

Particles and Universe

Lecture 9 Perturbative calculation

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I. Theory of elementary particles – description of events; Perturbative calculation; Renormalization

II. Probability of processes at high energy – a need for Higgs particle

Description of events

- Relativistic and quantum effects. To describe creation and annihilation of particles and mixing as well as decays → formalism of

Quantum Field Theory, QFT

(applied first to the QED)

- Technique of QFT

- perturbative calculation:

- the lowest order of pert. series

- quantum corrections: trees and loops

- Feynman diagrams

Theory of elementary particles

- 1948 – new phase of quantum mechanics (QFT); precise measurements → a need of more precise calculations
- Feynman method: diagrams and Feynman's rules for calculation – today the universal tool of particle physics
first application in QED
- QED describes interaction of electron and photons. Difficulties since quantum corrections **infinite!**
Method of removing them → renormalization **procedure.**
- Weak interaction – even more difficult situation. Proposal: **new interaction and new particles** → EW interaction with gauge bosons W/Z with **Higgs boson** - **renormalizable!**

Nobel prize: Glashow, Salam, Weinberg 1979 (W/Z)
t'Hooft, Veltman 1999 (renormalizability)

Relativistic effects

- free and virtual particles

Relativistic effects

(special theory of relativity)

- Einstein postulates:
 - Invariance of laws with respect to change of reference system (for the inertial systems)
 - Constant velocity of light c in the inertial systems
- Lorentz transformation relates measurements of two inertial reference frames
- The most „visible” relativistic effects
 - lifetime for particles moving with velocity close to c longer (*time dilatation*)
 - Muons born in the upper atmosphere arrive on the surface of Earth only due this effect (660 m)*
 - similar effect for space (*length (Lorentz) contraction*)

Lorentz and Poincare invariance

- Invariance with respect to rotations (in space-time) → **Lorentz invariance**
If in addition invariance with respect to translation in space and time → **Poincare invariance**

Not everything is relative !!

If decay (process) is forbidden in some reference frame it is forbidden in any frame

- However – for a description of a particular process
→ a particular system maybe more useful.

Forbidden processes - examples

- For a *free* electron process $e \rightarrow e \gamma$ is forbidden
Why?

Electron at rest has the lowest possible energy ($E = mc^2$), and cannot have lower energy after emission of photon. If it is not possible in one frame then it is not possible in any

It is possible in matter

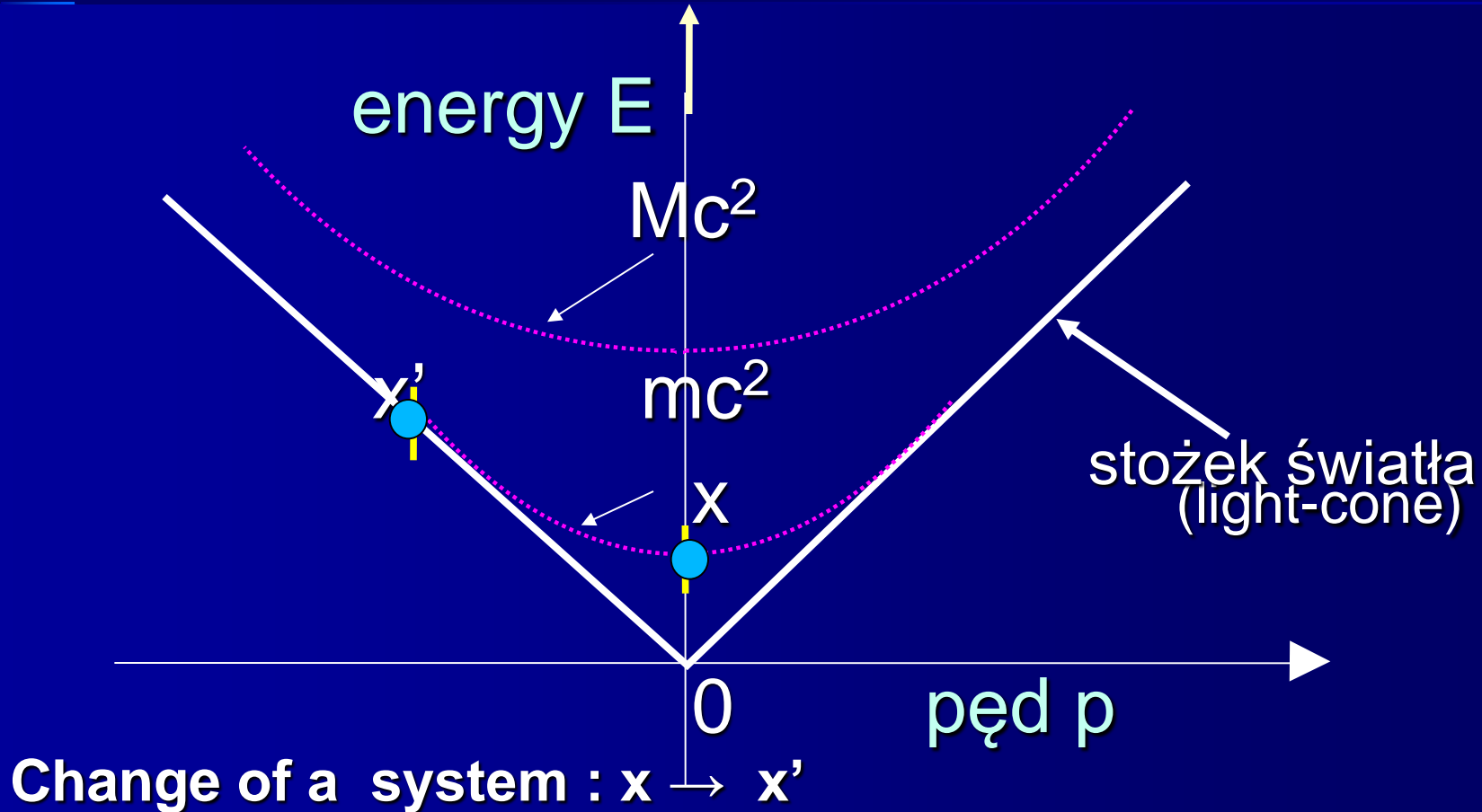
- For a *free* photon the decay $\gamma \rightarrow e^+e^-$ is not possible - observer moving towards a photon see it with a smaller momentum/energy ($E = pc$)...

