
 conservative approach (very good!):

CPX point not excluded



(point not excluded by FeynHiggs) or (point not excluded by CPH)

CPX holes "rely" heavily on large $BR(h_2 \rightarrow h_1 h_1)$

CPX holes: **CPH** has larger $BR(h_2 \rightarrow h_1 h_1)$ than **FeynHiggs**

⇒ reasons for differences under investigation

⇒ possibly higher-order effects that are not under control

(⇒ LEP analysis is currently the best strategy!)

A'': However, yet another thing ...

LEP analysis:

holes not excluded at the 95% C.L.

holes are excluded at the $\sim 75\%$ C.L.

\Rightarrow "combined" LEP/TeVatron analysis ??