

Metody eksperymentalne w fizyce wysokich energii

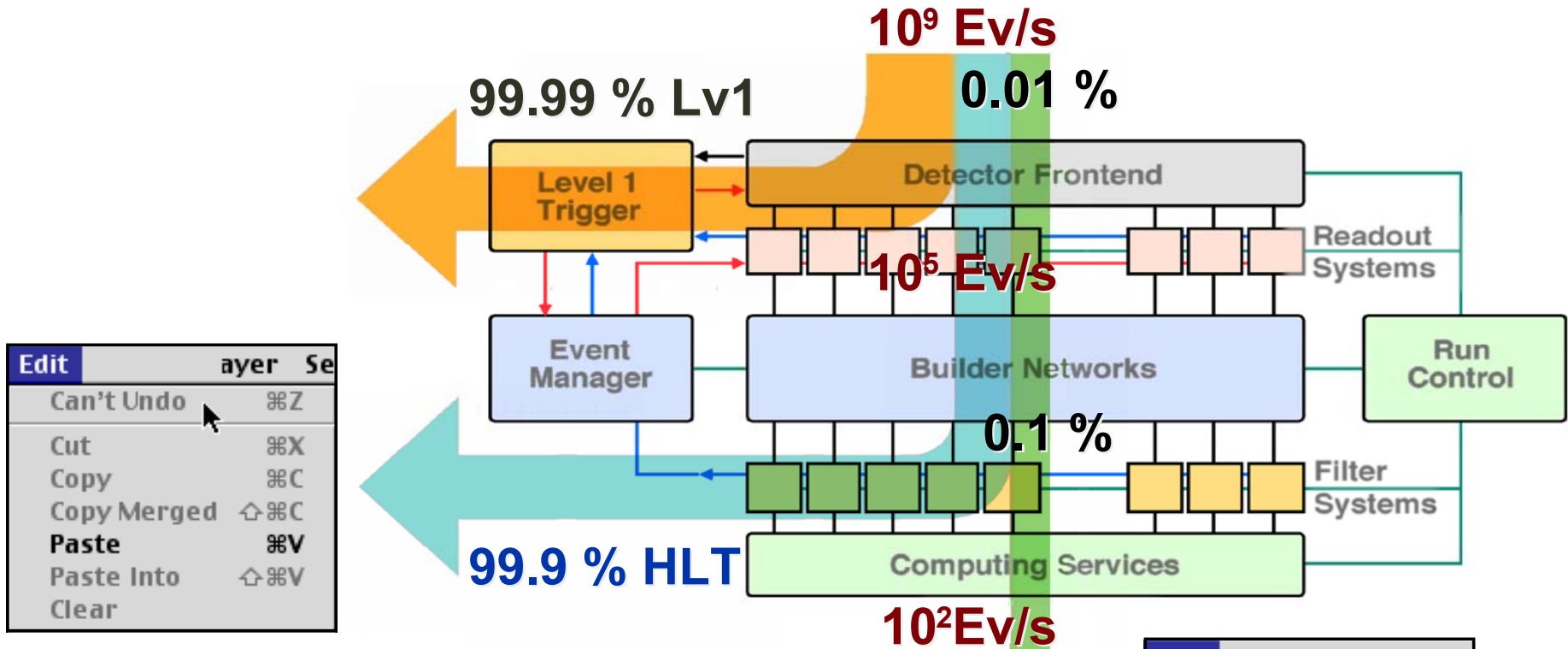
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Zakład Cząstek i Oddziaływań Fundamentalnych IFD

Wykład XII

- Rekonstrukcja przypadków

A parting thought



With respect to offline analysis:

Same hardware (Filter Subfarms)

Same software

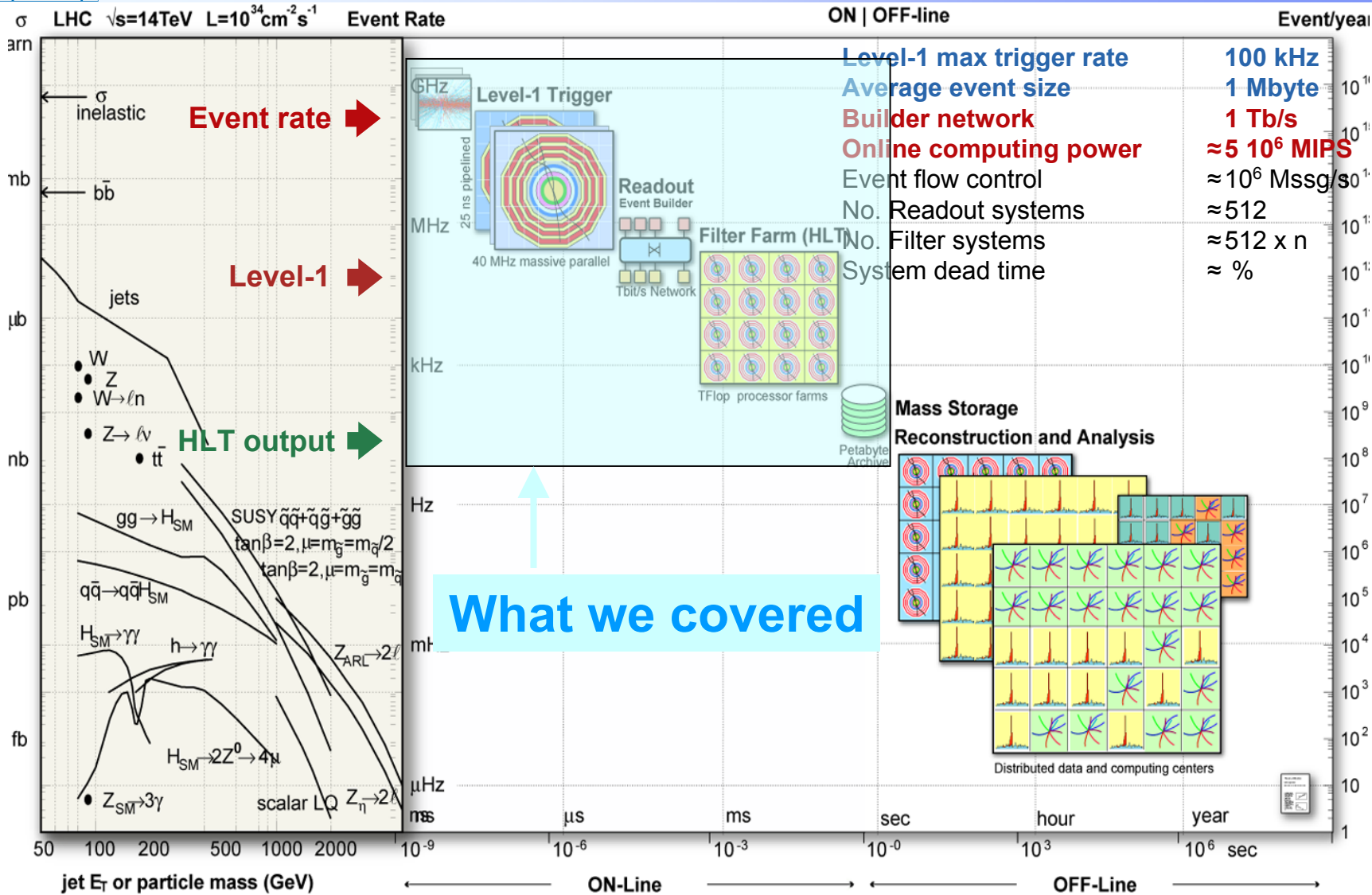
But different situations

Edit	ayer	Se
Can't Undo	⌘Z	
Cut	⌘X	
Copy	⌘C	
Copy Merged	⇧⌘C	
Paste	⌘V	
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Copy Merged	⇧⌘C	
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Paste Into	⇧⌘V	
Clear		



Online Physics Selection: summary



Dane odczytywane z detektora:

- "zapalone" cele kalorymetru
 - ID celi + sygnał**
 - + ew. informacja o profilu podłużnym lub poprzecznym w przypadku wielokrotnych odczytów
- "hity" w detektorach śladowych
 - ID/pozycja** drutu/paska/piksla
 - + ew. sygnał ($\sim dE/dx$)
 - + ew. czas (komory dryfowe, TPC)

Co potrzebujemy do analizy fizycznej:

- zidentyfikowane **fotony**: energia, kierunek emisji (\Rightarrow **wektor pędu**)
- zidentyfikowane **leptony** (elektrony/miony/taony): energia, pęd, znak
- zrekonstruowane **jety**: energia, pęd, masa niezmi., **zapach** (!)
- zrekonstruowany brakujący pęd poprzeczny, energia, masa...

Dodatkowe dane wejściowe

Wyznaczane na podstawie zebranych danych (w tym w odpowiedzi na detykowane triggery), parametrów pracy detektora itp.:

parametry kalibracji

parametry pozycjonowania

parametry efektywności

Niezbędne do rekonstrukcji danych.

Rzadko/powoli się zmieniają - w skalach dni, tygodni.

W najgorszym przypadku co napełnienie (fill).

Czasami raz na rok (od shutdownu do shutdownu).

Zazwyczaj przechowywane w dedykowanych bazach danych.

Kolejne kroki rekonstrukcji

Krok I: "normalizacja" na poziomie poszczególnych detektorów

kalorymetry

ID celi => pozycja w detektorze, z uwzględnieniem pozycjonowania
sygnał => depozyt energii w GeV, z uwzględnieniem kalibracji, tła

detektory śladowe

ID elementu, ew. czas dryfu

=> pozycja w detektorze, z uwzgl. pozycjonowania

Krok II: rekonstrukcja podstawowych "elementów"

wciąż na poziomie poszczególnych detektorów

kalorymetry

"klastry"/"wyspy" energii (pojedyncze cząstki ?)

detektory śladowe

"klastry"/punkty, ew. elementy/fragmenty toru (punkt + kierunek)

Krok III: **rekonstrukcja obiektów** na na poziomie globalnym

- rekonstrukcja torów
- rekonstrukcja wierzchołka pierwotnego i w. wtórnych
- rekonstrukcja jetów

Krok IV: **identyfikacja obiektów** (na poziomie globalnym
łączenie tor-klaster (i ew. korekta energii: PFA)

- identyfikacja leptonów
- znaczenie zapachów jetów (ciężkie kwarki)

Rekonstrukcja torów

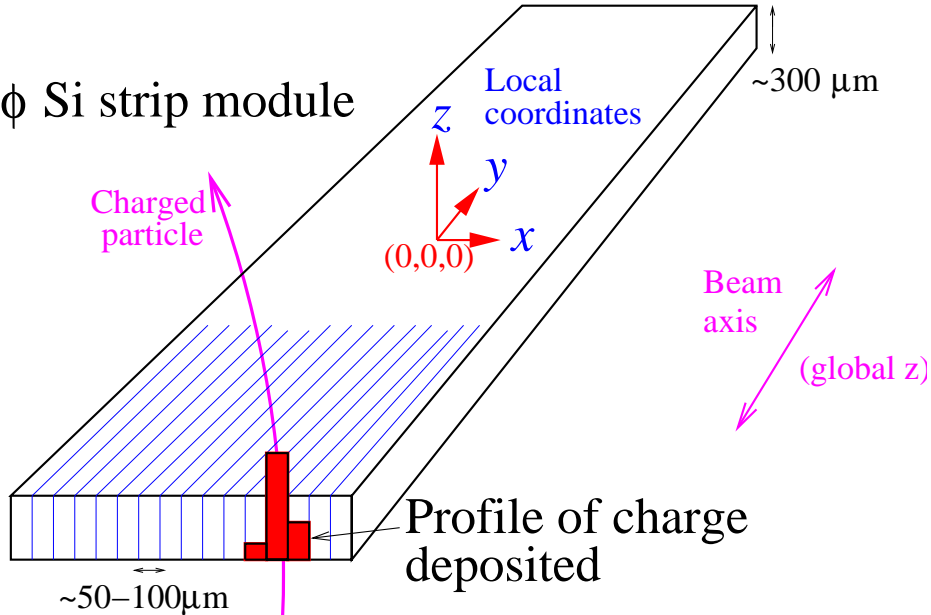
Basic Idea

When we talk about “tracking,” we want to do the following:

- ▶ Measure the true path of the charged particle, which let's us know...
- ▶ The momentum (3-momentum) if we know the magnetic field
- ▶ The sign of the charge of the particle
- ▶ With other constraints or assumptions, the “origin” in space of the particle
- ▶ Without some other detector though, we can't measure the mass independently just with a tracker

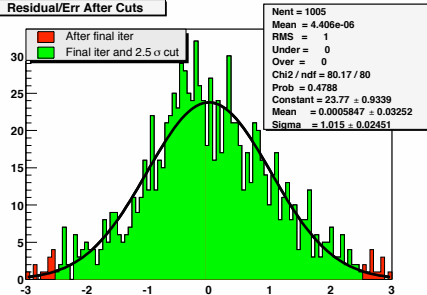
2D Point ($r\phi$ or rz): $\bar{x} = (\sum i * q_i) / \sum q_i$

$r\phi$ Si strip module

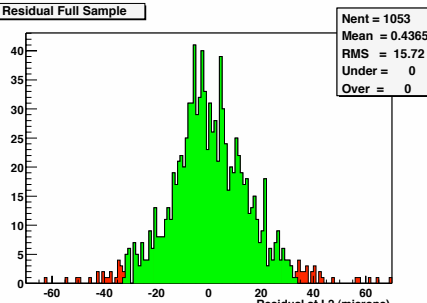


Typical Resolution of 50 μm Strips (CDF)

Residual/Err After Cuts



Residual Full Sample



Cluster Width	Resolution
1	12 μm
2	9 μm
3	14 μm
4+	22 μm

- ▶ rz strips are either shallow stereo (2°), or 90° stereo but larger pitch
- ▶ Intrinsic resolution of larger pitch stereo usually factor of 2+ worse than $r\phi$

Resolution of Si Strip Detector

- ▶ The resolution of a Si hit depends on the number of strips in the cluster (2-strip most precise)
- ▶ Charge drifts with Lorentz force: $q(\vec{E} + \vec{v} \times \vec{B})$
- ▶ Thin material: Landau distribution of charge [2]
- ▶ Track impact angle & position makes a difference [3]
- ▶ Delta rays (hard knock e^-) can bias charge distribution
- ▶ Dead channels, noise, V_{dep} , temperature all could affect this too
- ▶ Radiation damage changes
- ▶ Pileup from previous event
- ▶ Multiple particles passing through same strips

3D Space Point from Pixel Detector

CMS barrel
pixel module

Local
coordinates
 z
 y
 x
 $(0,0,0)$

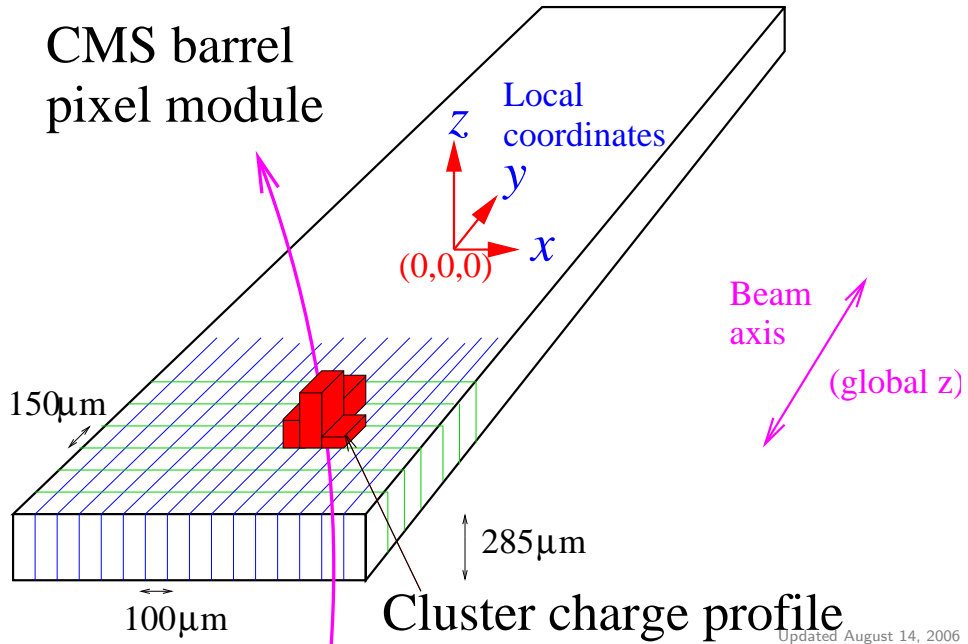
Beam
axis
(global z)

$150\mu\text{m}$

$285\mu\text{m}$

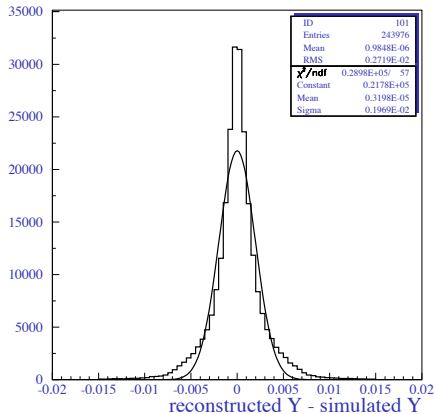
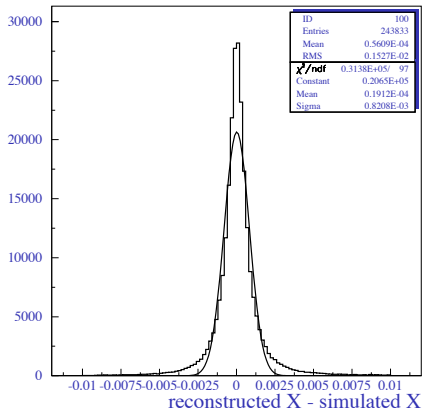
$100\mu\text{m}$

Cluster charge profile



Typical Pixel Resolution (CMS):

8 – 20 μm (with 100/150 μm pitch !)



Detektory pixlowe

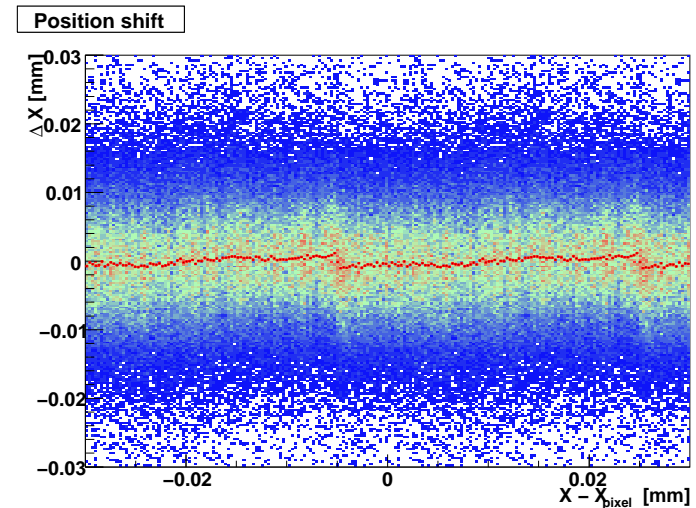
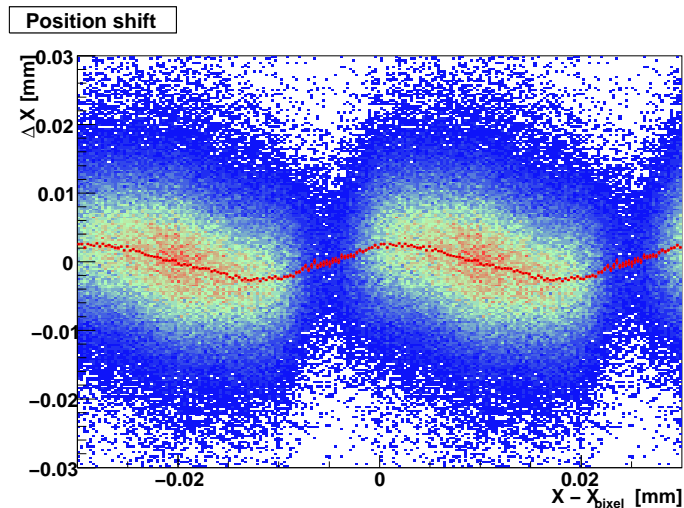
MAPS

Wyniki testów detektora MAPS, rozmiar pixla $30 \times 30 \mu m$

Błąd pozycji vs pozycja na pikslu.