Possible in the matter

Energy and momentum for a *free* particle

Relation defines the **mass shell** for a particle with mass m

$$E = c \sqrt{p^2 + m^2 c^2}$$



Particle on a mass shell = free particle

Inner particles in diagrams = virtual particles

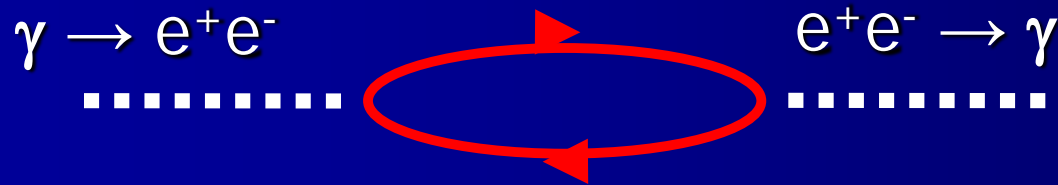
Quantum mechanics allows for

virtual particles for which

$$E^2 \neq p^2 + m^2 \quad (c=1)$$

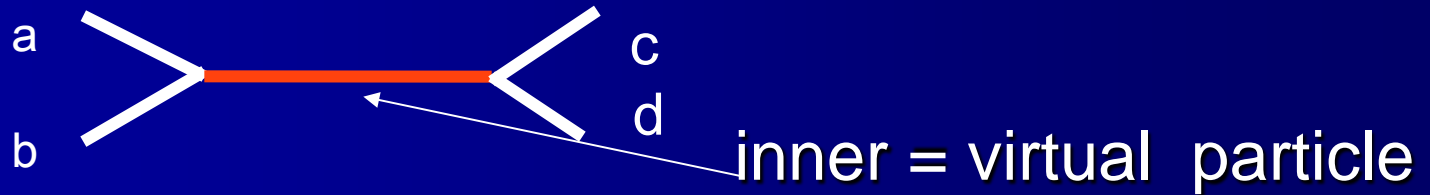
(particles *off* mass shell)

They exist for a while and are not observed directly
eg. they appear in the loops



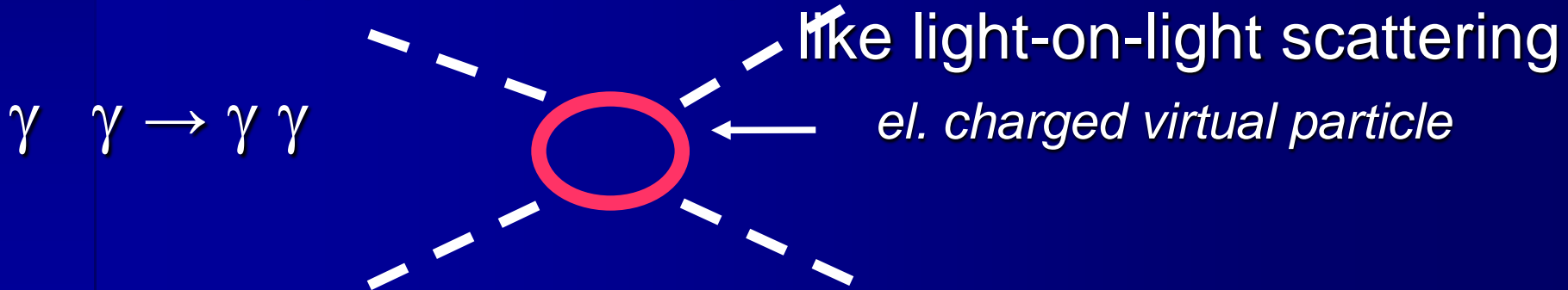
Examples- virtual particles

- In a collision of two particles a, b
two particles are produced c, d : $a+b \rightarrow c+d$