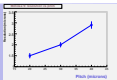
Pozycja wyznaczana metodą **środka ciężkości** (CoG)

Pozycja CoG po zastosowaniu poprawki (tzw. funkcja Eta)



Pojedyncza cząstka “zapala” średnio 4 pixle \Rightarrow systematyczny błąd pozycji zależy od punktu przejścia cząstki

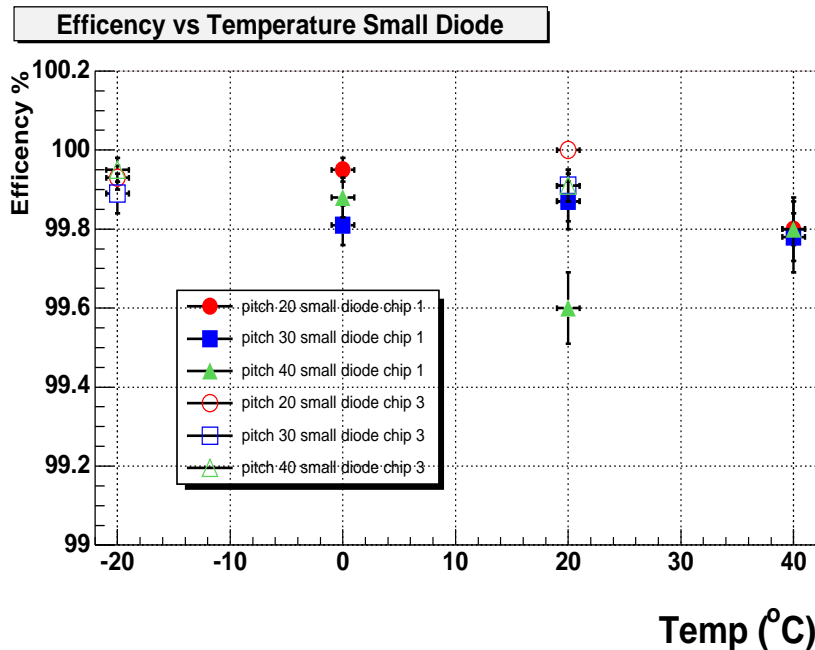
wyznaczana z danych poprawka prawie całkowicie eliminuje efekty systematyczne



Detection Efficiency & Spatial Resolution

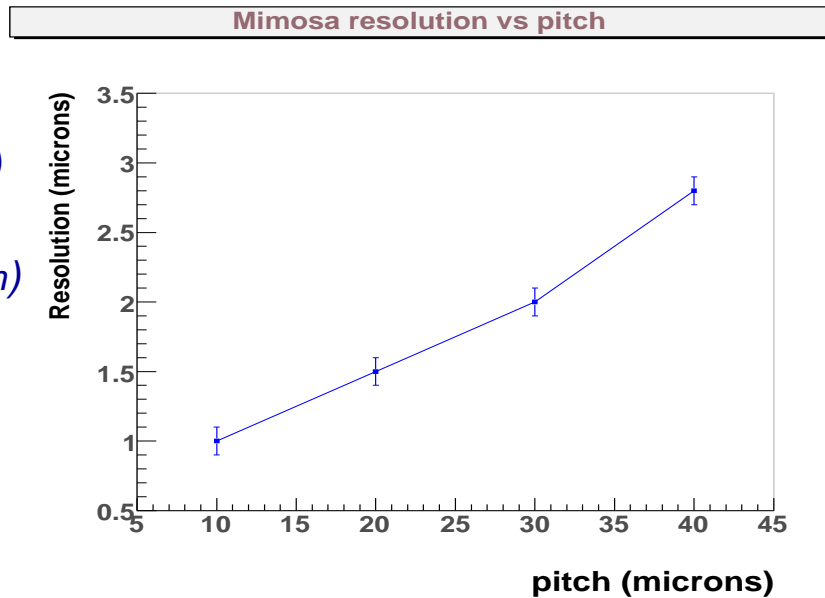
● Detection efficiency:

- * *Ex: MIMOSA-9 data (20, 30 & 40 μm pitch)*
- * $\epsilon_{det} \gtrsim 99.5\text{--}99.9\%$ repeatedly observed at room temperature (fake rate $\sim 10^{-5}$)
- * $T_{oper.} \gtrsim 40^\circ\text{C}$



● Single point resolution versus pixel pitch:

- * *clusters reconstructed with eta-function, exploiting charge sharing between pixels (12-bit ADC)*
- * $\sigma_{sp} \sim 1 \mu\text{m}$ (10 μm pitch) $\rightarrow \lesssim 3 \mu\text{m}$ (40 μm pitch)
- * *4-bit ADC simul. $\Rightarrow \sigma_{sp} \lesssim 2 \mu\text{m}$ (20 μm pitch)*
- * *measured binary output resolution (MIMOSA-16, -22):*
 $\sigma_{sp} \gtrsim 3.5 \text{ \& } 4.5 \mu\text{m}$ (18.4 & 25 μm pitch)



Pattern Recognition & Track Fitting

Typically, pattern recognition algorithms are either “inside-out” or “outside-in.”

- ▶ You have to start with some idea of the path of the particle to bootstrap your algorithm: a track seed
- ▶ Then you take this candidate, this seed, can try to find compatible hits in other layers
- ▶ Continue this process until you’ve met some criteria for what a “good track” should have
- ▶ Once you’ve got your hits for your track, try to do a good job of fitting your pseudohelix
- ▶ Pseudohelix because there is energy loss and multiple scattering
- ▶ There are also spurious hits from detector noise and low momentum, unreconstructable tracks. These will mess up your “true” helix

Helix Parameters

We can decompose the momentum of a track in spherical coordinates [1]:

$$p_x = p \cos \phi \sin \theta$$

$$p_y = p \sin \phi \sin \theta$$

$$p_z = p \cos \theta$$

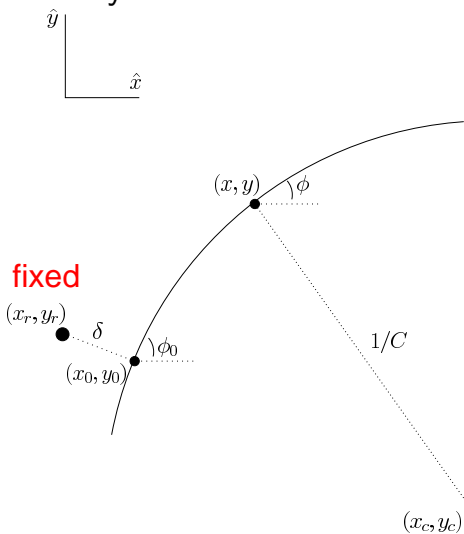
Different experiments choose different ranges for the angles, it's important that you figure out what they are using:

$$\phi \in [-\pi, \pi] \quad \theta \in [0, \pi]$$

There must also be some “reference point” in space to uniquely define our helix: (x_r, y_r, z_r)

Helix Parameters in $x - y$ Plane

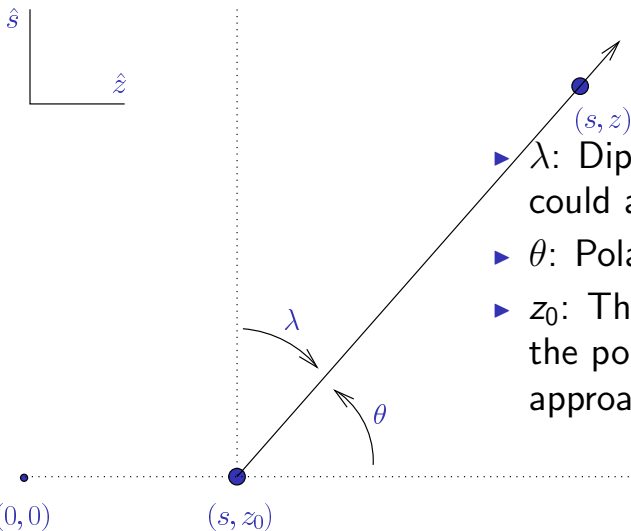
This parameterization is more closely related to things we actually measure with our trackers



- ▶ C : Curvature of the track. Signed with charge.
- ▶ ϕ_0 : Azimuthal angle of the momentum at the point of closest approach
- ▶ δ : Distance of closest approach. (Also signed, but differently.)

3 parameters in x-y

$$z = z_0 + s \tan \lambda$$



- ▶ λ : Dip angle of track, or could also use
- ▶ θ : Polar angle of track
- ▶ z_0 : The z of the track at the point of closest approach in $x - y$

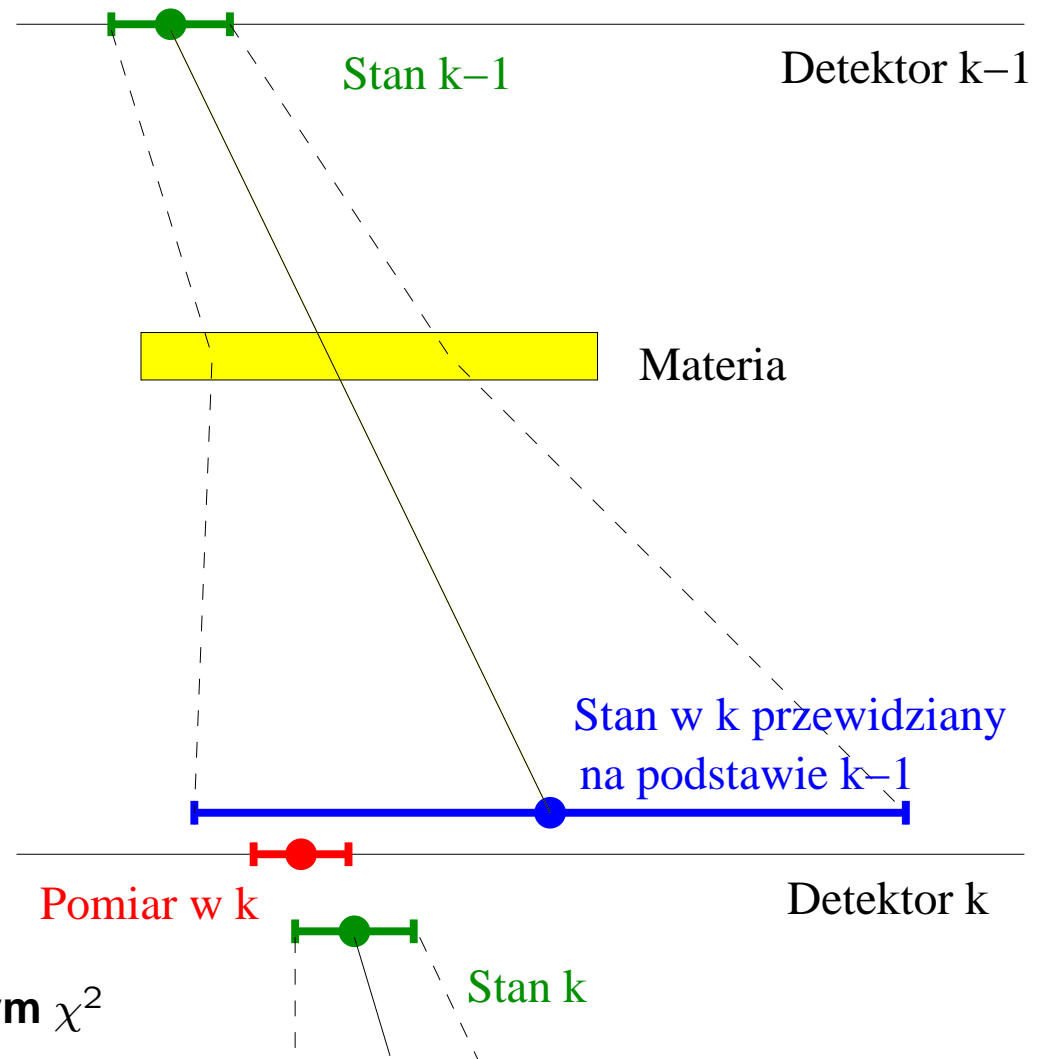
+ 2 parameters in r-z

Kalman Filter

- ▶ The Kalman Filter [4] is an iterative procedure
- ▶ You start with a seed track. For example, a pair of hits in the pixel detector that line up within 5σ of the primary interaction region
- ▶ Then you add points on successive layers, taking into account projected error from current hypothesis track and multiple scattering
- ▶ Finally, if this track passes some requirements for a minimum number of hits, you refit it (smooth) again with the filter using a better starting point

Filtr Kalmana - sposób na dopasowywanie torów

- Algorytm oblicza parametry toru używając punktów “pasujących” do toru:
- Punkty są dodawane zaczynając od najbardziej zewnętrznych
- **Stan początkowy** jest brany z zewnętrznych detektorów śladowych
- **Stan obliczony dla detektora k-1** jest **ekstrapolowany** do detektora k z uwzględnieniem wpływu materii.
- **Punkty w znalezione w detektorze k** są porównywane z **przewidywaniem** i są dodawane jeżeli nie wprowadzają dużego wzrostu χ^2 przy dopasowaniu toru - obliczany jest **stan k**
- Jeżeli więcej niż jeden punkt “pasuje” brane są pod uwagę wszystkie możliwości i ostatecznie wybierany jest tor z największą liczbą punktów i najmniejszym χ^2



Tracking Resolution

The resolution of the various helix parameters depends on a number of things

- ▶ The number of hits used from various subdetectors (hopefully, more is better)
- ▶ The momentum of the particle. Higher momentum particles deflect less from multiple scattering
- ▶ The polar angle of the track (η)
- ▶ Quality of alignment
- ▶ Presence of other tracks!
- ▶ Detector noise

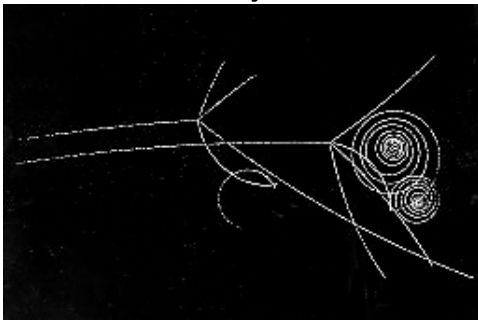
High Multiplicity Strategies

- ▶ If your inner layers are being swamped, use an outside-in algorithm instead of an inside-out algorithm
- ▶ Try to boost efficiency for higher p_T tracks by reconstructing them first, and then remove these hits from consideration
- ▶ Upgrade your detector

Rekonstrukcja wierzchołków

Vertexing: Basic Idea

The basic idea of vertexing is to figure out where the particles came from. We can associate tracks to particle decays and interactions this way.



Vertexing Strategy

- ▶ A vertex is a point where more than one particle comes from
- ▶ If there is more than one track coming from the same place, then the helices should cross each other, right?
- ▶ Look for places where helices cross
- ▶ Caveat: The track parameters $\vec{p}_i = (C, \phi, d_0, \tan \lambda, z_0)$ are different in different parts of the helix

Vertexing Algorithm

- ▶ If helixes cross, that means the track parameters must be the same at some point
- ▶ In 2D, this is as simple as looking for crossing circles
- ▶ But you can take full advantage of full track parameterization and covariance matrix to look for vertexes
- ▶ Most methods [5] are built on some kind of χ^2 of track parameters

$$\chi^2 = \sum_i^{N_{\text{trk}}} (\vec{x} - \vec{g}(\vec{p}_i))^T \mathbf{J}^T \mathbf{M}_i^{-1} \mathbf{J} (\vec{x} - \vec{g}(\vec{p}_i))$$

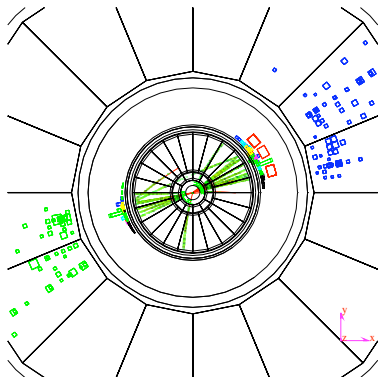
Vertex Algorithm Continued

- ▶ Again, we have a problem of pattern recognition
- ▶ The χ^2 can be defined for any collection of tracks
- ▶ Vertex algorithms are usually iterative however, pruning or down-weighting tracks which make large contributions to the χ^2 (outliers)
- ▶ Clumps, or clusters of tracks are what is searched for in this way

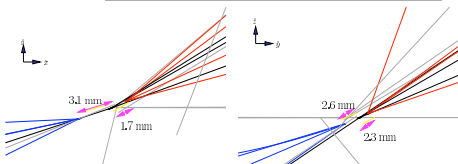
Rekonstrukcja wierzchołków wtórnych

B-Tagging: Basic Idea

Run # 441525 Event # 1504 Total Energy : 110.38 GeV



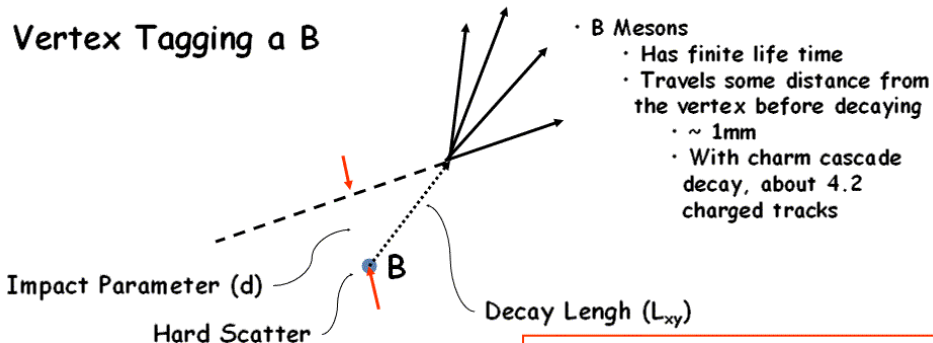
Transverse imbalance : .0290		Longitudinal imbalance : -.2966	
Thrust : .9175	Major : .2847	Minor : .0800	
Event DAQ Time : 800000		1	



- ▶ B hadrons have lifetimes and decay lengths distinct from other species
- ▶ Decay length is measurable in a given event by finding a vertex (“secondary”) and taking the distance to the “primary” vertex
- ▶ Can look for B hadrons by finding these vertexes which are compatible with known properties of B hadrons

Example B-Tag (D0)

Vertex Tagging a B

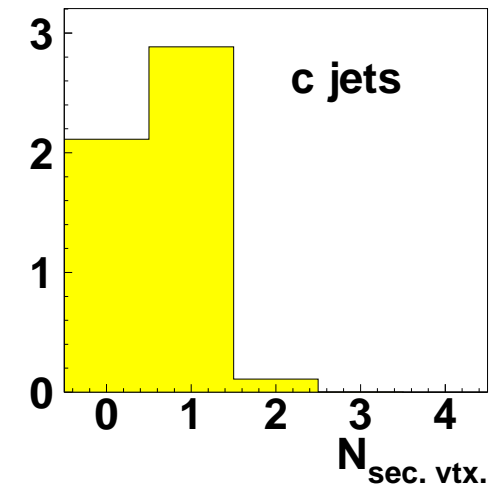
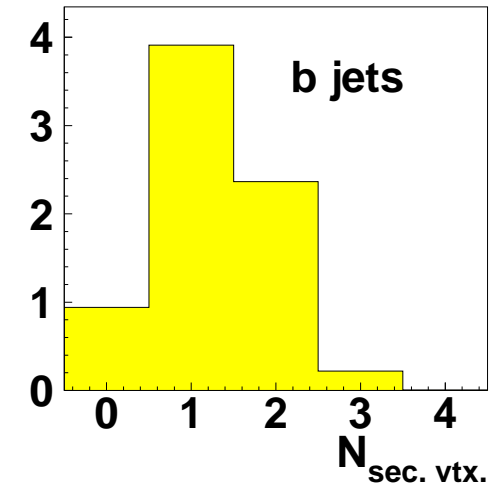
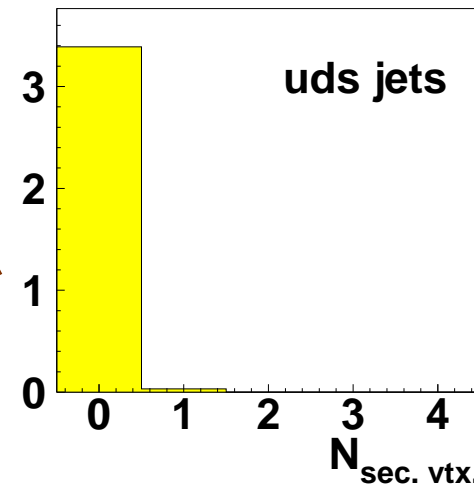
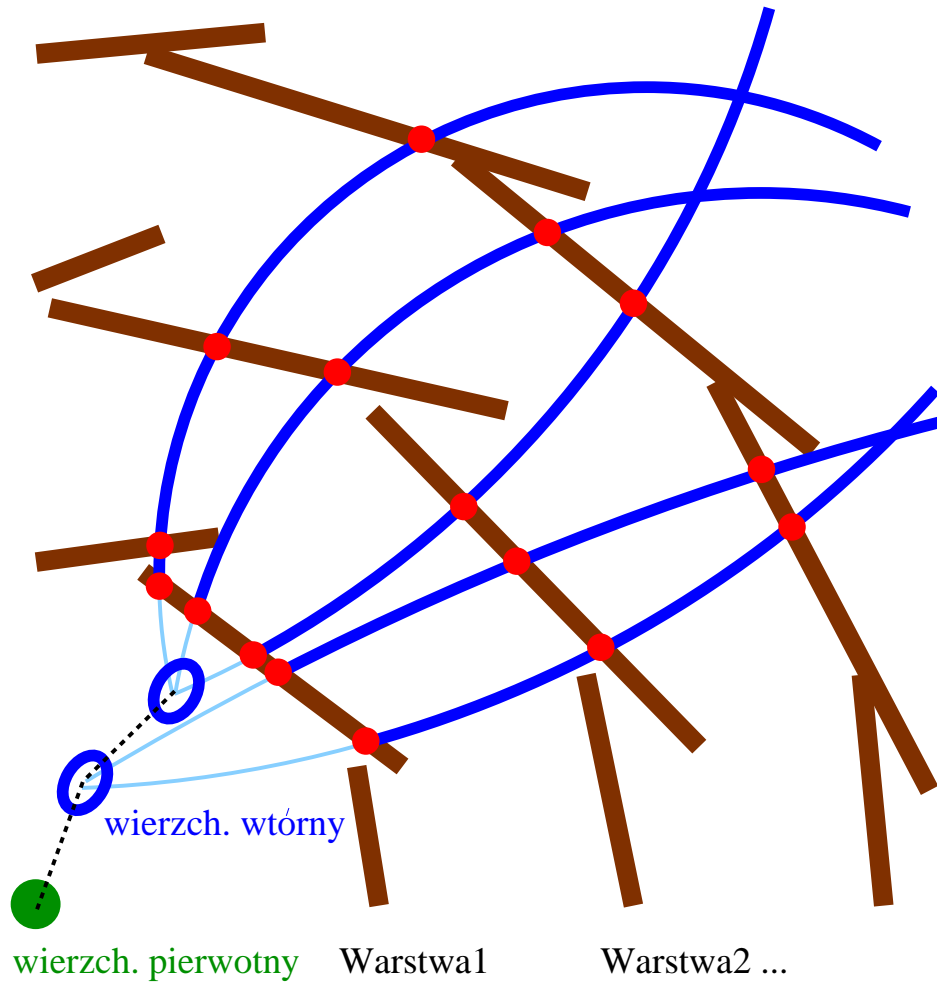


Several algorithms under active development

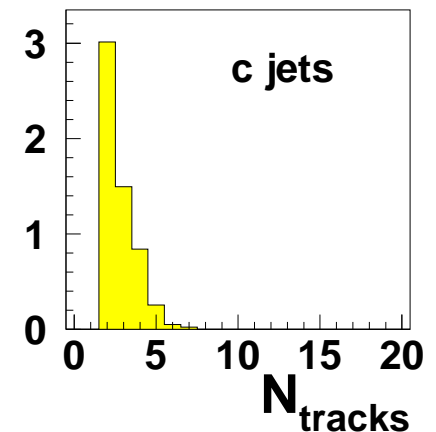
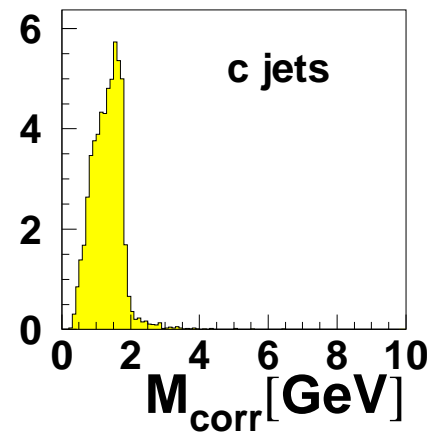
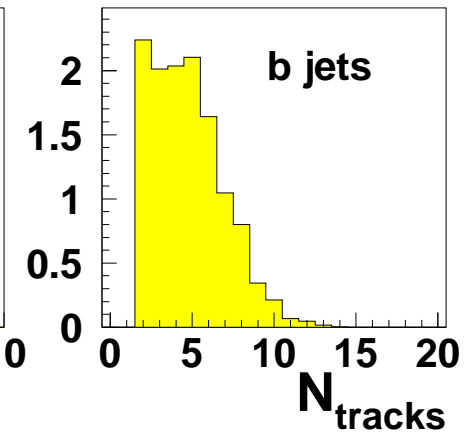
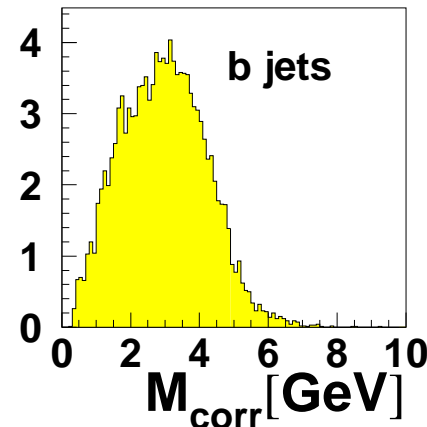
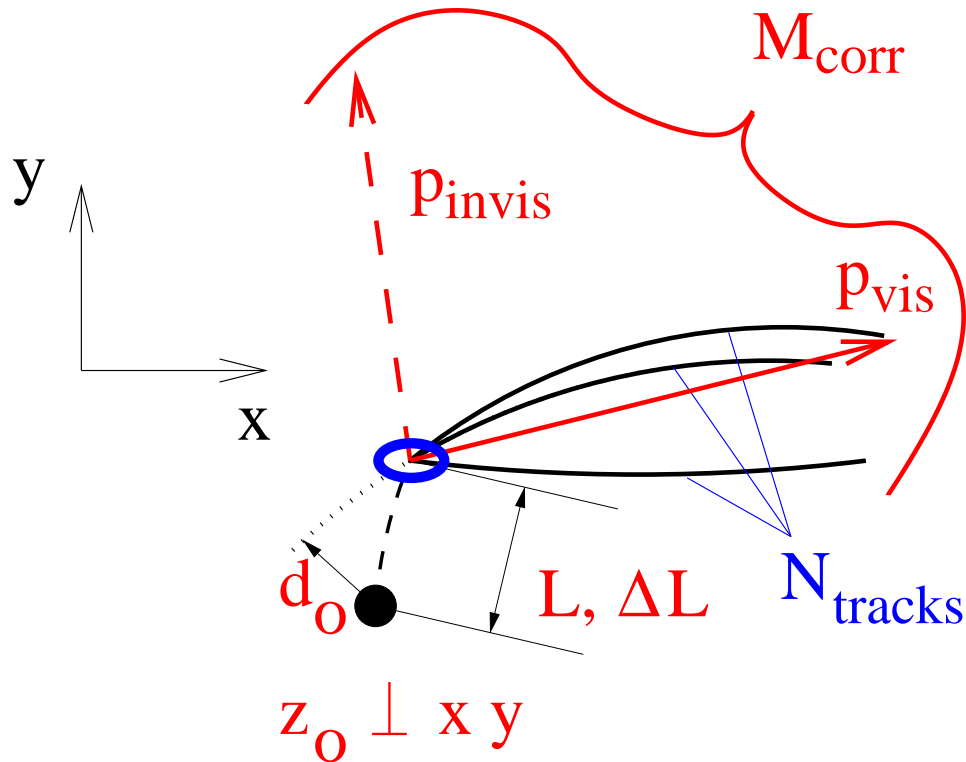
Impact Parameter Resolution	$d/\sigma(d)$
Decay Length Resolution	$L_{xy}/\sigma(L_{xy})$

“Jet tagging” - 1

- Dla każdego dżetu wierzchołki wtórne są wyznaczane za pomocą ZVTOP

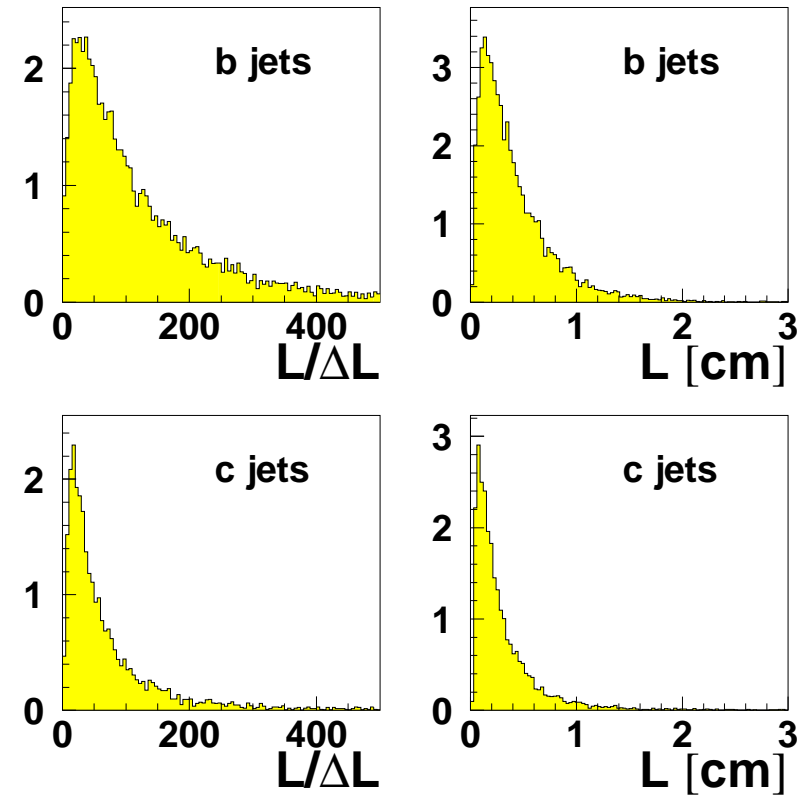
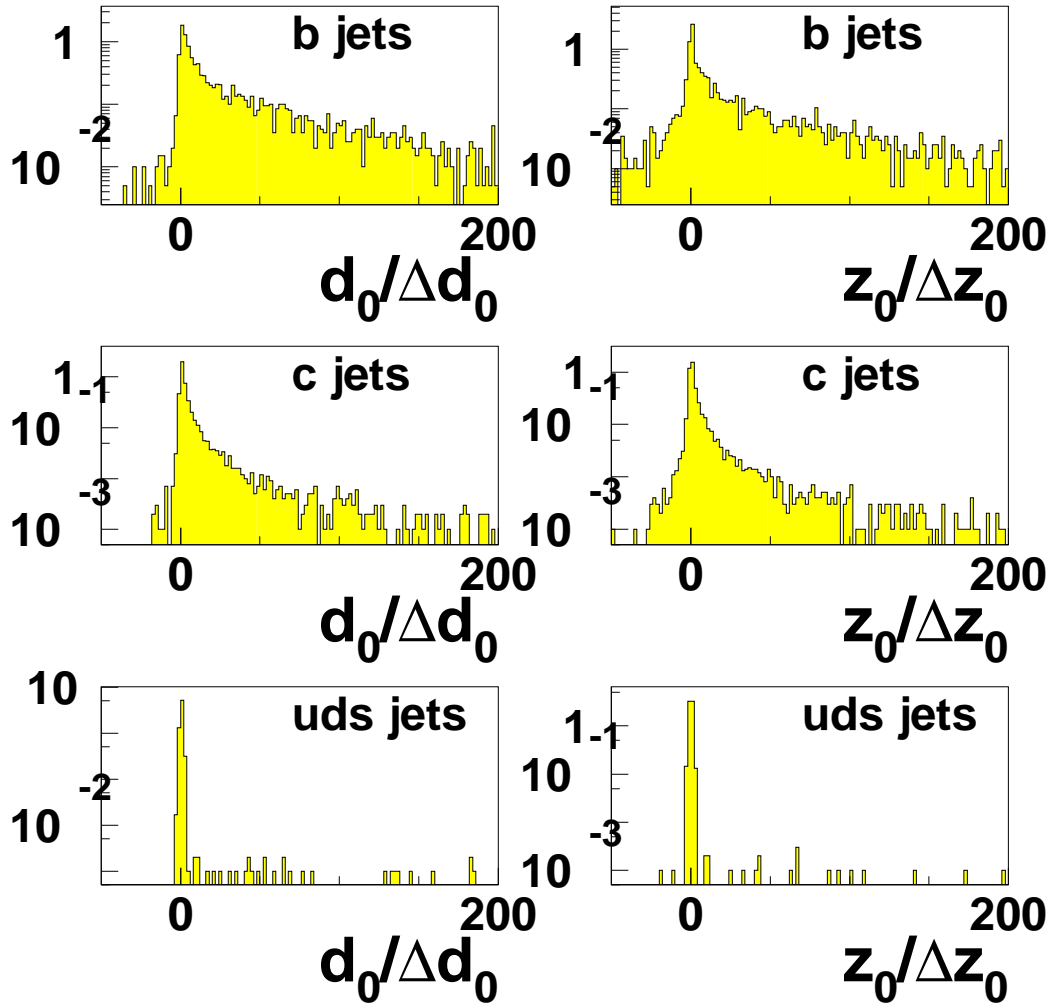
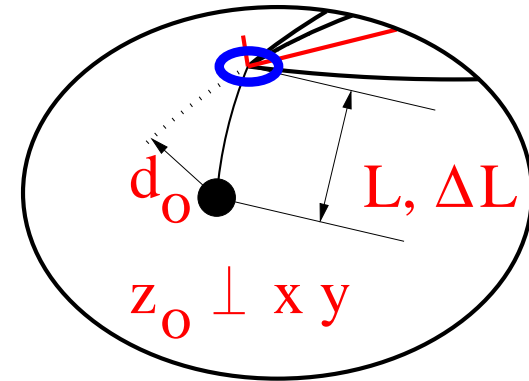


“Jet tagging” - 2



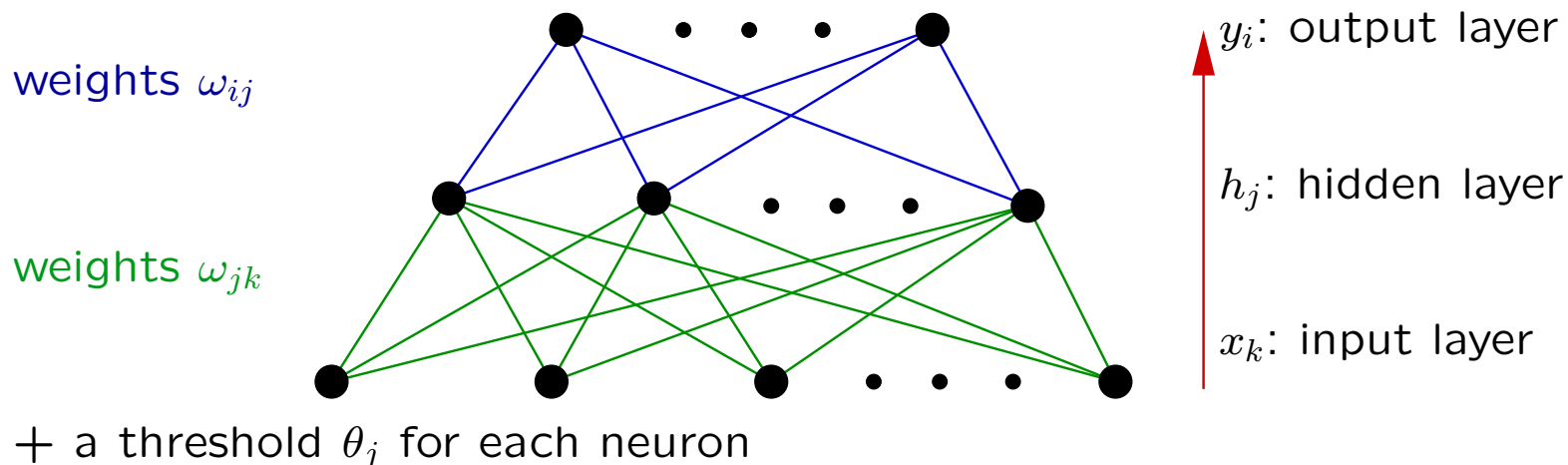
- Dla każdego dżetu **sieć neuronowa** (na podstawie powyższych wielkości) zwraca **b-tag** i **c-tag** $\in (0, 1)$ - wielkości mówiące na ile jet jest podobny do dżetu “b” lub “c”

“Jet tagging” - 3



What is a neural network?

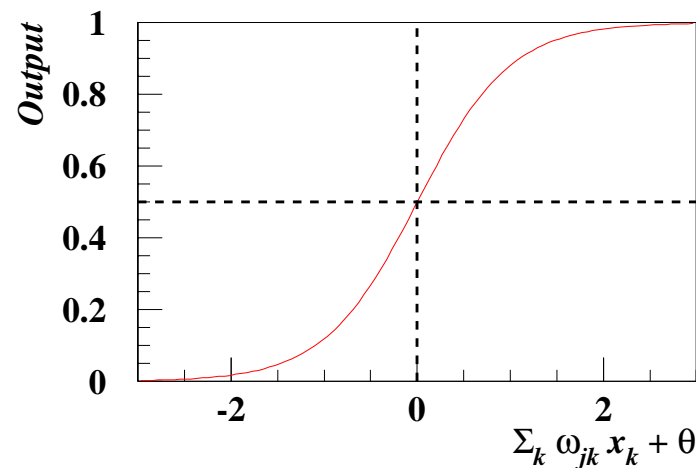
Feed-Forward Neural Network



Output of a neuron:

$$O_j = g\left(\sum_k \omega_{jk} x_k + \theta_j\right)$$

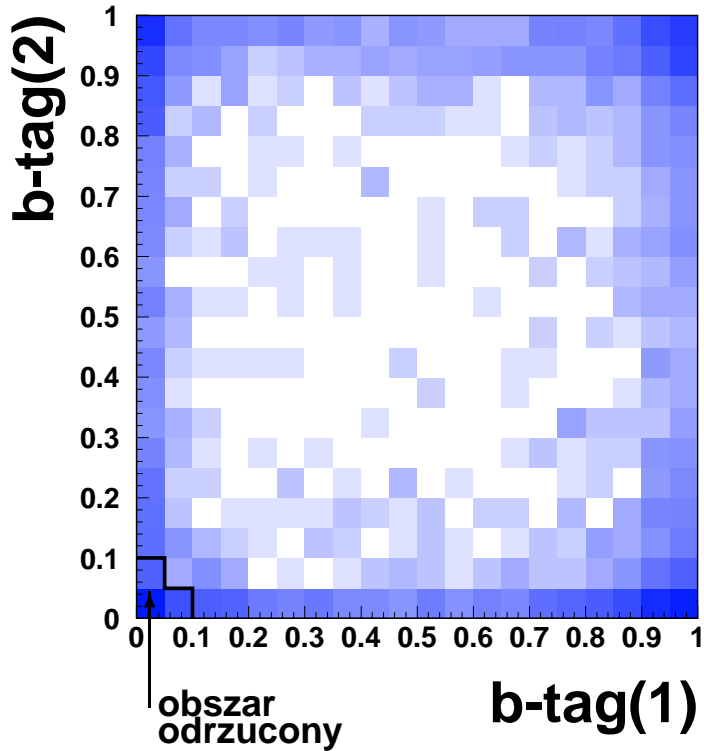
$$\text{with } g(x) = \frac{1}{1 + e^{-2x}}$$



Training: adjust weights
→ NN_{out} el. ~ 1 and had. ~ 0

b-tag dla $h \rightarrow \text{jet}(1) \text{jet}(2)$

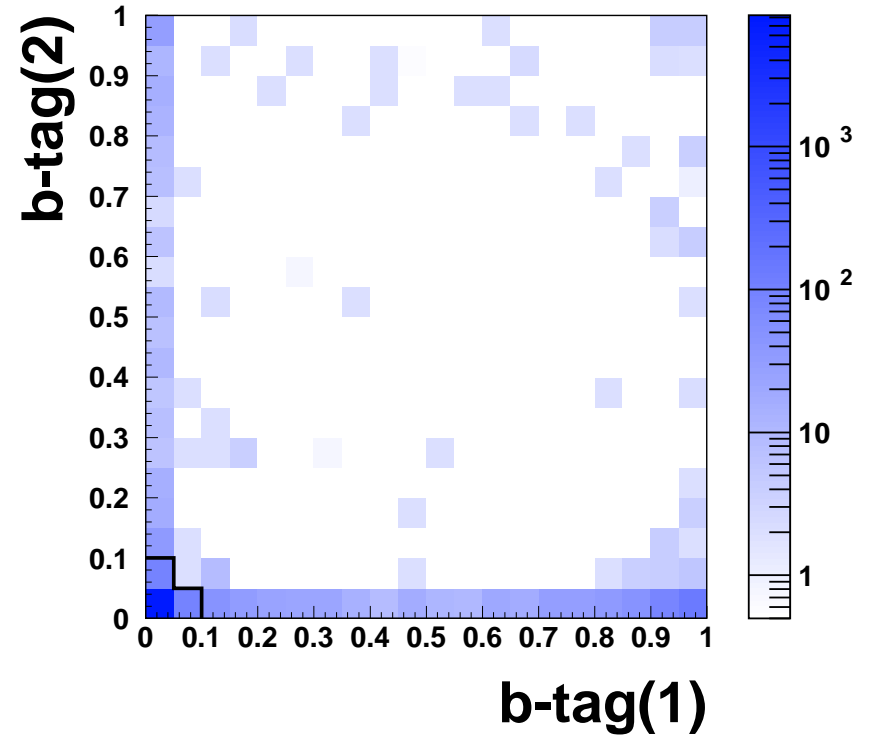
Signal: $h \rightarrow b\bar{b}$



$E(1) > E(2)$

500 fb^{-1}

Background



Cięcia optymalizujące $\Delta\Gamma/\Gamma$

Sygnal: 2414 przypadków

Tło: 785 przypadków

Precyzja pomiaru:

$$\frac{\Delta\Gamma(h \rightarrow b\bar{b})}{\Gamma(h \rightarrow b\bar{b})} = 2.3\%$$