There are various channels
(= various virtual exchange) *like two split experiments*

- Some processes are possible due to them—



$$\gamma \gamma \rightarrow \gamma \gamma$$

photons are neutral and can not interact with themselves

Probability of processes

Probability amplitude and probability of processes

- Knowing properties of particles we can calculate probability of processes they are involved in
- Initial particles we produce or prepare, but the final particles we can only observe – but there are many ways from the initial to the final particles
(*channels of processes*)
- We calculate *probability amplitudes for each channel* → interference of the amplitudes
- Feynman diagrams for all channels of process → we **sum probability amplitudes not probabilities!**

$$|A|^2 = |\sum A_i|^2$$

Question which channel is like a question - which slit a photon used

Infinities

- Description of the process - all channels should be included, with possible virtual particles
- Virtual particles may have various energies, in loops even arbitrary energies and all these cases should be included.
- If contributions are not damped – a problem
 - > infinite probabilities for a given process!
- ... → problems with spin 1 particles, damping smaller for larger spin

(problem with graviton even worst)

QED:

infinities and renormalization

- Problems with photon (spin 1)? – In QED infinite contributions from various diagrams cancel - as was shown in 1948 by Feynman, Tomonaga, Schwinger (Nobel 1964) →
renormalization procedure
(introduced by Kramers in 1938)
- QED gives finite predictions - very precise, eg. for anomalous magnetic moment

Perturbative calculation

- Prediction for a process – we need to sum various contributions - with many vertices
- In principle -> series of infinite number of terms...
 - it is OK if next terms small

(small perturbation, perturbation series)

- Eg. in QED small coupling constant

$$\alpha = e^2/(4 \pi) \sim 1/137$$

so next term in perturbative expansion $\sim \alpha^2$ - 1%

- More terms in the perturbation series – including higher orders in the expansion in the coupling constant (higher powers of the coupling constant)
 - > higher precision of theoretical prediction

Anomalous magnetic moment for muon

Anomalous magnetic moment for muon (or $g-2|_{\mu}$)

The magnetic moment μ proportional to the spin s

spin and magn. moment

- vector quantities

$$\vec{\mu} = g_{\mu} \frac{e\hbar}{2m_{\mu}c} \vec{s} ; \quad g_{\mu} = 2 (1 + a_{\mu})$$

For the fundamental particle with spin $\frac{1}{2}$ the simplest act of the el-(magnetic) interaction $\rightarrow g = 2$ (Dirac 1925) so deviation (or $g-2$) is called the anomalous magnetic moment

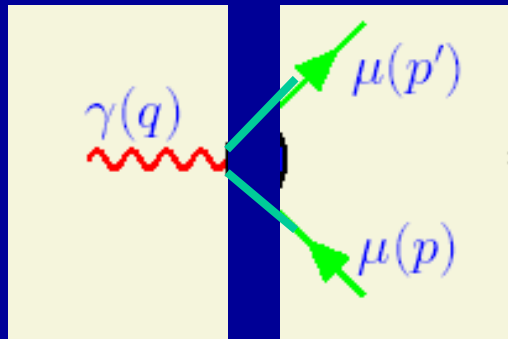
First measurement for electron in 1922! Next in 1948r

Stern, Gerlach 22: $g_e = 2$; **Kusch, Foley 48:** $g_e = 2 (1.00119 \pm 0.00005)$

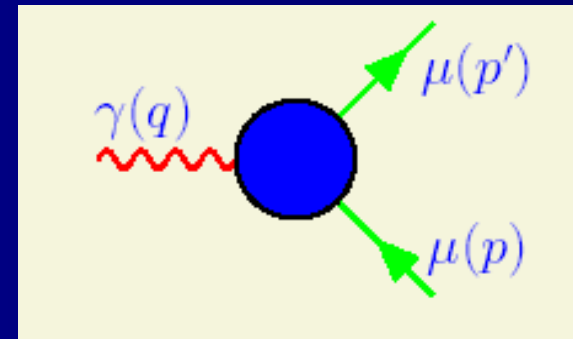
Origin of a_μ

(for muon)

From extra interactions ...