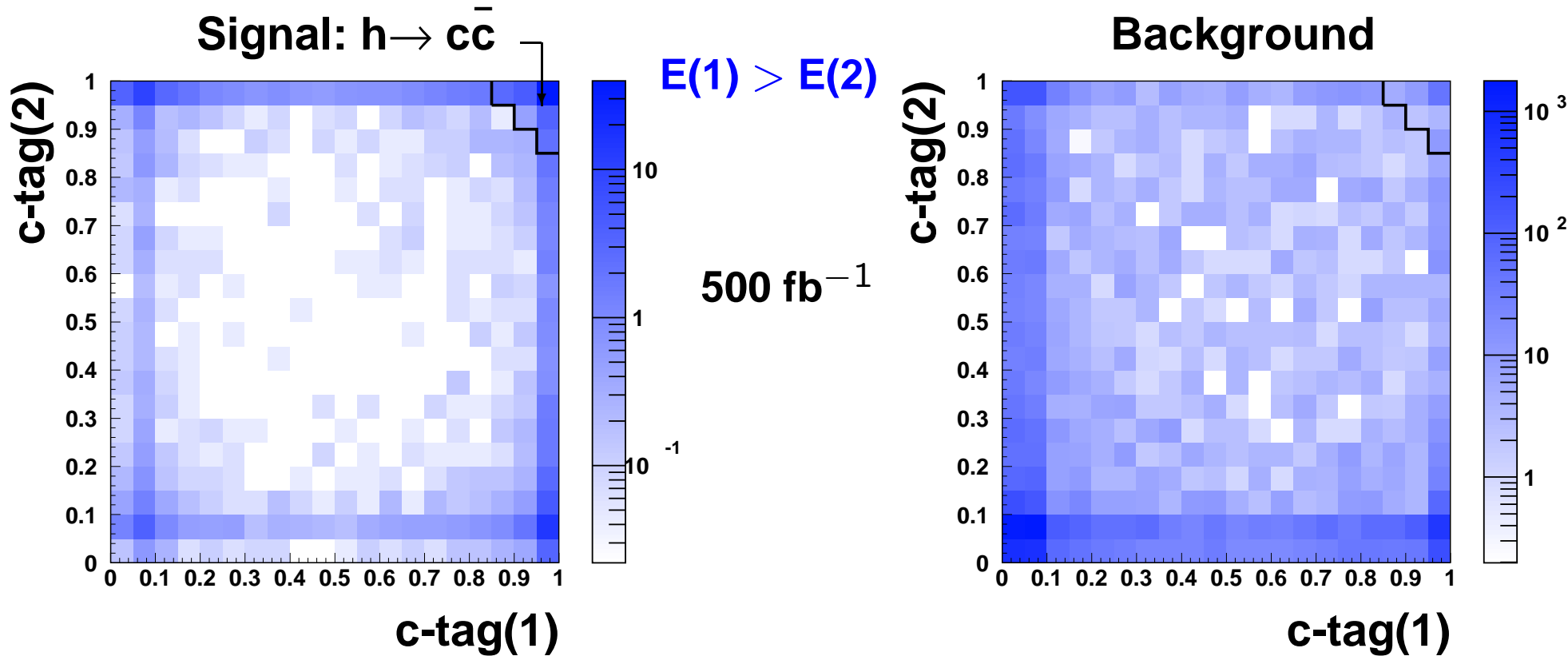
Przykładowe rozkłady dla :

warstwy: 1 - 5

grubość warstwy $50\mu\text{m}$

rozdzielczość warstwy $2\mu\text{m}$

c-tag dla $h \rightarrow \text{jet}(1) \text{jet}(2)$



Cięcia optymalizujące $\Delta\Gamma/\Gamma$

Sygnal: 42 przypadków

Tło: 79 przypadków

Precyzja pomiaru:

$$\frac{\Delta\Gamma(h \rightarrow c\bar{c})}{\Gamma(h \rightarrow c\bar{c})} = 26\%$$

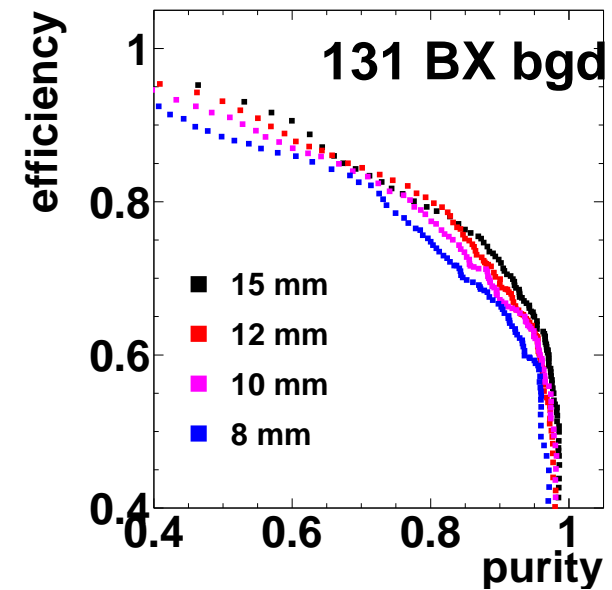
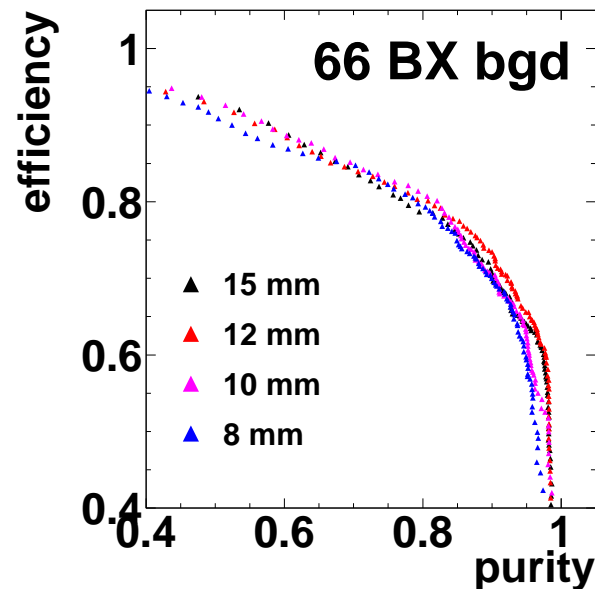
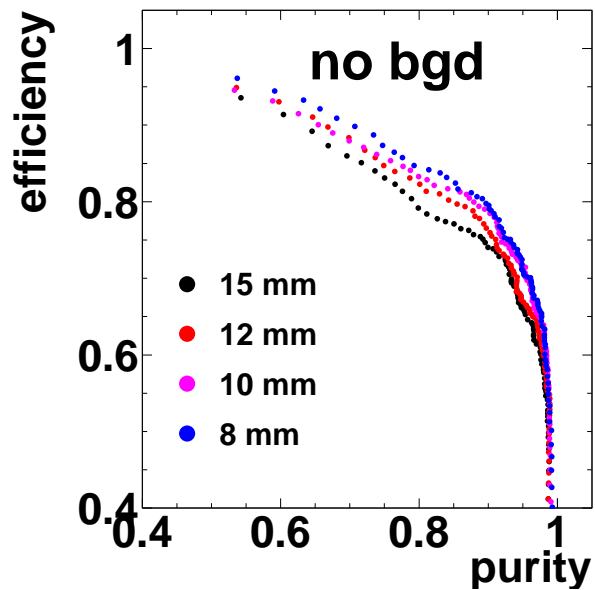
Przykładowe rozkłady dla :

warstwy: 1 - 5

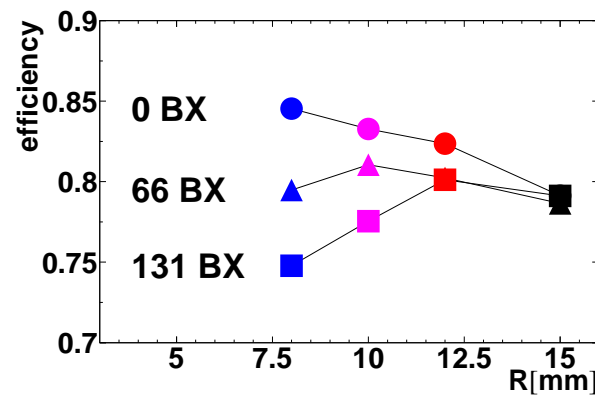
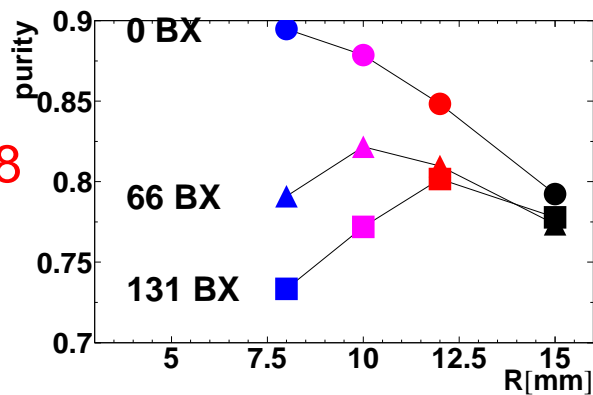
grubość warstwy $50\mu m$

rozdzielczość warstwy $2\mu m$

Wpływ tła na identyfikację kwarków b



Efficiency 0.8

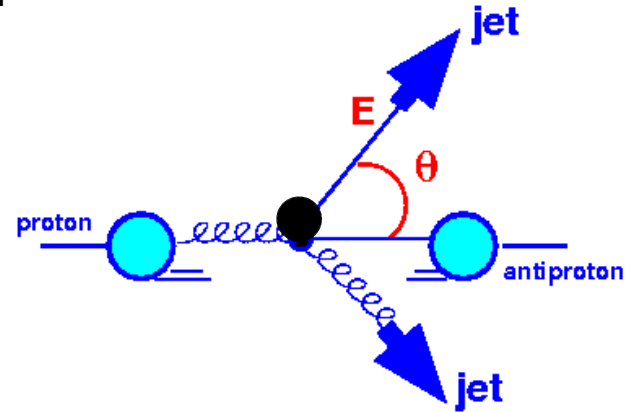
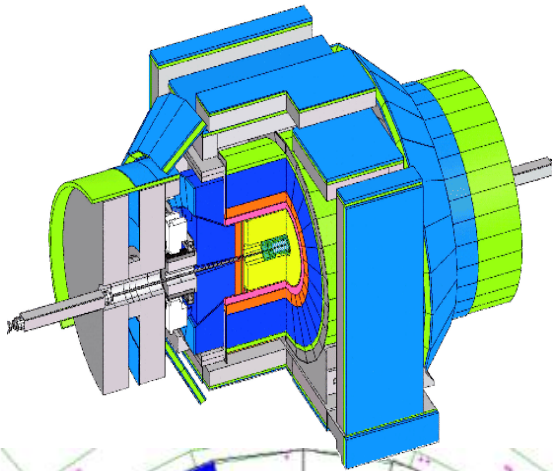


Purity 0.8

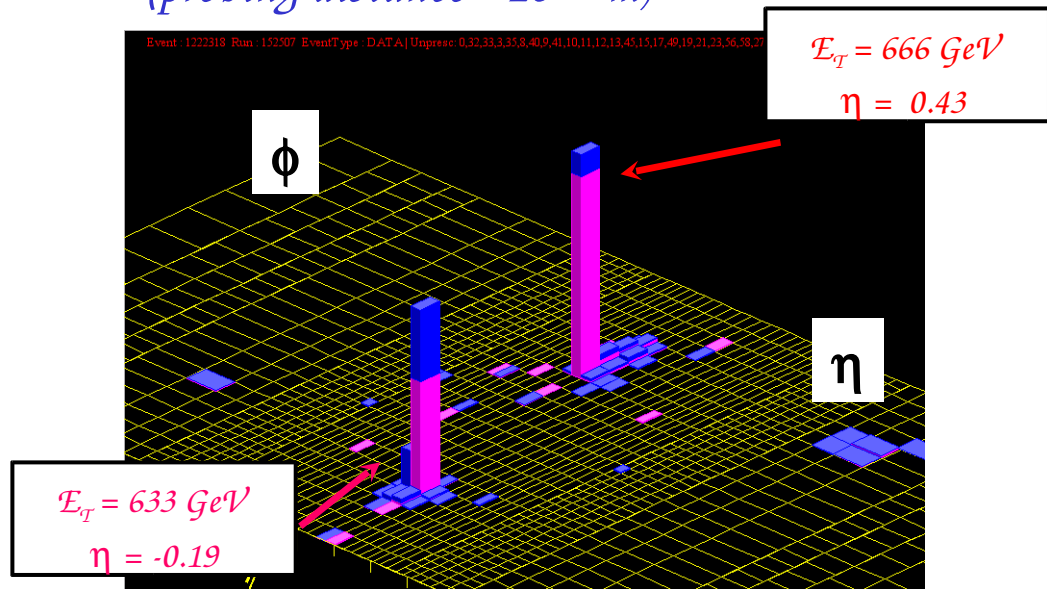
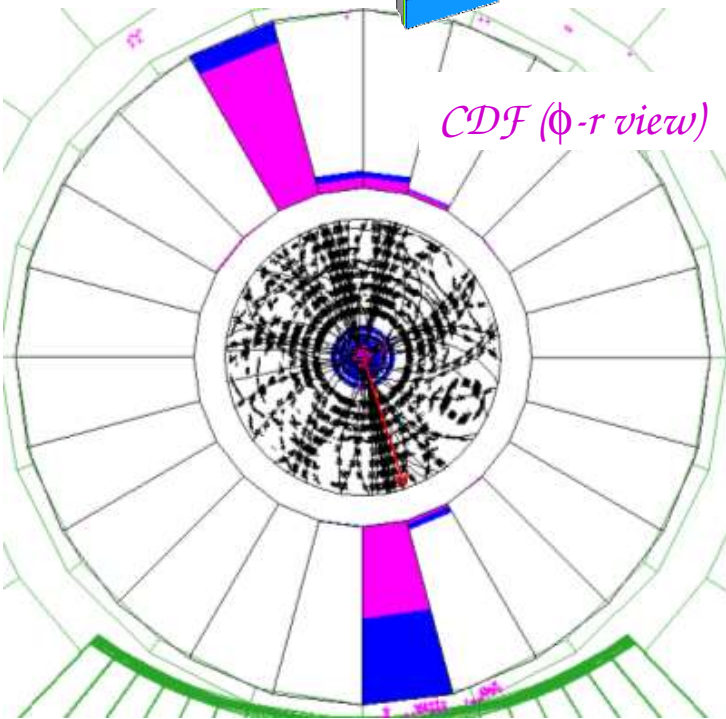
Rekonstrukcja jetów

What do Jets Look Like ?

(the highest P_T jet event from CDF)

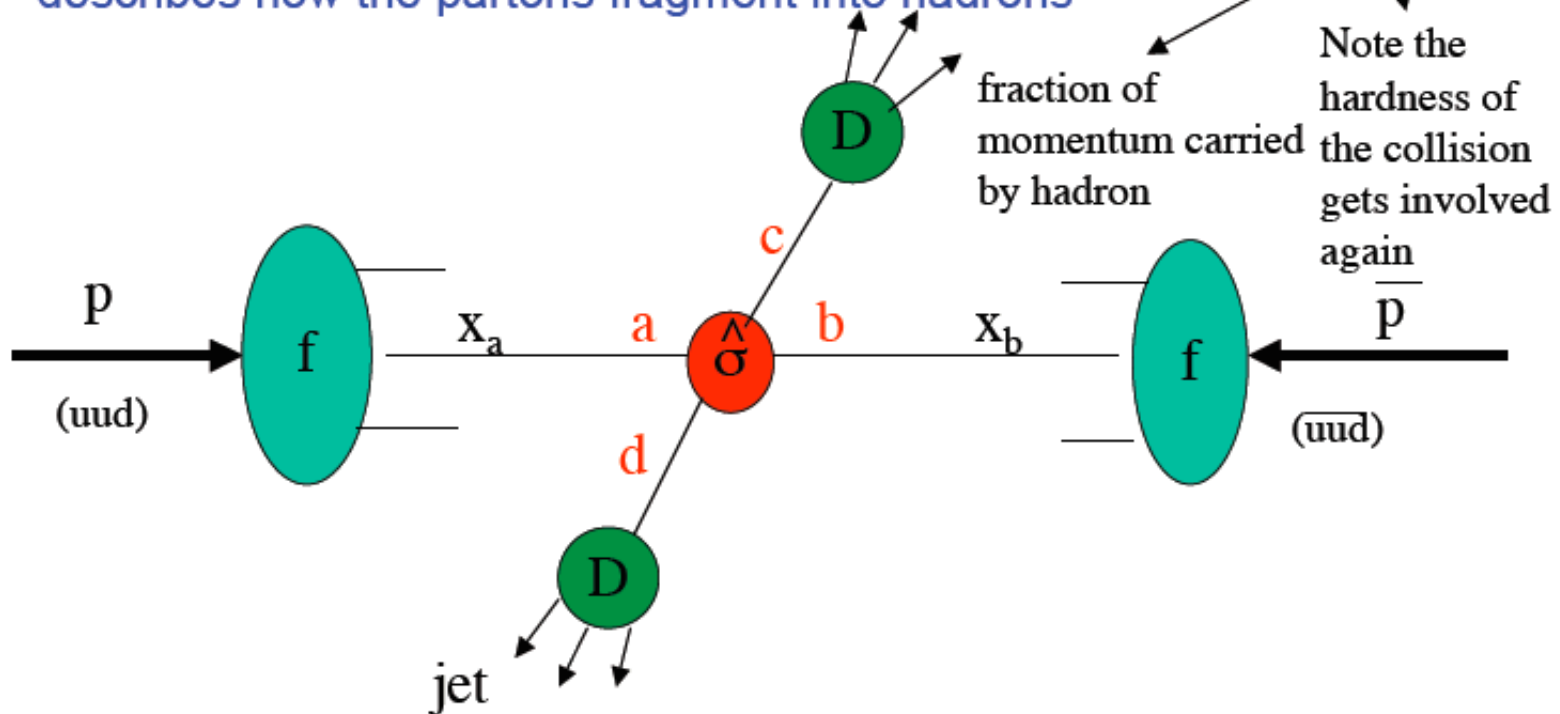


Dijet Mass = 1.36 TeV
(probing distance $\sim 10^{-19}$ m)



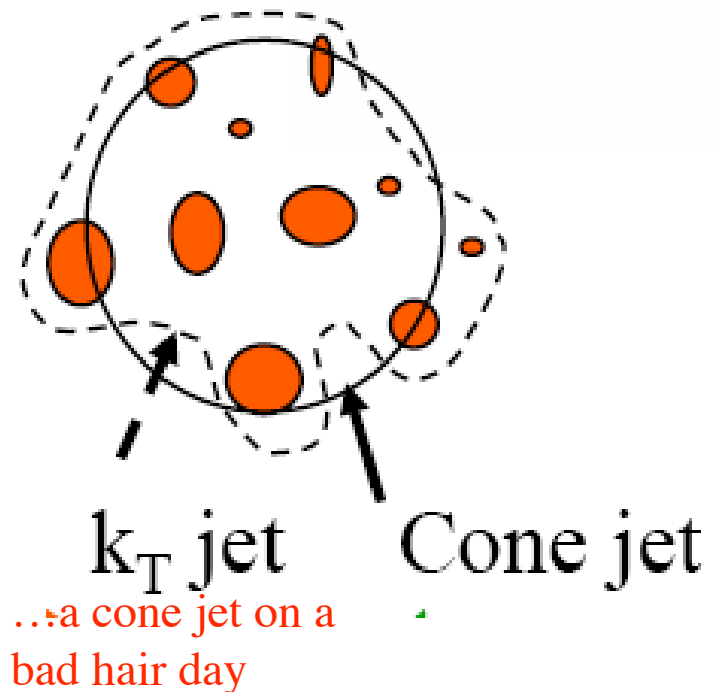
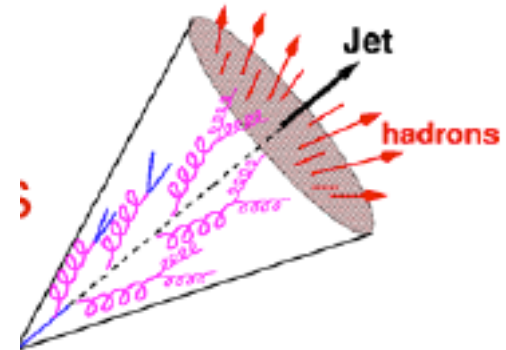
2->2 hard scattering

- c and d are the outgoing partons, but problem: they carry color; we don't observe the quarks or gluons themselves (don't see naked color) but instead the partons fragment into jets of hadrons (pions, kaons, and protons)
- it's the jets of hadrons that we can observe experimentally; $D(z, Q^2)$ describes how the partons fragment into hadrons



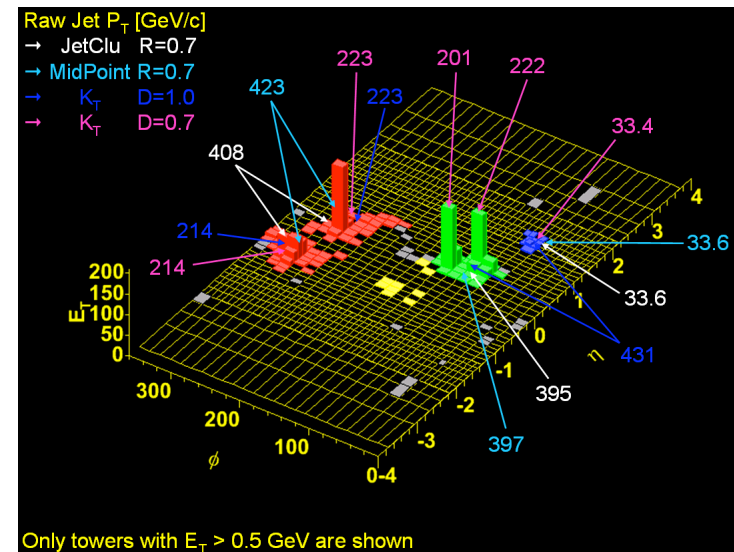
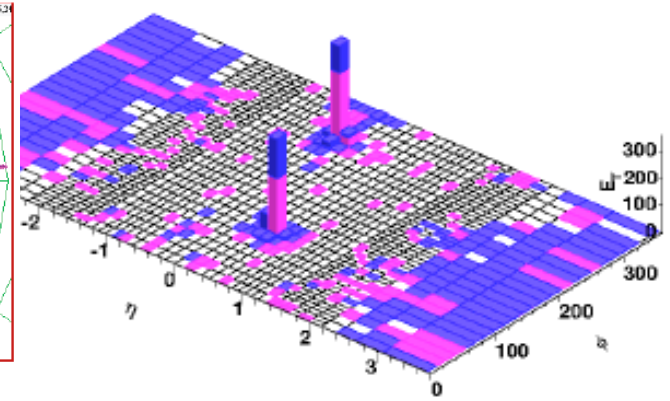
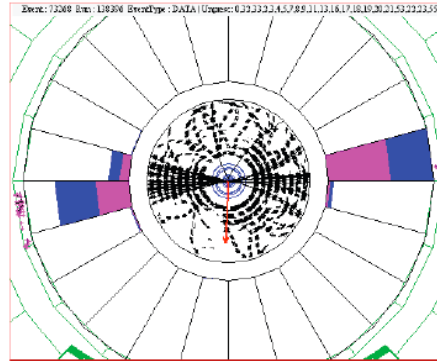
What is a jet?