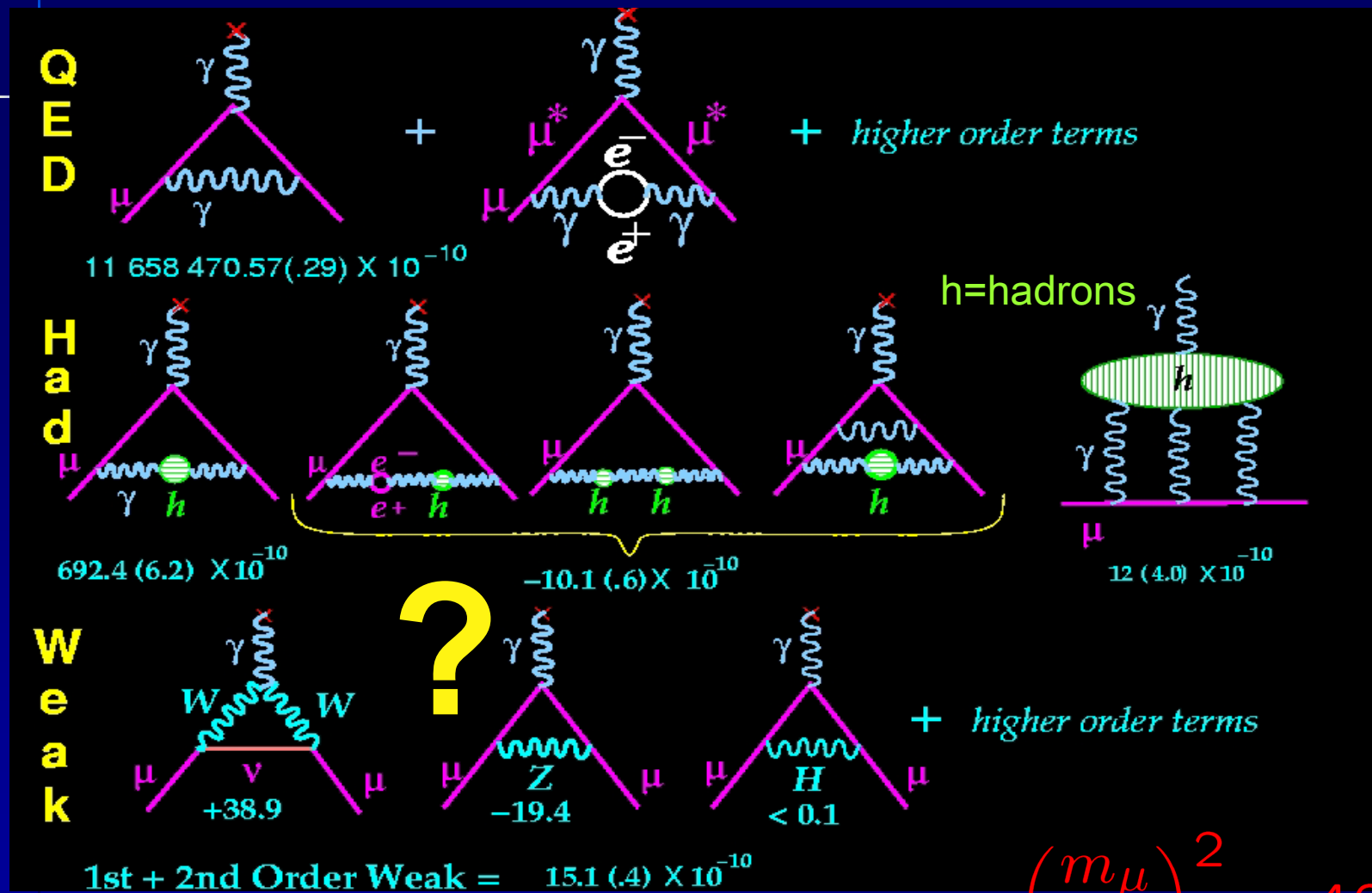
$g = 2$ if only this act



$a = g-2 \neq 0$ if more actions

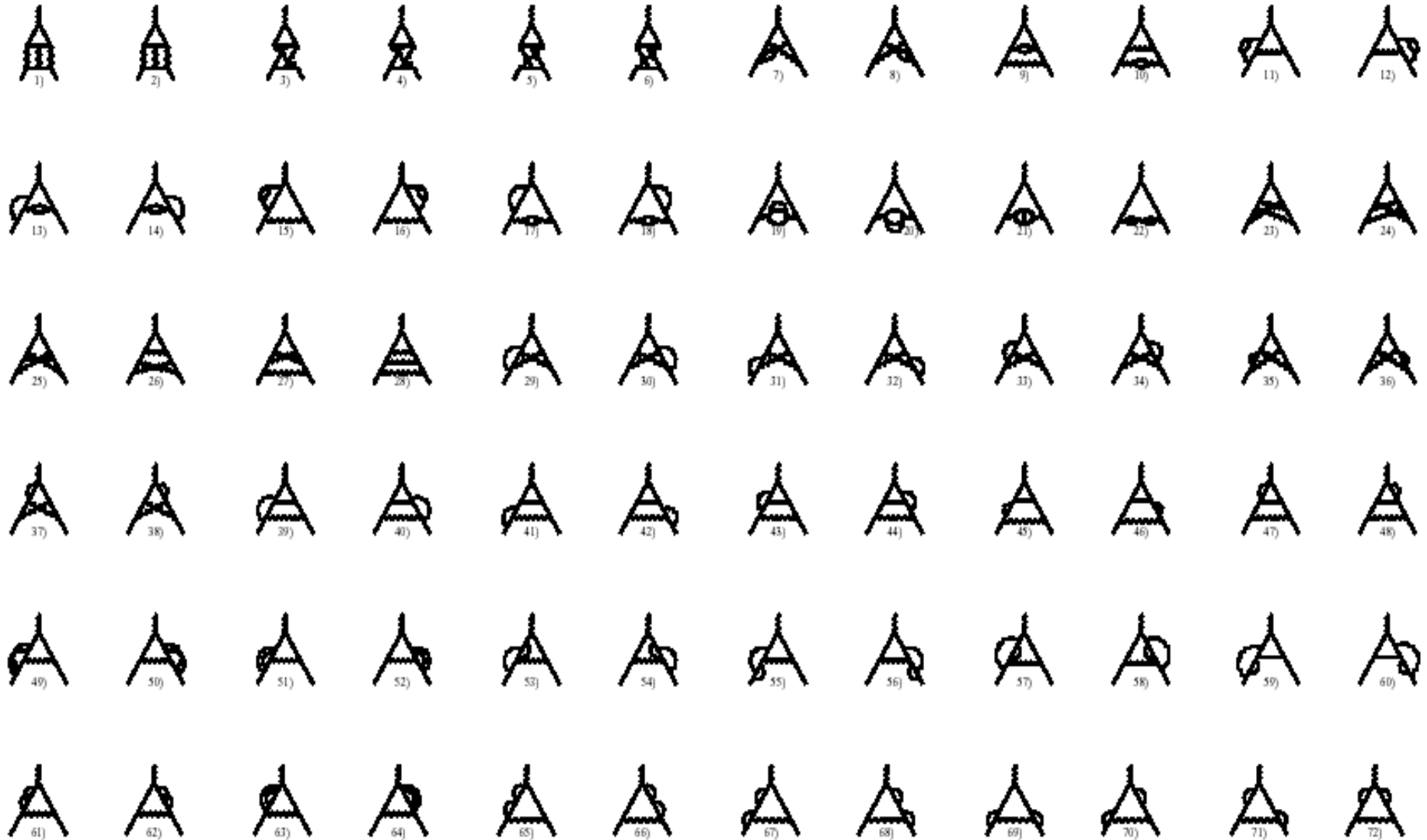
foton z pędem q , mion z pędem p , mion z pędem p' , $p'=p+q$
(dokładniej to są 4-pędy: (energia, 3-pęd))