- Jets are the experimental signatures of quarks and gluons
- Jets manifest themselves as localized clusters of energy
- It is the role of the jet algorithm to identify and measure the properties of a jet
- A jet algorithm can either measure
 - ◆ closeness in momentum space:
 k_T algorithm
 - ▲ most often used at LEP and HERA
 - ◆ closeness in coordinate space:
cone algorithm
 - ▲ most often used at the Tevatron
 - ◆ at the LHC, hopefully both will be equally used
- Can apply these jet algorithms to calorimeter towers or particles or partons...and would like to get a similar answers, as much as possible

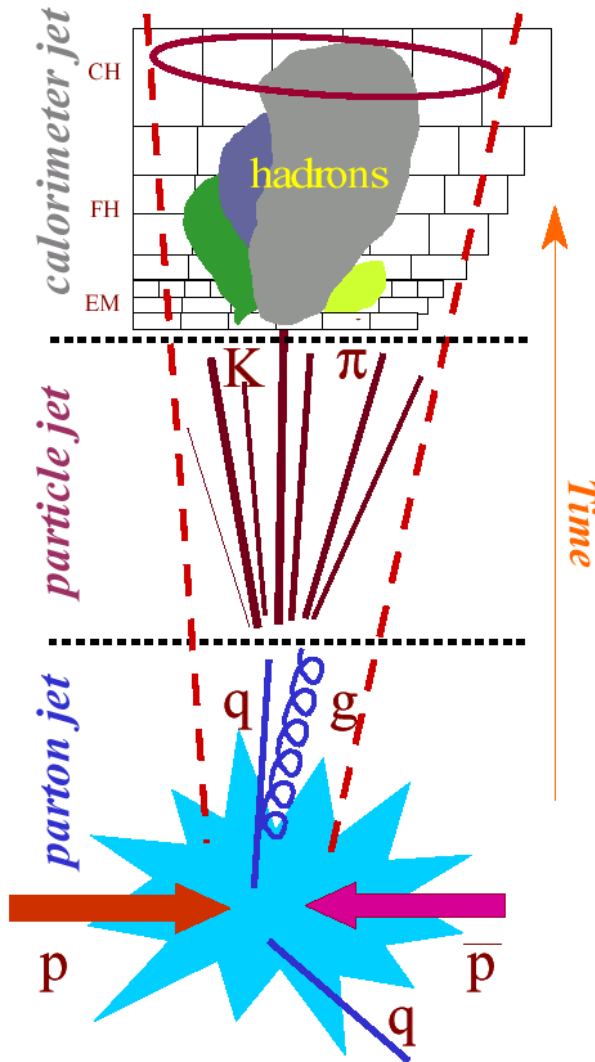


Jet algorithms

- For some events, the jet structure is very clear and there's little ambiguity about the assignment of towers/particles to the jet
- But for other events, there is ambiguity and the jet algorithm must make decisions that impact precision measurements
- If comparison is to hadron-level Monte Carlo, then hope is that the Monte Carlo will reproduce all of the physics present in the data and influence of jet algorithms can be understood
 - ◆ more difficulty when comparing to parton level calculations



Jet Finding



- **Calorimeter jet (cone)**

- ◆ jet is a collection of energy deposits with a given cone R : $R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$
- ◆ cone direction maximizes the total E_T of the jet
- ◆ various clustering algorithms

- correct for finite energy resolution
- subtract underlying event
- add out of cone energy

- **Particle jet**

- ◆ a spread of particles running roughly in the same direction as the parton after hadronization

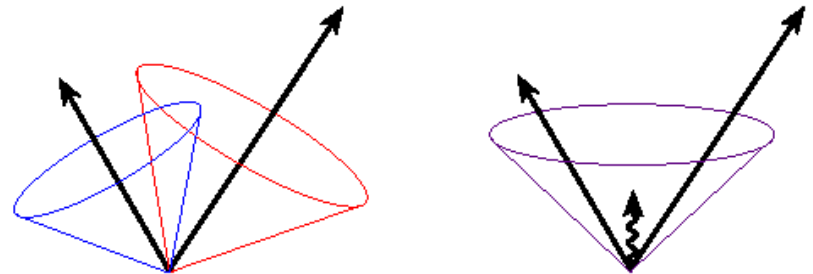
Jet Algorithms: Cone

- **Cone algorithms**

- ◆ draw a cone of fixed size around a seed
- ◆ compute jet axis
- ◆ draw a new cone around the new jet axis and recalculate axis and new E_T
- ◆ iterate until stable

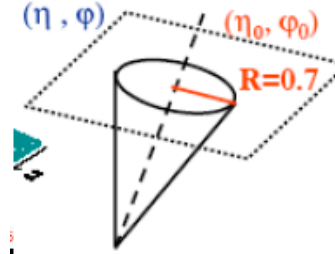
→ **In addition:**

- ◆ add additional midpoint seeds between pairs of close jets
- ◆ split/merge after stable proto-jets found



Midpoint cone algorithm

- Generate p_T ordered list of towers (or particles/partons)
- Find proto-jets around seed towers (typically 1 GeV) with $p_T >$ threshold (typically 100 MeV)



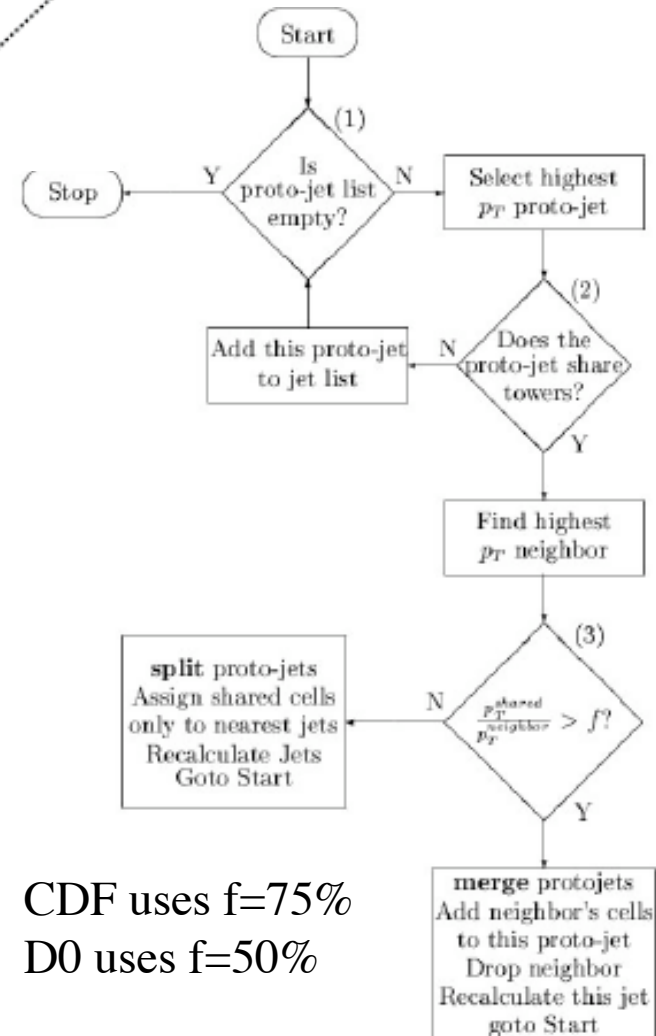
- ◆ include tower k in cone if

$$k \in C \text{ iff } \sqrt{(y_k - y_C)^2 + (\phi_k - \phi_C)^2} \leq R_{\text{cone}}$$

$$p_C = (E_C, \vec{p}_C) = \sum_{k \in C} (E_k, \vec{p}_k), \quad \bar{y}_C \equiv \frac{1}{2} \ln \frac{E_C + p_{z,C}}{E_C - p_{z,C}}, \quad \bar{\phi}_C \equiv \tan^{-1} \frac{p_{y,C}}{p_{x,C}}$$

- ◆ iterate if $(\bar{y}_C, \bar{\phi}_C) \neq (y_C, \phi_C)$
- ◆ NB: use of seeds creates IR-sensitivity

- Generate midpoint list from proto-jets
 - ◆ using midpoints as seed positions reduces IR-sensitivity
- Find proto-jets around midpoints
- Go to splitting/merging stage
 - ◆ real jets have spatial extent and can overlap; have to decide whether to merge the jets or to split them
- Calculate kinematics (p_T, y, ϕ) from final stable cones



CDF uses $f=75\%$

D0 uses $f=50\%$

Jet Algorithms: k_T

For each object and pair of objects:

order all d_{ii} and d_{ij} :

$$d_{ii} = k_{T,i}^2$$

$$d_{ij} = \underbrace{\min(k_{T,i}^2, k_{T,j}^2)}_{\text{Soft}} \frac{\Delta R_{ij}^2}{D^2}$$

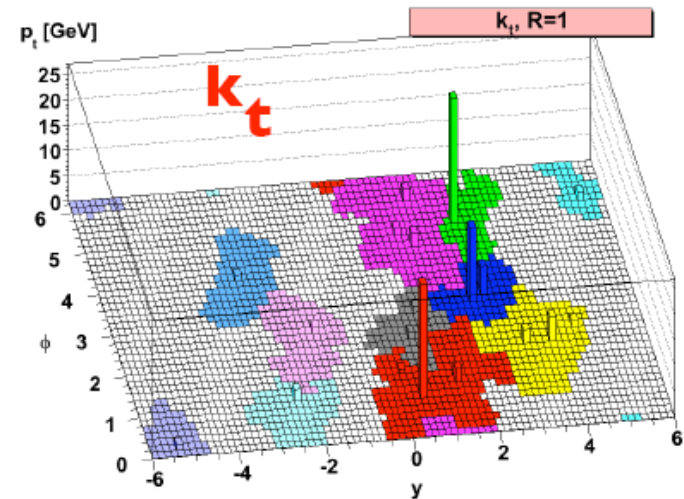
Collinear (if $\Delta R \ll 1$)

Resolution parameter ($D=1$)

If $d_{\min} = d_{ij}$
 \Rightarrow merge particles

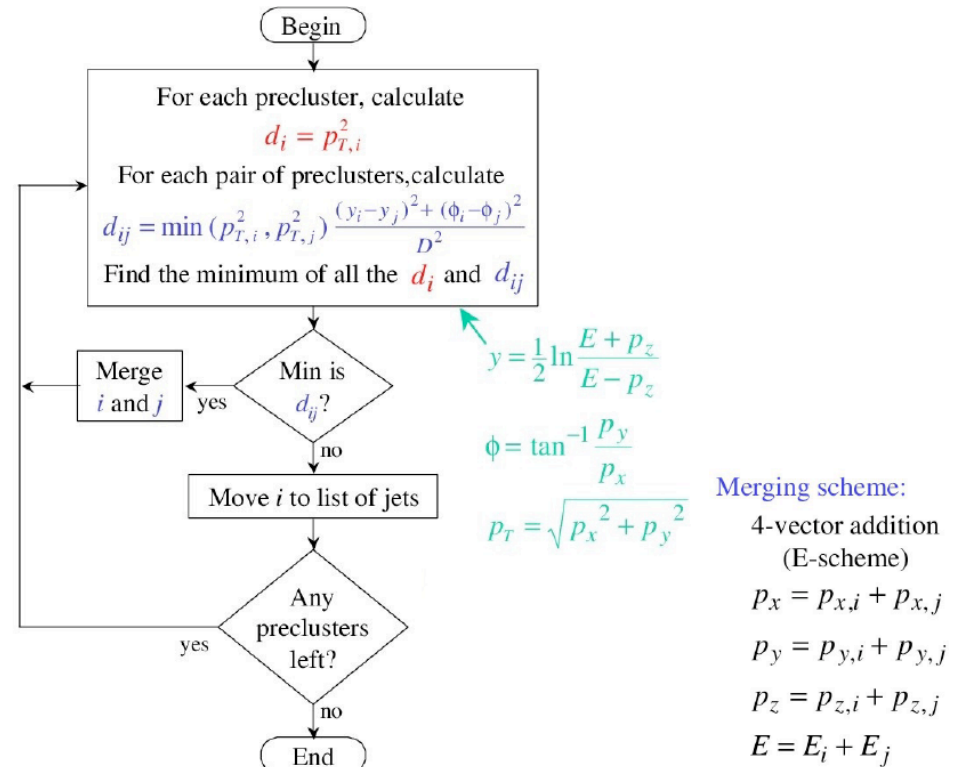
If $d_{\min} = d_{ii}$
 \Rightarrow jet

- theoretically favored, no split-merge
- to reduce computation time, start with 0.2×0.2 pre-clusters



k_T algorithm

- The k_T jet algorithm successively merges pairs of partons, particles or calorimeter towers in order of increasing relative transverse momentum
- The algorithm typically contains a parameter D that controls the termination of the merging and characterizes the approximate size of the resulting jets
- Since the k_T algorithm fundamentally merges nearby particles, there is a correspondence of jets reconstructed in a calorimeter to jets reconstructed from individual hadrons, leptons, and photons
- As the jet does not have a fixed area, the underlying event subtraction is more problematic



Rekonstrukcja energii jetów

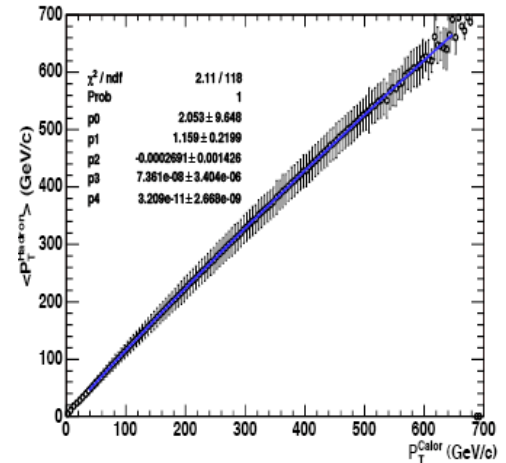
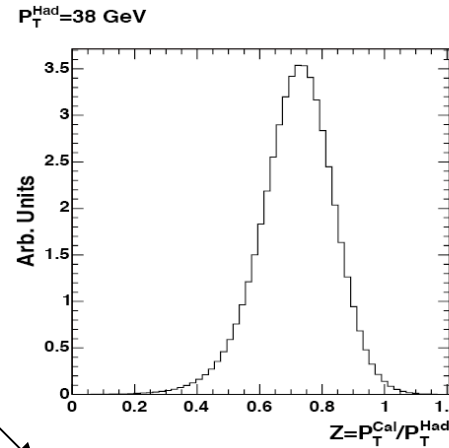
Jet Corrections

- Need to correct from calorimeter to hadron level (different response of calorimeter to EM and HAD energy)
 - ◆ and for resolution effects
- And from hadron to parton level for other observables (such as comparisons to parton level cross sections)

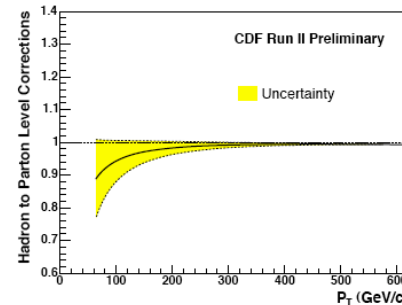
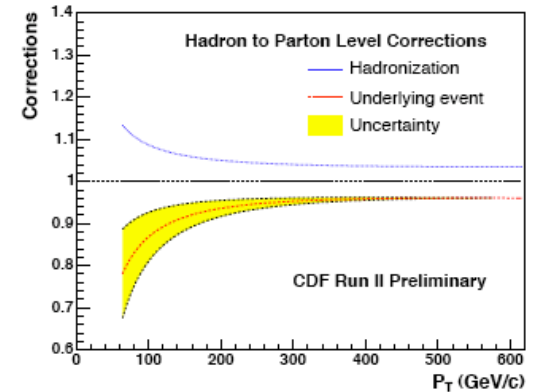
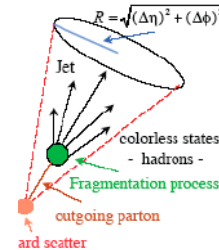
- ◆ underlying event and out-of-cone
 - ▲ can correct data to parton level or theory to hadron level...or both and be specific about what the corrections are

- ◆ note that loss due to hadronization is basically constant at 1 GeV/c for all jet p_T values at the Tevatron (for a cone of radius 0.7)
 - ▲ for a cone radius of 0.4, the two effects cancel to within a few percent

- ◆ interesting to check over the jet range at the LHC



partons in cone give rise to hadrons outside the cone



for cone of 0.7, UE correction wins

Response correction

One possibility:

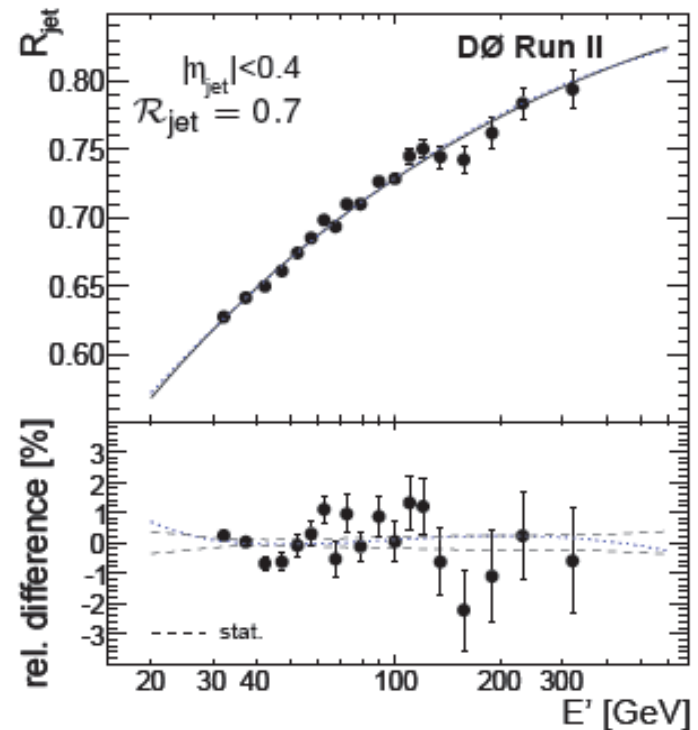
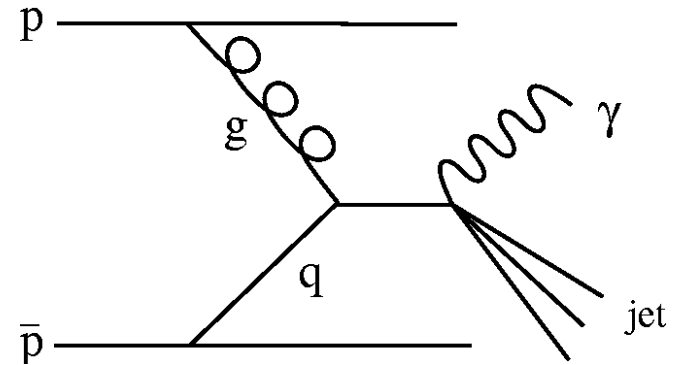
use γ +jet events

- γ well measured and calibrated

- missing transverse energy projection:

$$R_{\text{had}} = 1 + \frac{\vec{E}_T \cdot \vec{p}_{T,\gamma}}{\vec{p}_{T,\gamma}^2}$$

The 25% correction to the response is the largest of all energy scale contributions



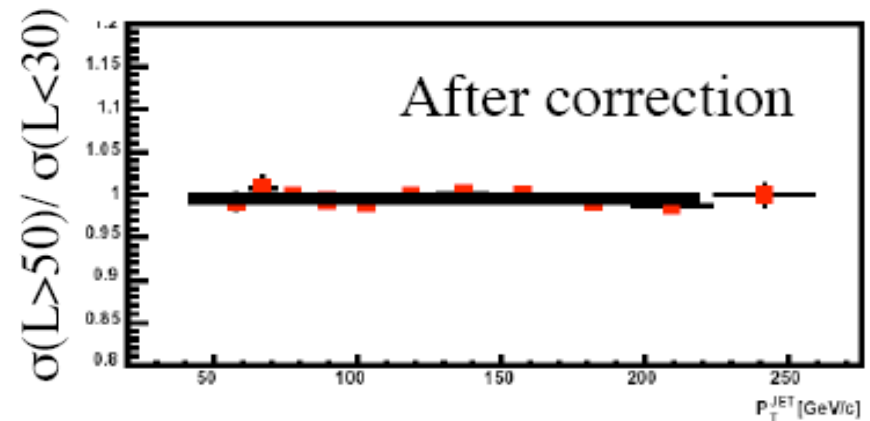
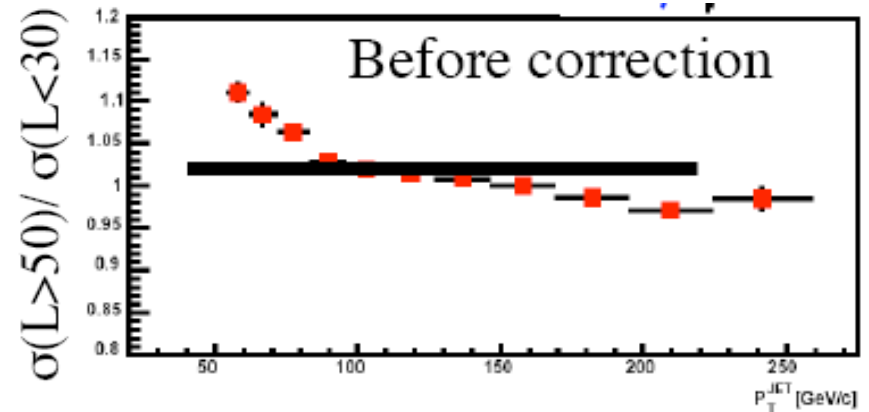
JES for k_T algorithm

Problem of k_T algorithm in hadron colliders: Multiple interactions

- k_T algorithm is attracted to energy deposits and picks up the energy from MI, UE etc.

- k_T algorithm has now been used by CDF to measure a cross-section for Run2:

- empirical energy correction factor used using the fact that the cross-section is luminosity independent

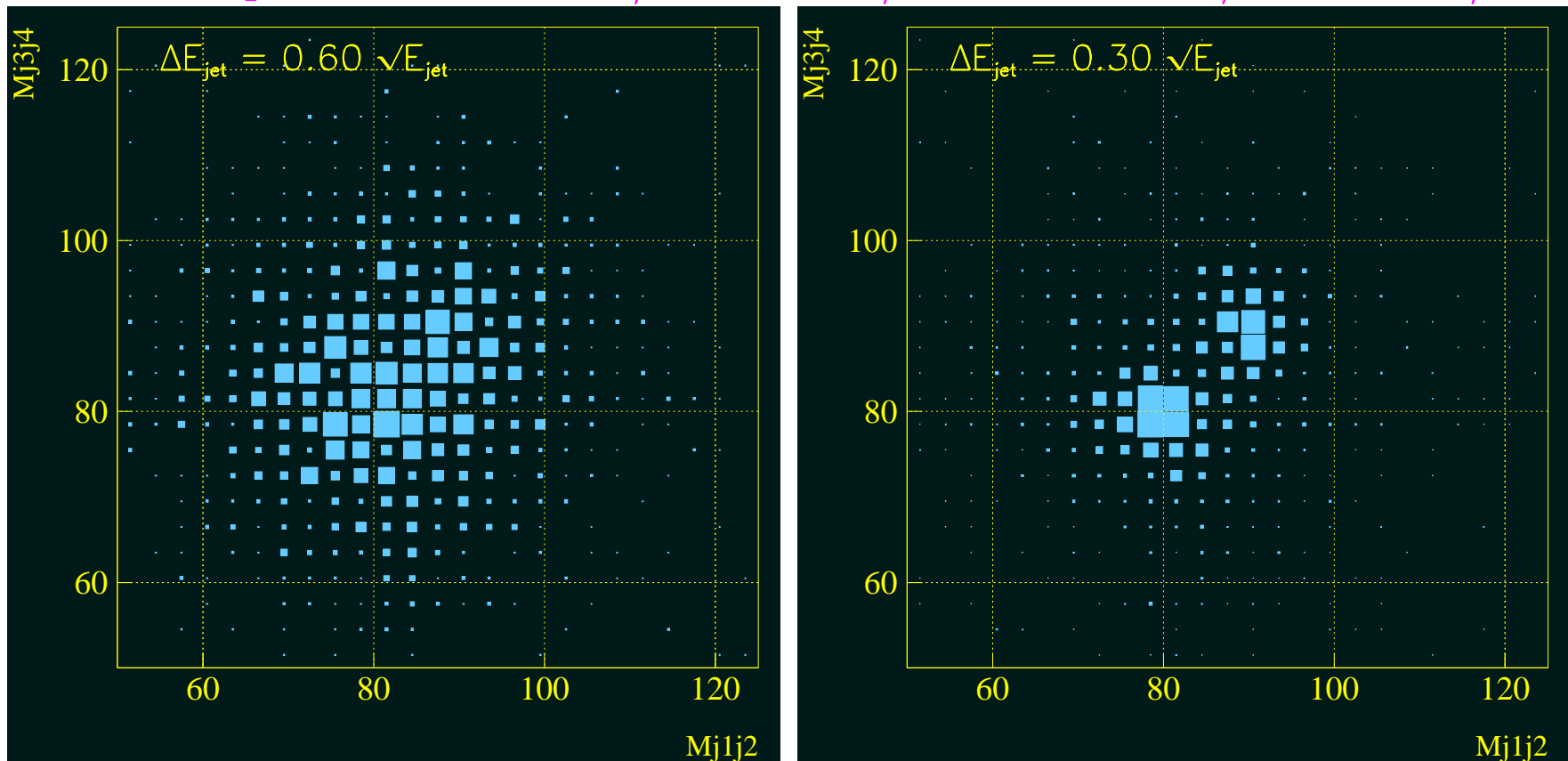


Rekonstrukcja obiektów

Energy flow in jets

- Some processes where WW and ZZ need to be separated without beam constraints (e.g. $e^+e^- \rightarrow \nu\nu WW, \nu\nu ZZ$)
- This requires a resolution of about $\Delta E/E = 30\%/\sqrt{E}$

WW-ZZ separation for $\Delta E/E = 60\%/\sqrt{E}$ and $\Delta E/E = 30\%/\sqrt{E}$



Jet Reconstruction @ ILC

- **Q.** How to achieve the best attainable jet energy resolution?
- **A.** Since the momentum resolution for the charged particle measured by trackers is much better than the energy resolution of calorimeters, the best energy resolution is obtained by reconstructing momenta of individual particles **avoiding double counting** among **Trackers** and **Calorimeters**.
 - Charged particles (~60%) measured by Tracker.
 - Photons (~30%) by electromagnetic CAL (ECAL).
 - Neutral hadrons (~10%) by ECAL + hadron CAL (HCAL).

$$\mathbf{E}_{\text{TOT}} = \mathbf{p}_e + \mathbf{p}_\mu + \mathbf{p}_{\text{charged hadron}} + \mathbf{E}_\gamma + \mathbf{E}_{\text{neutral hadron}}$$

[tracks only] [calorimeter only]

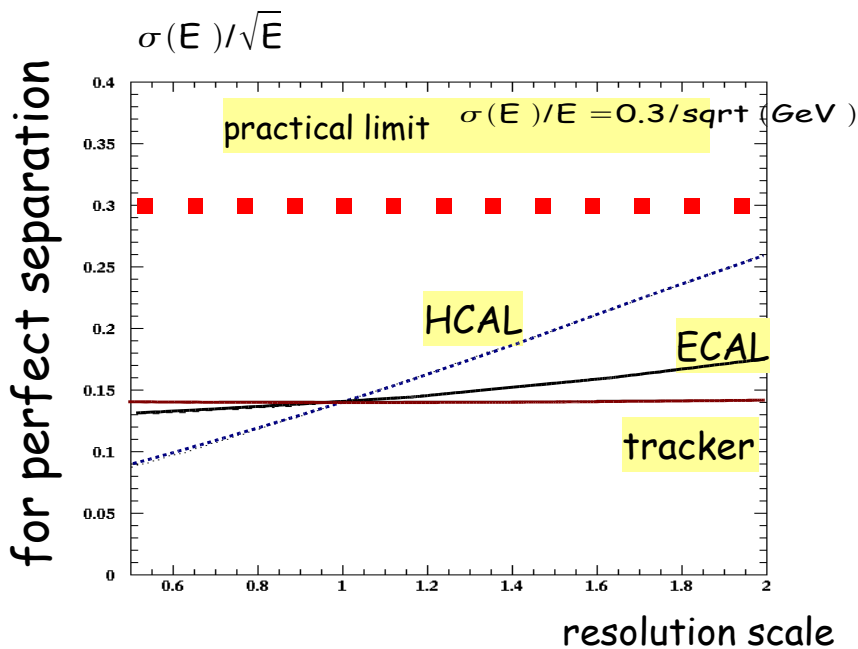
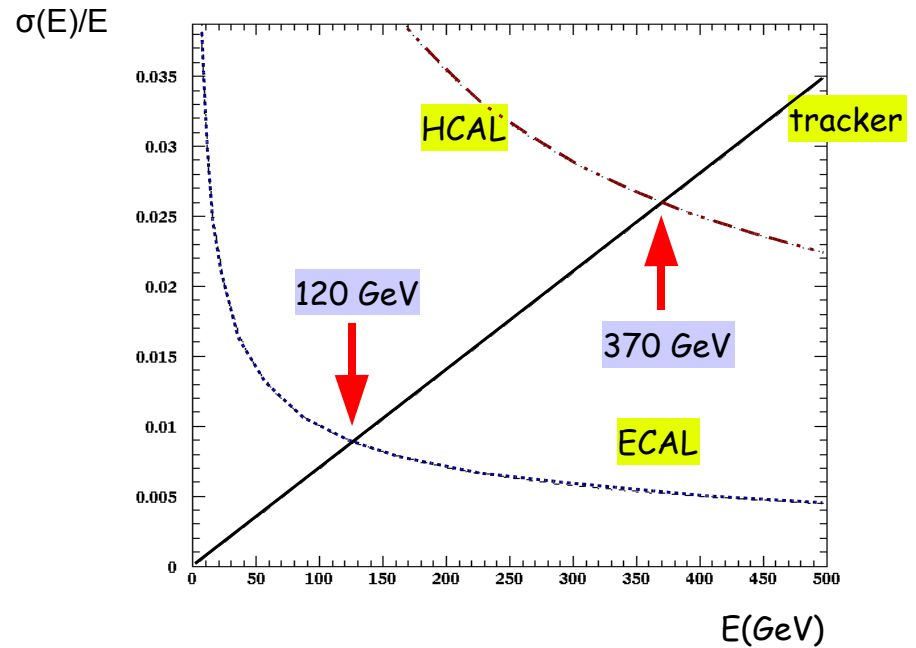
Particle Flow Algorithm (PFA)

Particle Flow: Basics

$$\sigma(\text{Jet}) = \sqrt{\sum \epsilon_T^2 E_i^4 + \sum \epsilon_{\text{ECAL}}^2 E_i + \sum \epsilon_{\text{HCAL}}^2 E_i}$$

Resolution is dominated by **HCAL**
and by
"confusion" term

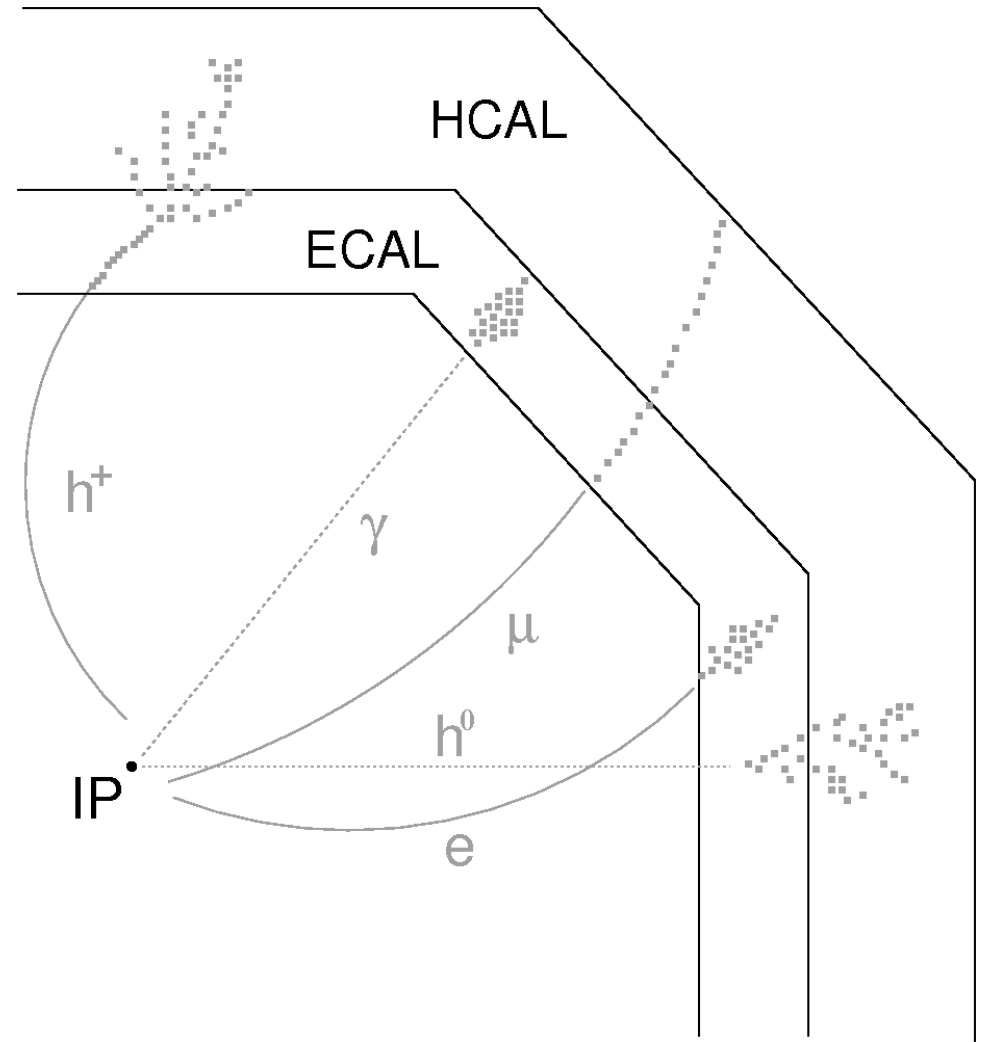
Resolution tracker - Calorimeter



- design detector to
- minimize confusion term
 - minimize the role of the HCAL
 - for the rest: build the best HCAL possible

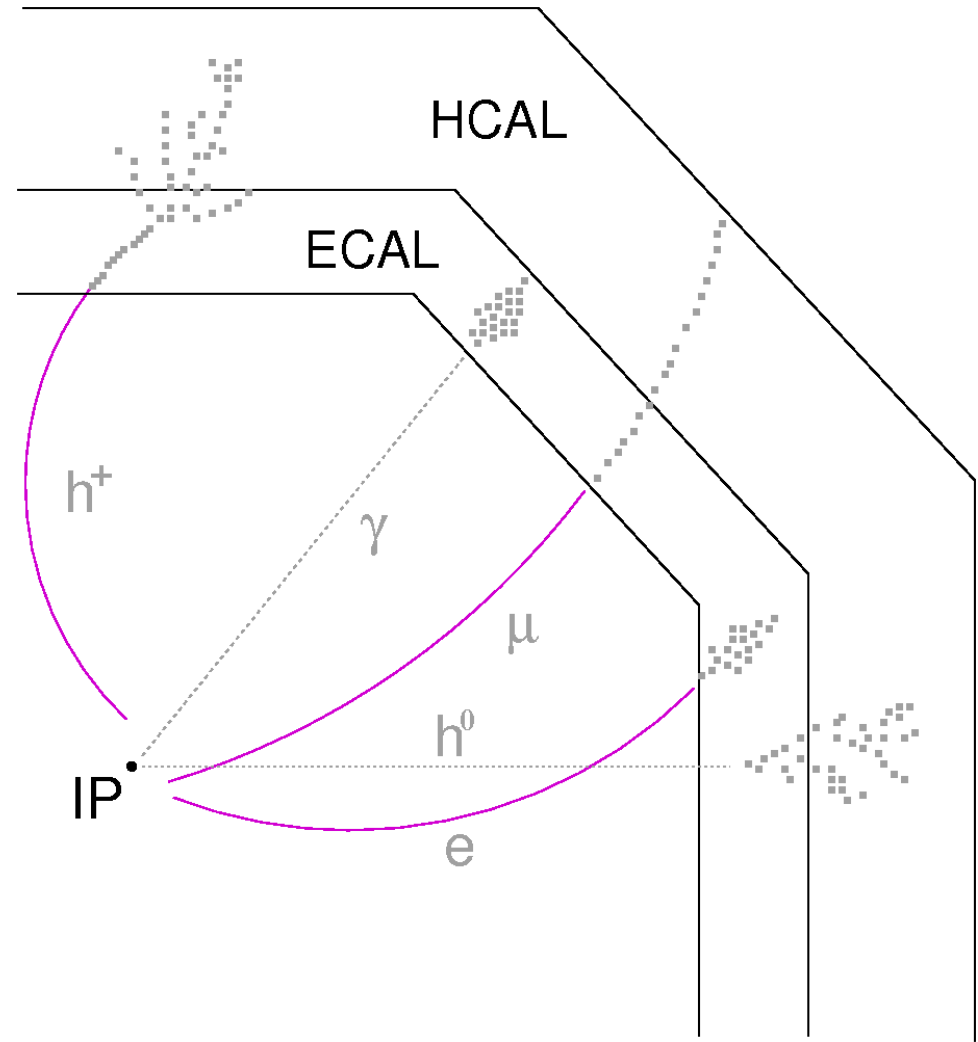
Effect of changing the
resolutions by a scale factor