Standard Model : QED, hadronic (h) and EW (W/Z i H) contributions



e vs. μ : relative contribution of heavier things $\left(\frac{m_\mu}{m_e}\right)^2 \simeq 40,000$

3- order of pert. calculation (QED)



QED Contribution a_μ^{QED}

$$\begin{aligned} a_\mu^{\text{QED}} \cdot 10^{10} = \sum C_i \left(\frac{\alpha}{\pi}\right)^i &= 11614097.3 \text{ (1-loop)} \\ &+ 41321.8 \text{ (2-loop)} \\ &+ 3014.2 \text{ (3-loop)} \\ &+ 38.1 \text{ (4-loop)} \\ &+ 0.4 \text{ (5-loop)} \end{aligned}$$

Terms up to α^3 are known analytically, a recent more accurate numerical calculation of the α^4 terms and the leading $\log \alpha^5$ terms gave

(T. Kinoshita and M. Nio, 2005; A.L. Kataev, 2006):

$$a_\mu^{\text{QED}} = (116584719.4 \pm 1.4) \cdot 10^{-11}.$$

From the latest value of a_e (G. Gabrielse et al., 2006; M. Passera, 2006):

$$\alpha^{-1} = 137.035999710(96), \quad a_\mu^{\text{QED}} = (116584718.09 \pm 0.14 \pm 0.08) \cdot 10^{-11}.$$

The errors are due to: a/ $\mathcal{O}(\alpha^5)$, b/ α

Jegerlehner'07

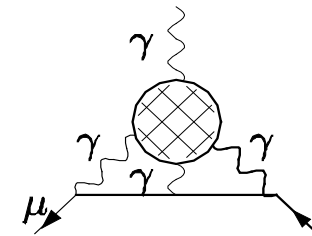
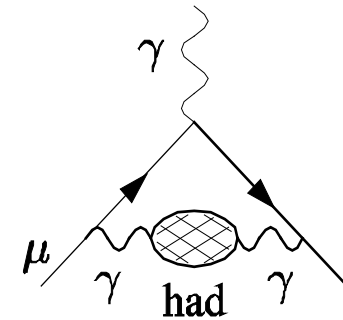
$g-2|_{\mu}$: exp-theory (MS)