Track-Based Particle Flow Concept



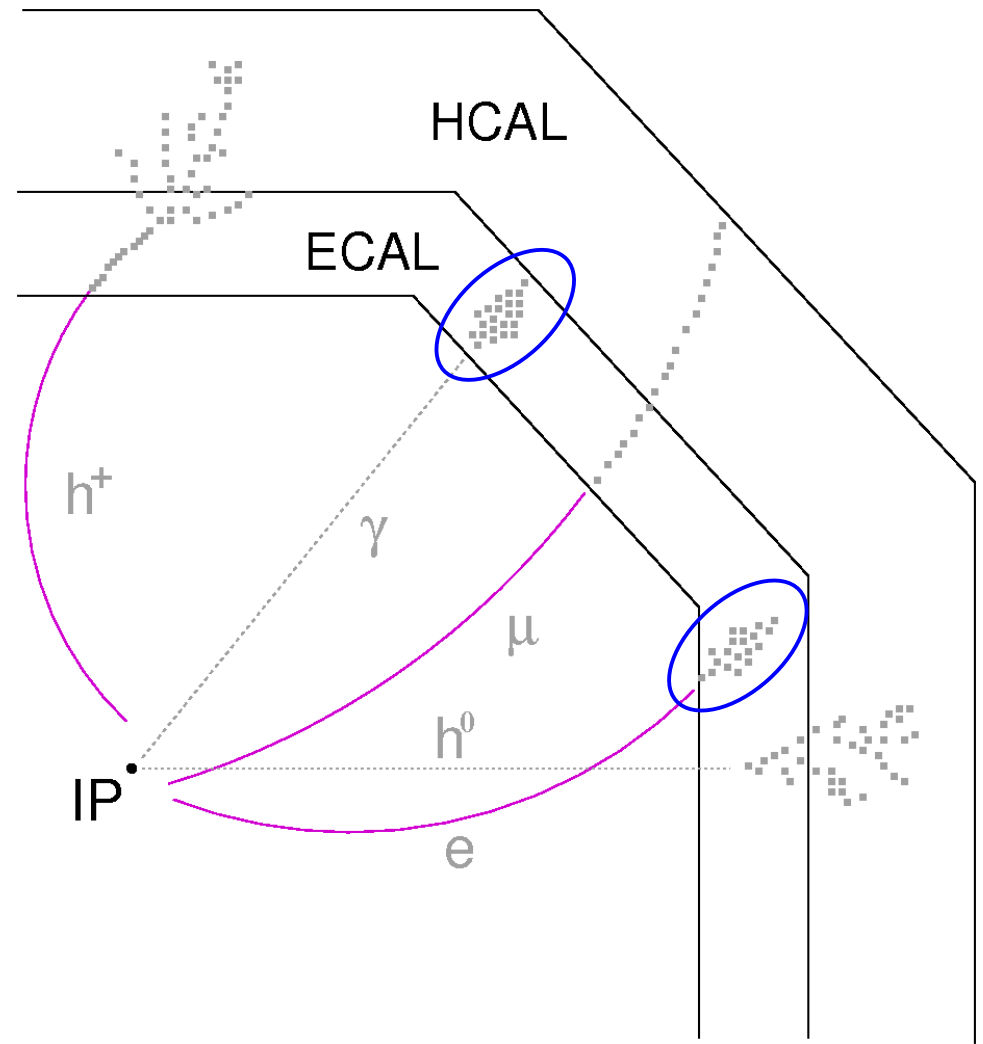
Track-Based Particle Flow Concept

1. tracking (Silicon and TPC)



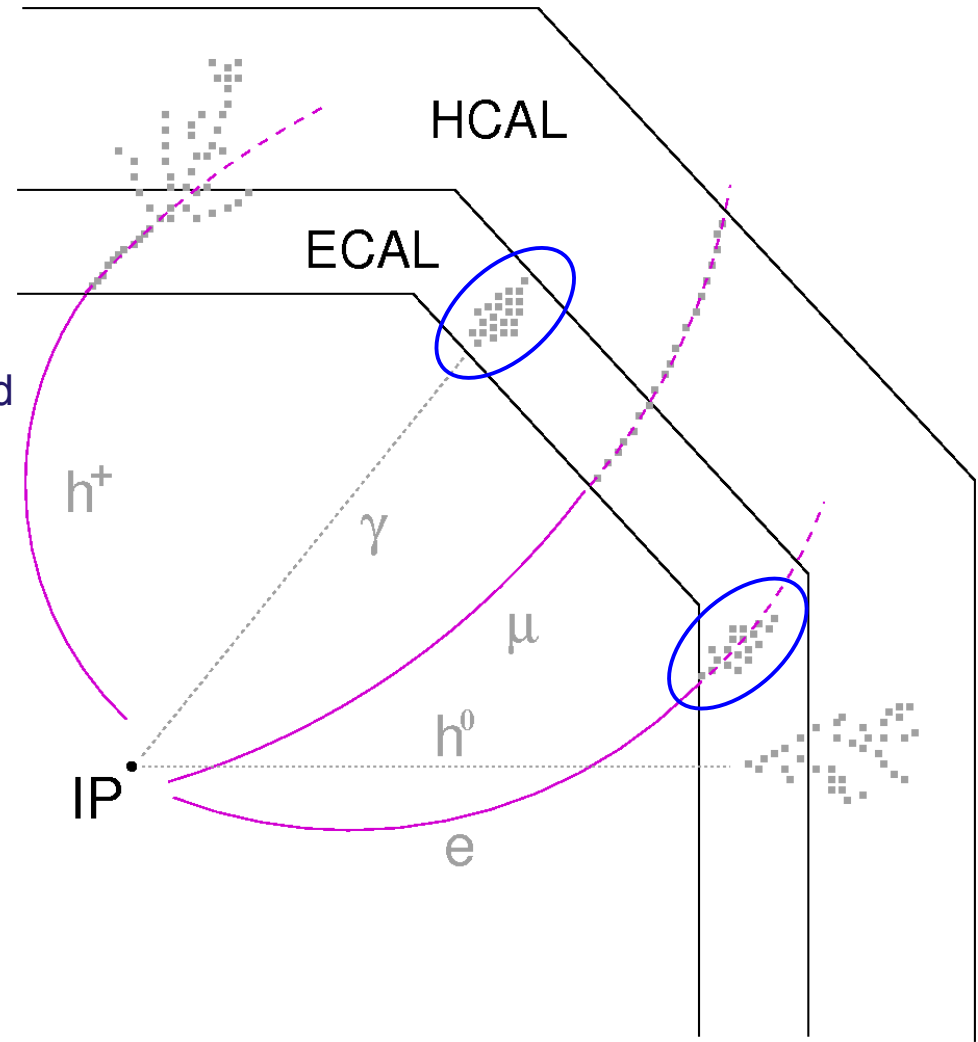
Track-Based Particle Flow Concept

1. tracking (Silicon and TPC)
2. find photon candidates



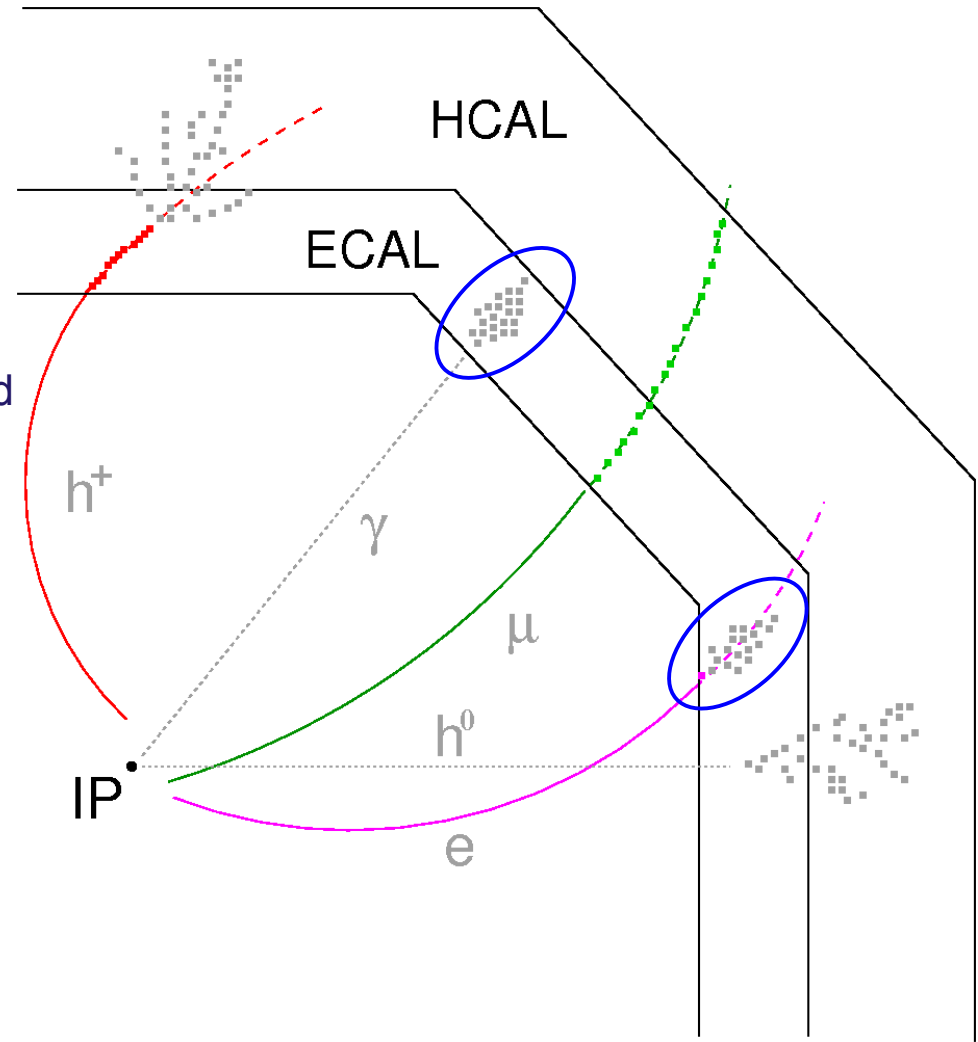
Track-Based Particle Flow Concept

1. tracking (Silicon and TPC)
2. find photon candidates
3. extrapolate tracks into Calorimeter
 - different models, with and w/o energy loss, multiple scattering, ...
 - dedicated Geometry description needed



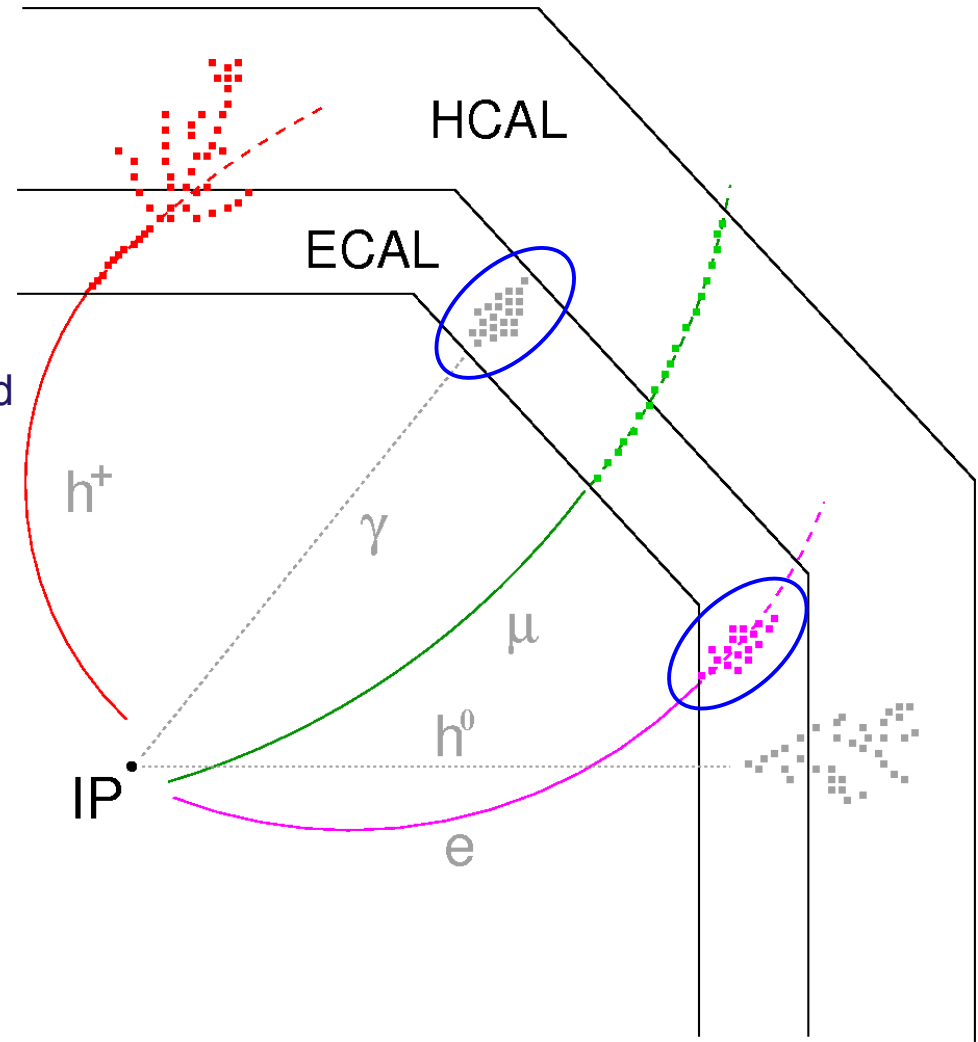
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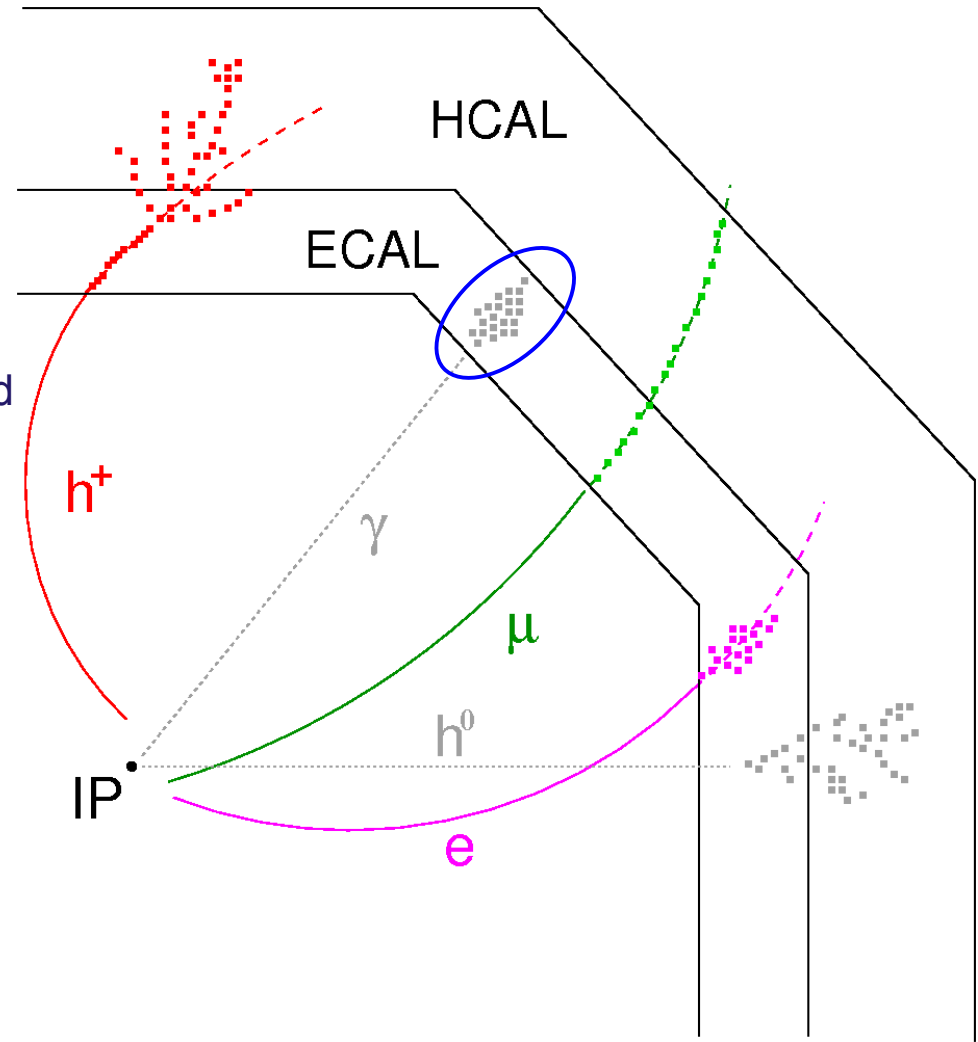
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5. clustering (ECAL and HCAL)
 - variable, depending on track and photon candidates
 - different algorithms



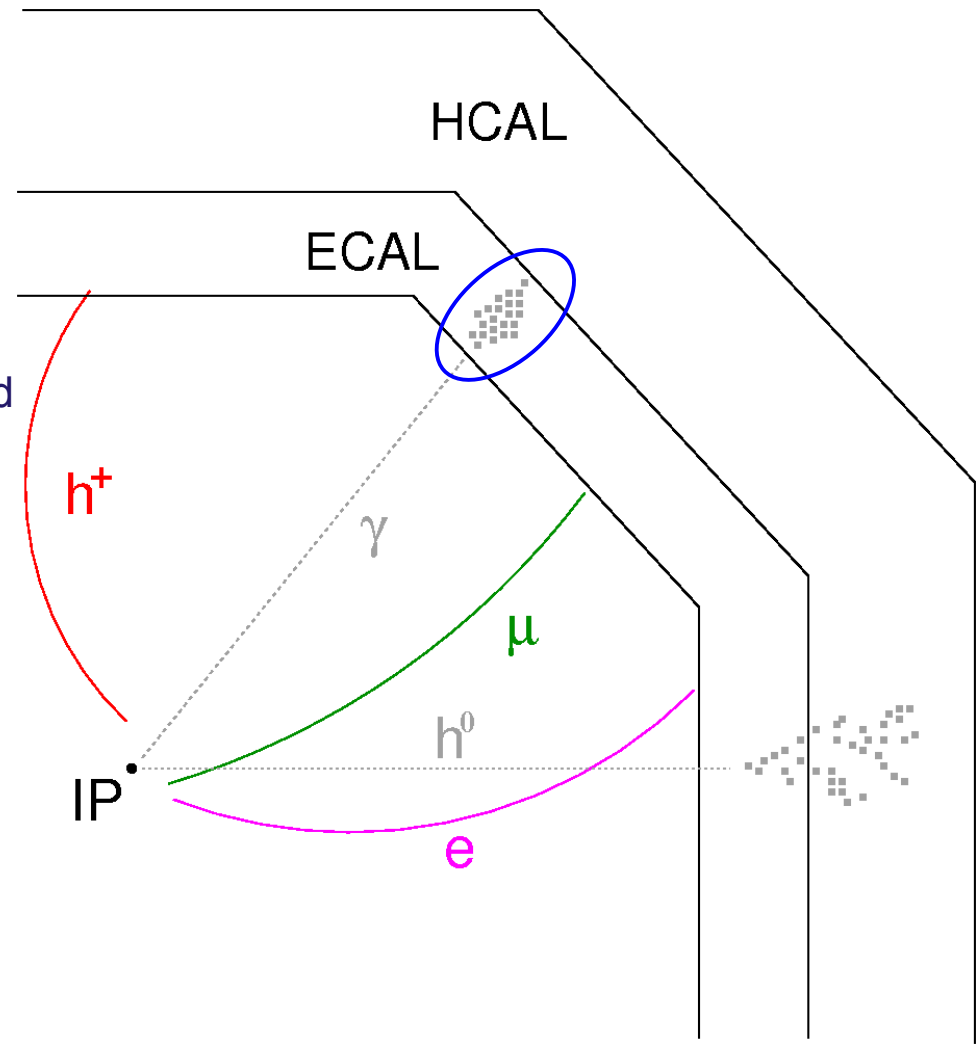
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6. particle ID for $e^{+/-}$, $h^{+/-}$



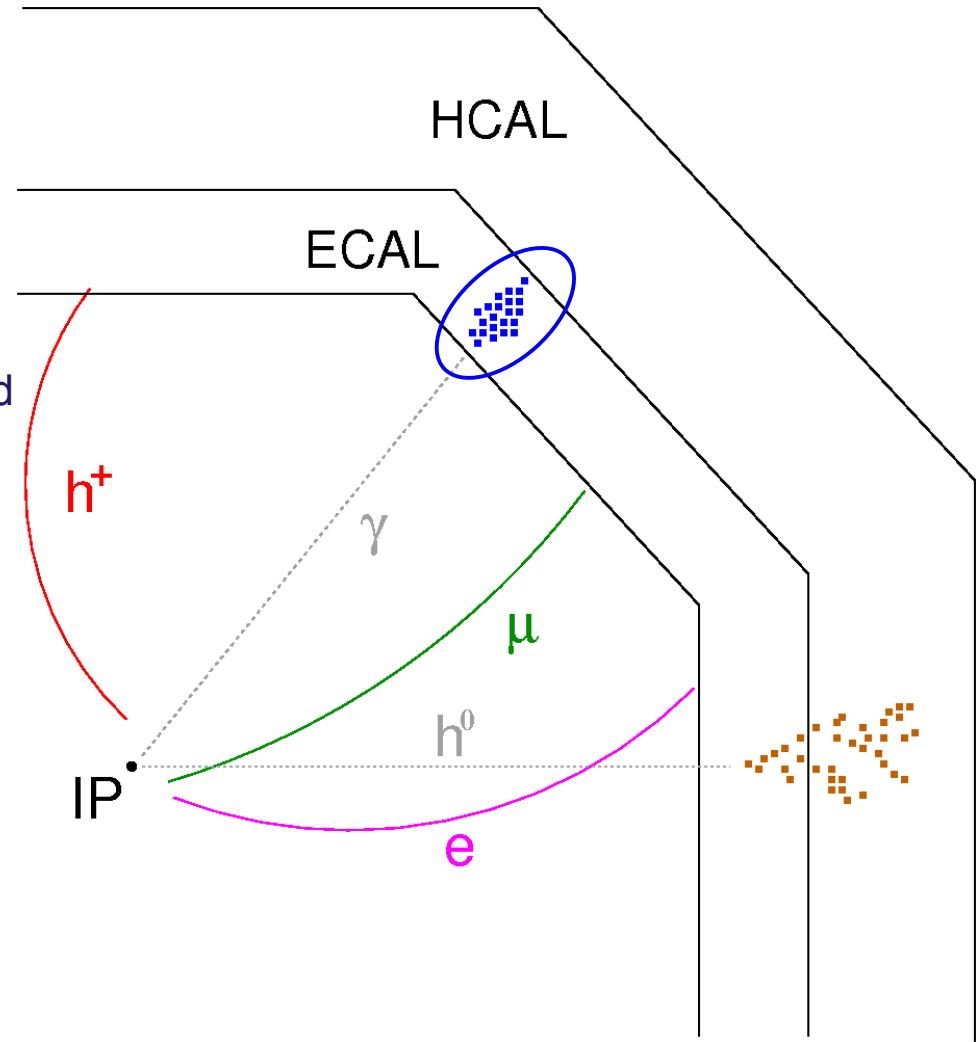
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6. particle ID for $e^{+/-}$, $h^{+/-}$
7. remove 'charged' Calorimeter hits