EXP

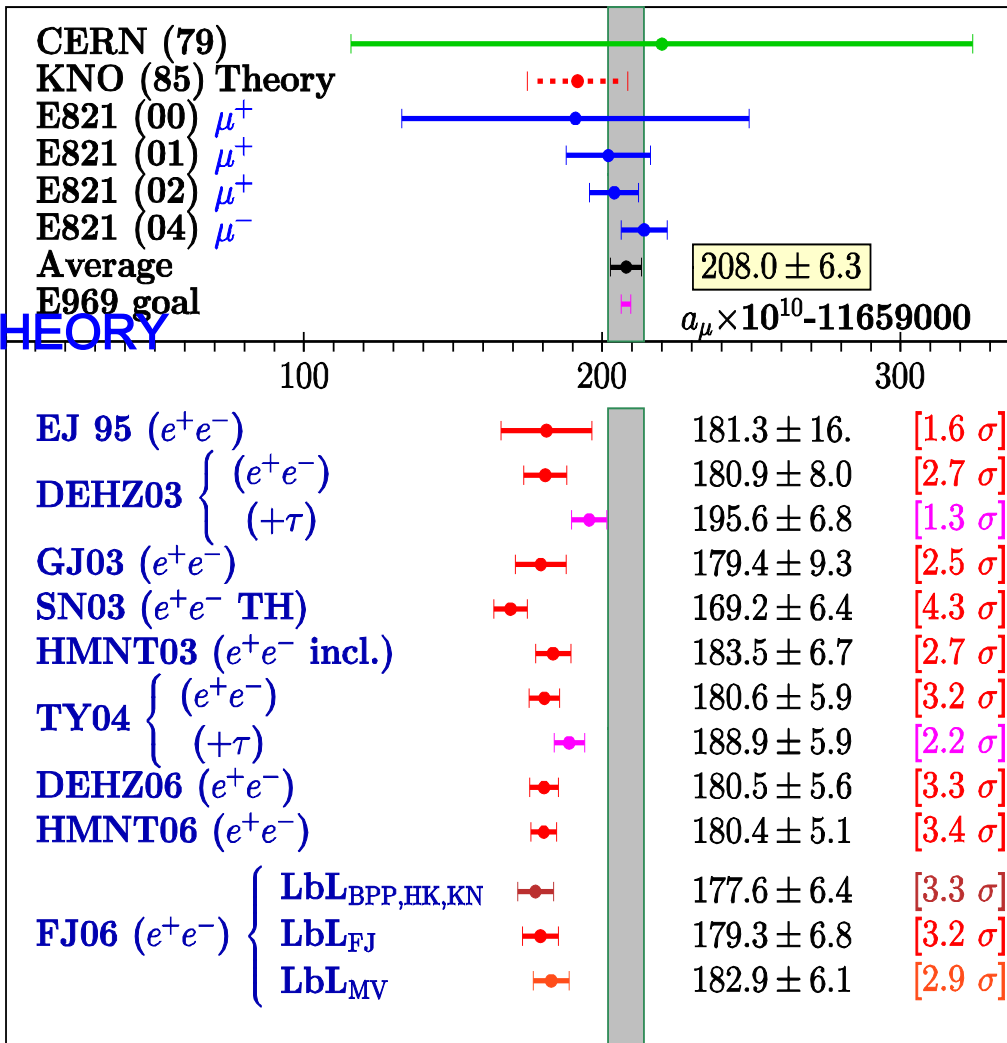
New Physics?

$$\delta a_{\mu} = (287 \pm 91) 10^{-11}$$

(3.2 σ)



THEORY



New calculation in SM (2013)

Dermisek, Rava – May 2013

The discrepancy between the measured value of the muon anomalous magnetic moment [17] and the SM prediction,

3.4 σ

$$\Delta a_{\mu}^{exp} = a_{\mu}^{exp} - a_{\mu}^{SM} = 2.7 \pm 0.80 \times 10^{-9}, \quad (28)$$

Experiment Brookhaven 2006

[17] G. W. Bennett *et al.* [Muon G-2 Collaboration], “Final Report of the Muon E821 Anomalous Magnetic Moment Measurement at BNL,” *Phys. Rev. D* **73**, 072003 (2006) [hep-ex/0602035].

The most precise quantity...

Brookhaven,
USA

In 2014

transport to
Fermilab
(Chicago)

<http://muon-g-2.fnal.gov/bigmove/>

The BNL muon storage ring

Pomiar g-2
dla mionu



www.g-2.bnl.gov

Small, very precise experiment at low energies - www.g-2.bnl.gov

Renormalization

Idea

- QED – infinities only in expressions containing mass and electric charge. Using quantities measured in experiment we can remove infinities as follows:

A – expression for physical quantity; let 1st order correction Δ contains an infinite contribution Δ' ,

$$A = e_0(1 + \Delta) + \dots = e_0(1 + \Delta' + \dots)(1 + \Delta'' + \dots) = e_{\text{fiz}}(1 + \Delta'' + \dots),$$

so the prediction for quantity A finite (Δ'' - finite)
(+...higher terms)

For magnetic moments infinities are at the intermediate stages of calculations – but they are the same for electron and muon.

We can express them by each other and for electric charge use the exp. value.