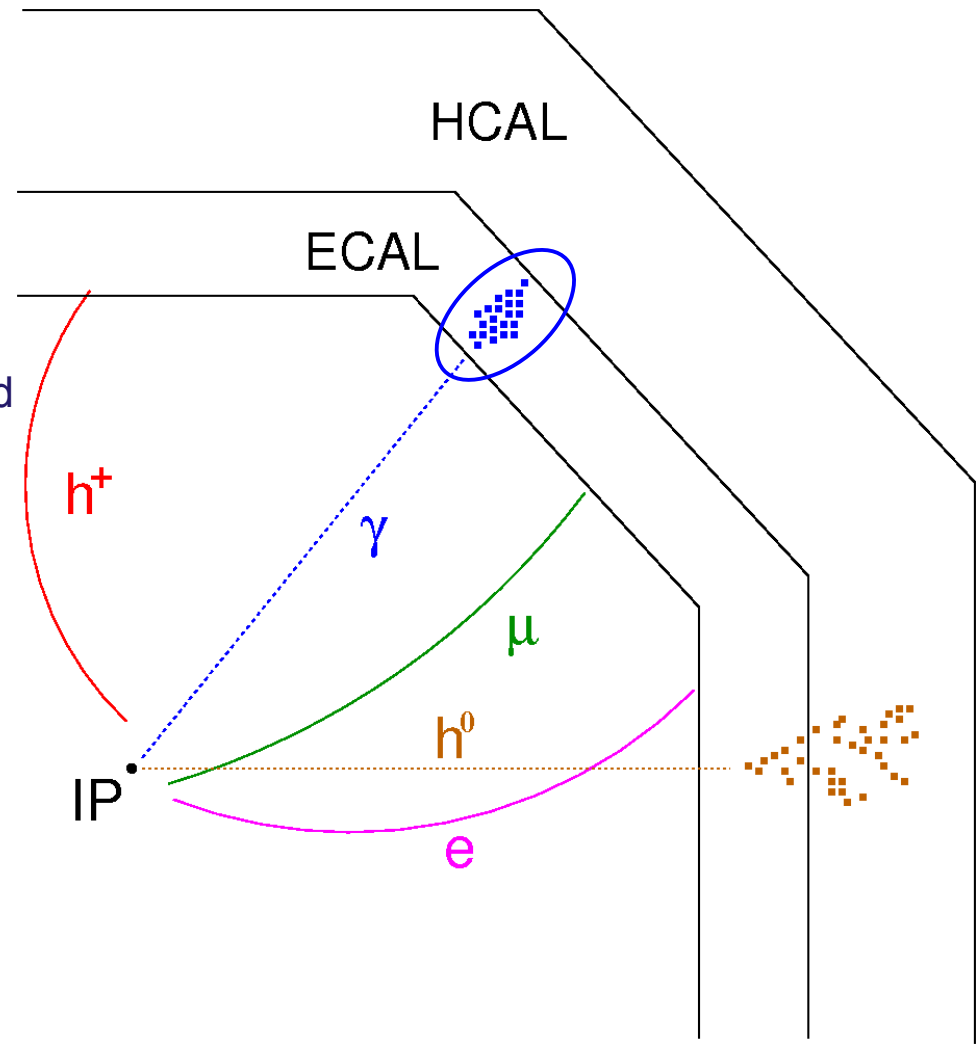
Track-Based Particle Flow Concept

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 - variable, depending on track and photon candidates
 - different algorithms
6. particle ID for $e^{+/-}$, $h^{+/-}$
7. remove 'charged' Calorimeter hits
8. clustering on 'neutral' hits



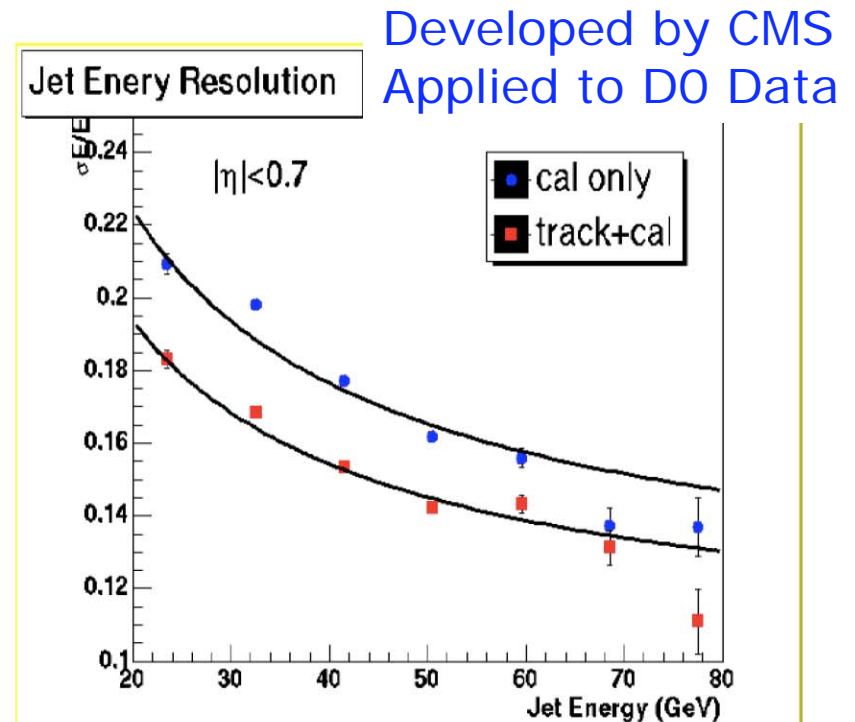
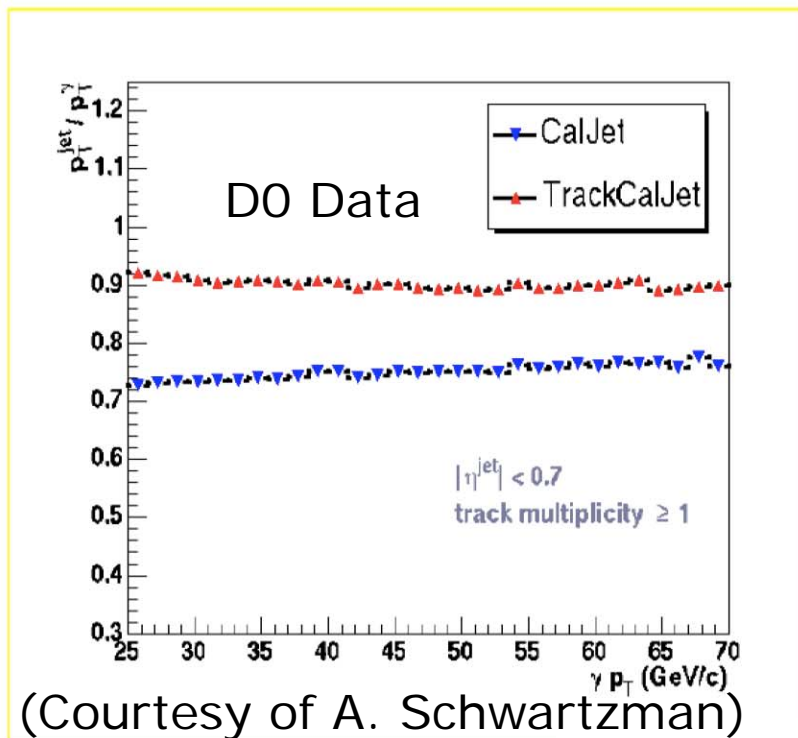
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7. remove 'charged' Calorimeter hits
8. clustering on 'neutral' hits
9. particle ID for photons and h^0



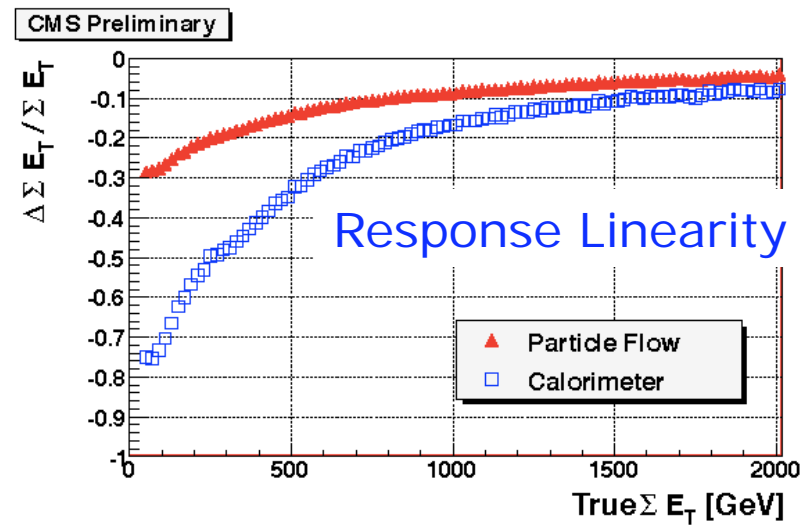
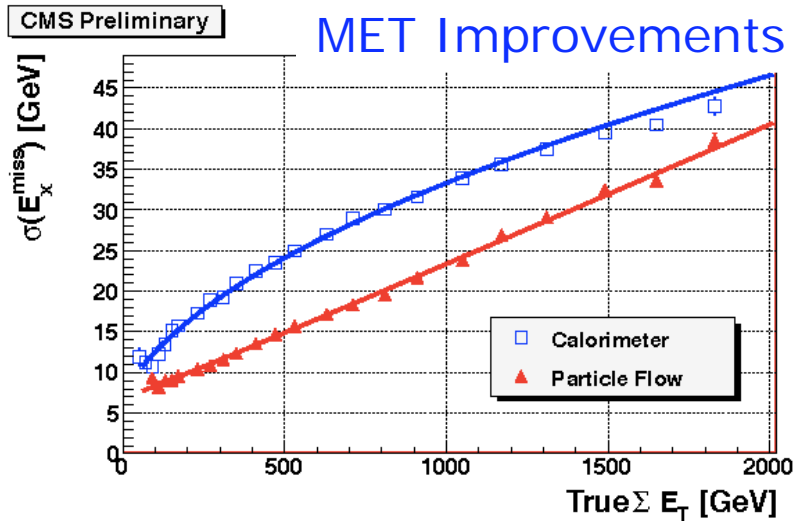
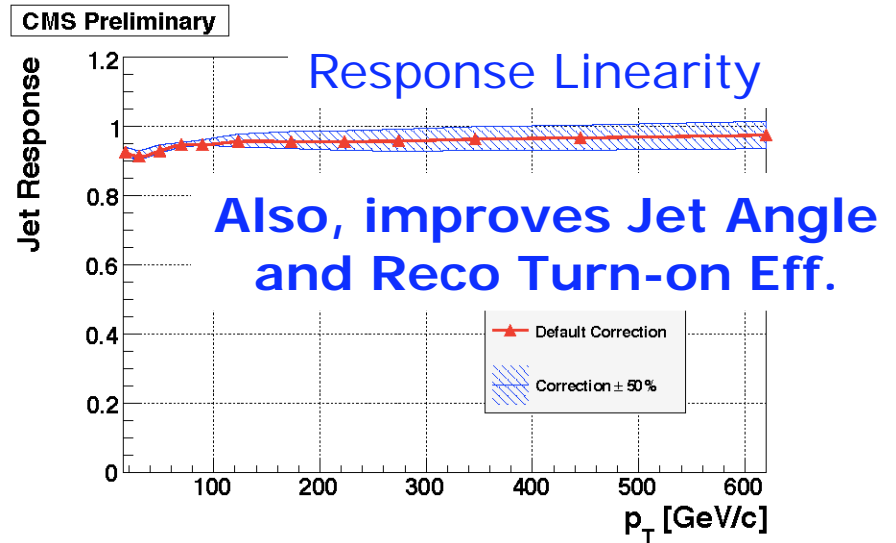
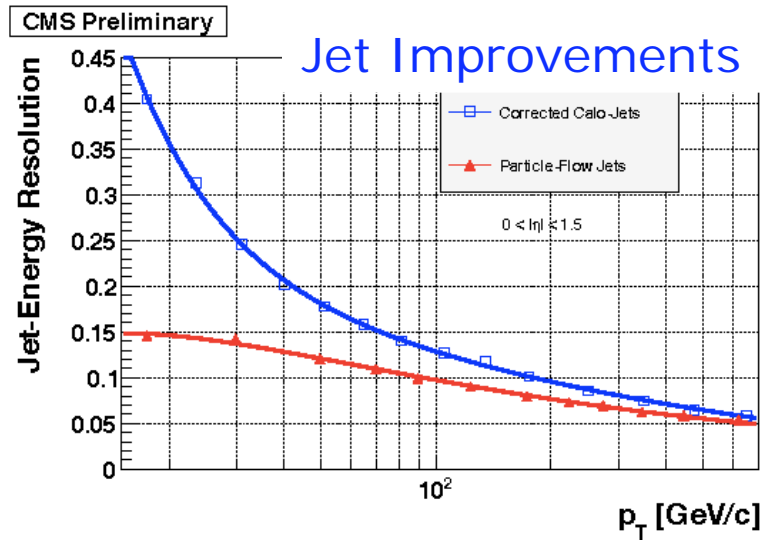
Jet Plus Track

- No track-cluster matching required
 - Expected e/π non-linearity is subtracted from calorimeter measurement and replaced with track momentum



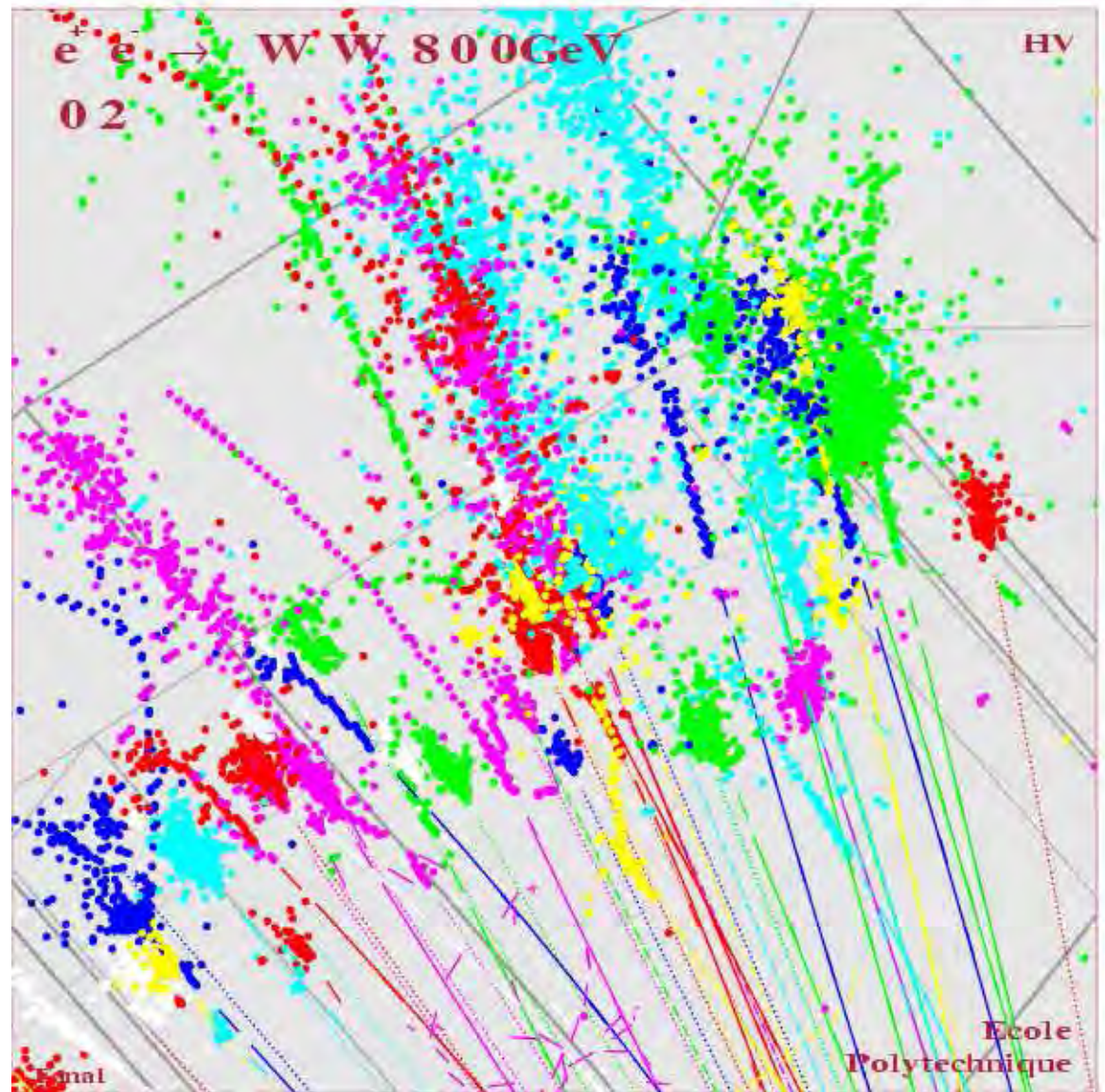
Magnitude of improvement is much larger
for calorimeters with large e/π non-linearity

Particle Flow



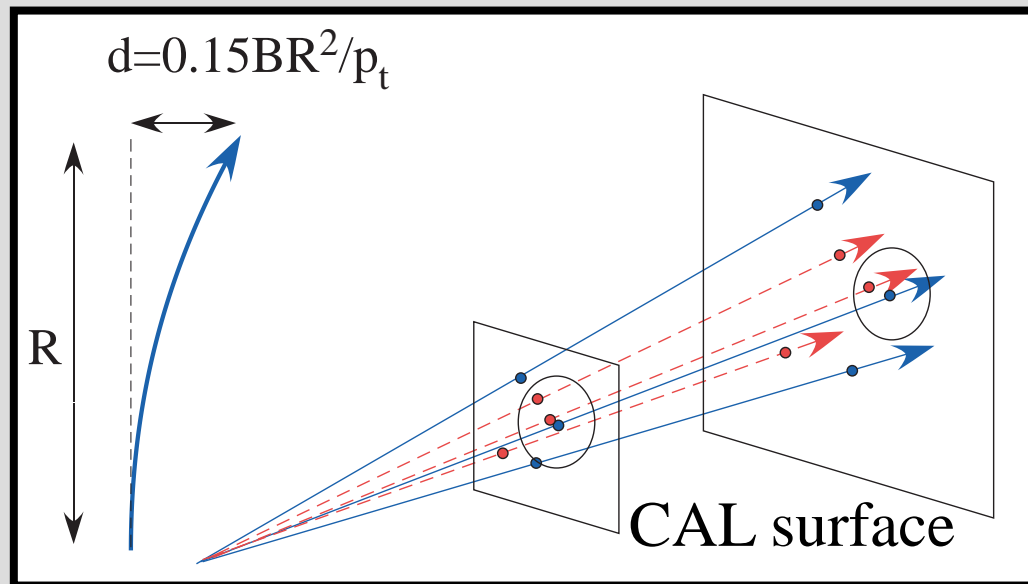
Main problem: Confusion

- At high energy jets are very narrow
- ⇒ Tracks are very close at the calorimeter
- Need very fine granularity of calorimeter and sophisticated software to separate showers
- Energy resolution still dominated by confusion term



Particle Flow Algorithm

- In order to get good energy resolution by PFA, separation of particles is important. → Reduce the density of charged and neutral particles at calorimeter surface.



Often quoted “Figure of Merit”

$$\frac{BR^2}{\sqrt{\sigma^2 + R_M^2}}$$

B : Magnetic field

R : CAL inner radius

σ : CAL granularity

R_M : Effective Moliere length

- For transverse separation of particles at the ECAL surface, stronger B-field and/or large ECAL radius are preferable.

* Fine segmentation of CAL is also important for pattern recognition.

Radius vs. B-field