Trick very useful

- Electric charge is a *free (initial)* parameter of the QED; the same for the electron mass m_e
- Theory does not predict values of these parameters – we can hide infinities in them
- If infinities only in expressions for free parameters of the theory → **theory is renormalizable and provide definite predictions for physical processes**
- **It is not satisfactory, but works OK**

Non-renormalizability

- There are theories with infinities not only in free parameters
- For long time it was considered that theories with spin-1 particles, but different than a photon, are non-renormalizable (eg. Fermi theory for weak int.).
- Today we know that theories with spin-1 particles are renormalizable if they are **gauge theories**, (moreover – non-abelian **Yang-Mills theories**)

Gravitation: gauge theory, but not renormalizable

Gauge transformation- a reminder

Schrodinger equation, based on $E = (\bar{p}^2/2m + V)$

- $i \partial \psi(\bar{\mathbf{x}},t) / \partial t = - 1/2m \bar{\nabla}^2 \psi(\bar{\mathbf{x}},t)$
(for potential $V=0$ and using $\hbar = 1$)

- $E \rightarrow i \partial_0$ $\partial^\mu = (\partial_0, - \bar{\nabla})$
 $\bar{p} \rightarrow - i \bar{\nabla}$

- local transformation, phase $\alpha(\bar{\mathbf{x}},t)$
 $\psi'(\bar{\mathbf{x}},t) = e^{i \alpha(\bar{\mathbf{x}},t)} \psi(\bar{\mathbf{x}},t)$

$|\psi(\bar{\mathbf{x}},t)|^2$ probability does not change, but

invariance of equation only if we add interaction !

Predictions for larger energies

Probability of processes for large energies

- Bad, if probability of processes rises with energy (bigger than 100 % !?)
- QED: Compton scattering – at the lowest order of pert. calculation two diagrams, each rising with energy but the sum of diagrams – OK (cancelation)
- Cancelation results from the structure of the theory (gauge symmetry)

Calculation of probability using Feynman rules

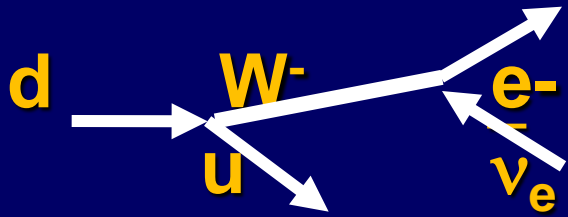
- To each line and vertices in Feynman diagram a factor is assigned. Here we track only the energy E .
 - Incoming or outgoing photon (and each **spin 1 particle**) – a factor E
Virtual photon (spin 1 particle) – a factor $1/E^2$
 - Incoming or outgoing **spin 1/2 particle** - a factor \sqrt{E} ,
virtual spin 1/2 particle – a factor $1/E$
 - Incoming or outgoing **spin 0 particle** – a factor 1,
virtual spin 0 particle – a factor $1/E^2$
 - Additional factors from couplings

Product of factors → probability amplitude A

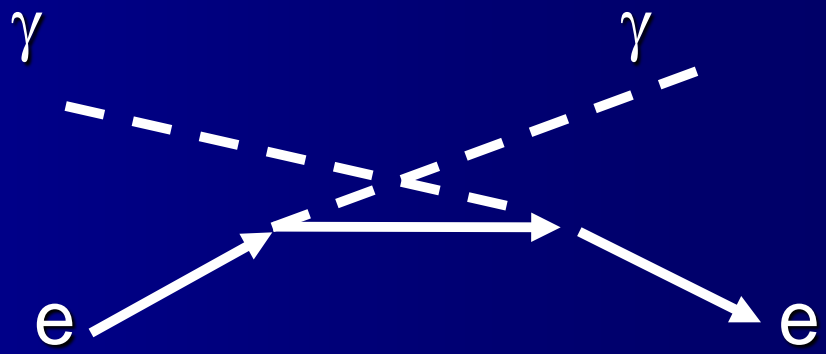
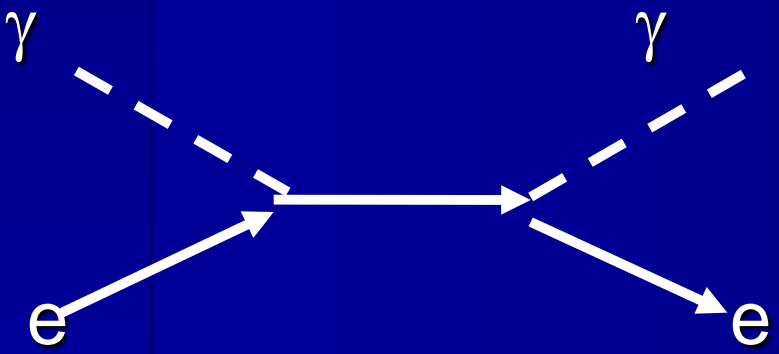
(probability = $|A|^2$)

Electroweak interaction

Decay $d \rightarrow u e \bar{\nu}_e$



- W boson is very massive (80.4 GeV), so in decay of quark d (mass \sim MeV) it is very virtual (far off mass-shell)
- Boson W has spin 1 \rightarrow problem with renormalizability
- To understand this problem, we analyse first Compton scattering $\gamma e \rightarrow \gamma e$ (QED), two diagrams:



Calculation of the probability amplitude for the Compton process

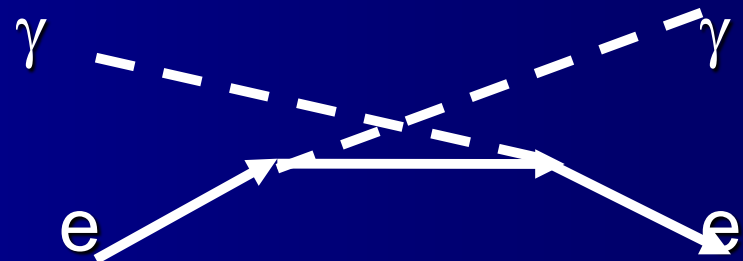
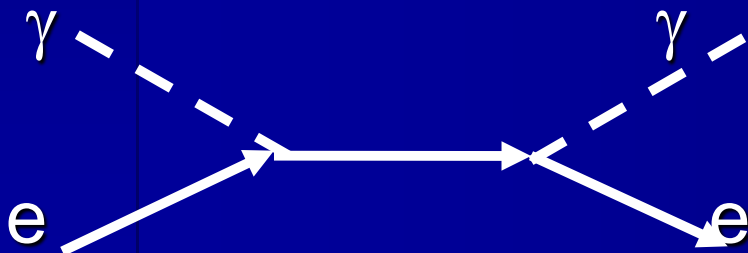
For this process the amplitude rises with the energy as

$$E^2 (\sqrt{E})^2 1/E = E^2$$

(for probability E^4)

Bad behaviour for the individual diagrams, but in sum cancelation and the final amplitude does not grow with energy

$\gamma e \rightarrow \gamma e$

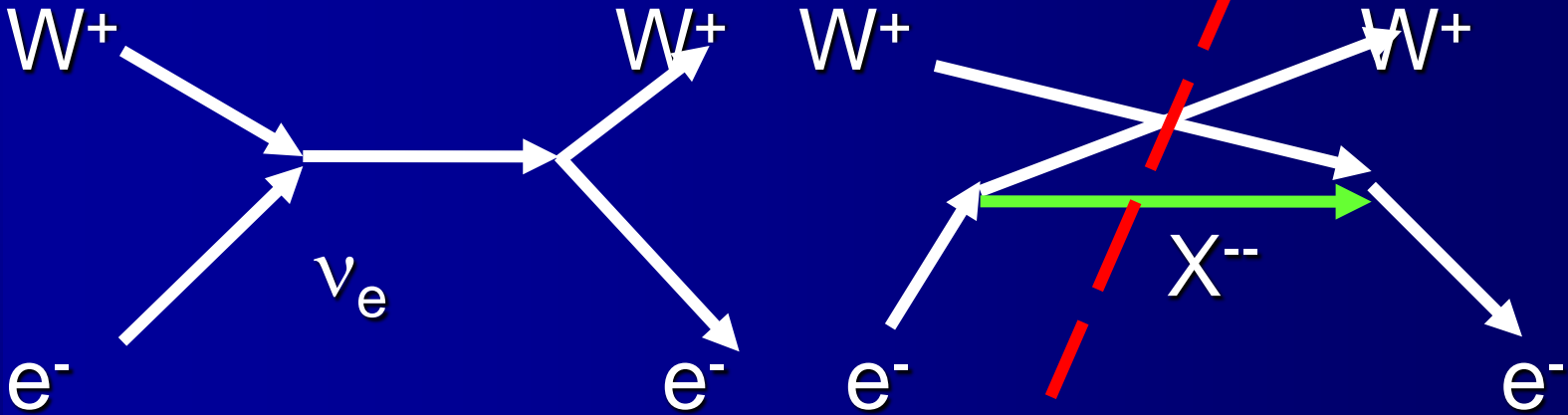


Scattering $W e \rightarrow eW$



Tu strzałki na liniach W
oznaczają pędy

■ Exchange γ by W:

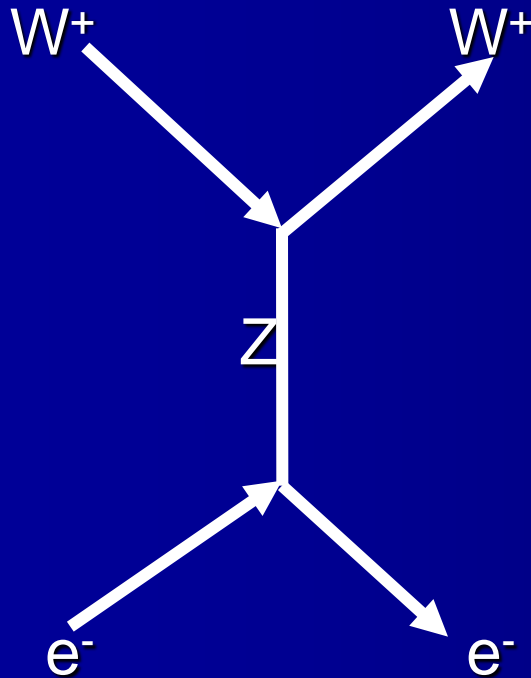


Charge conservation \rightarrow X^{--} double charged
particle – not observed !

Bad high energy behaviour $E^2 (\sqrt{E})^2 1/E = E^2$

Z boson needed!

- Using diagram with Z boson



Behaviour $E^2 (\sqrt{E})^2 (\text{vertex } WWZ \sim E)/E^2 = E^2$, and a proper sign of coupling WWZ – cancelation like for the Compton process!

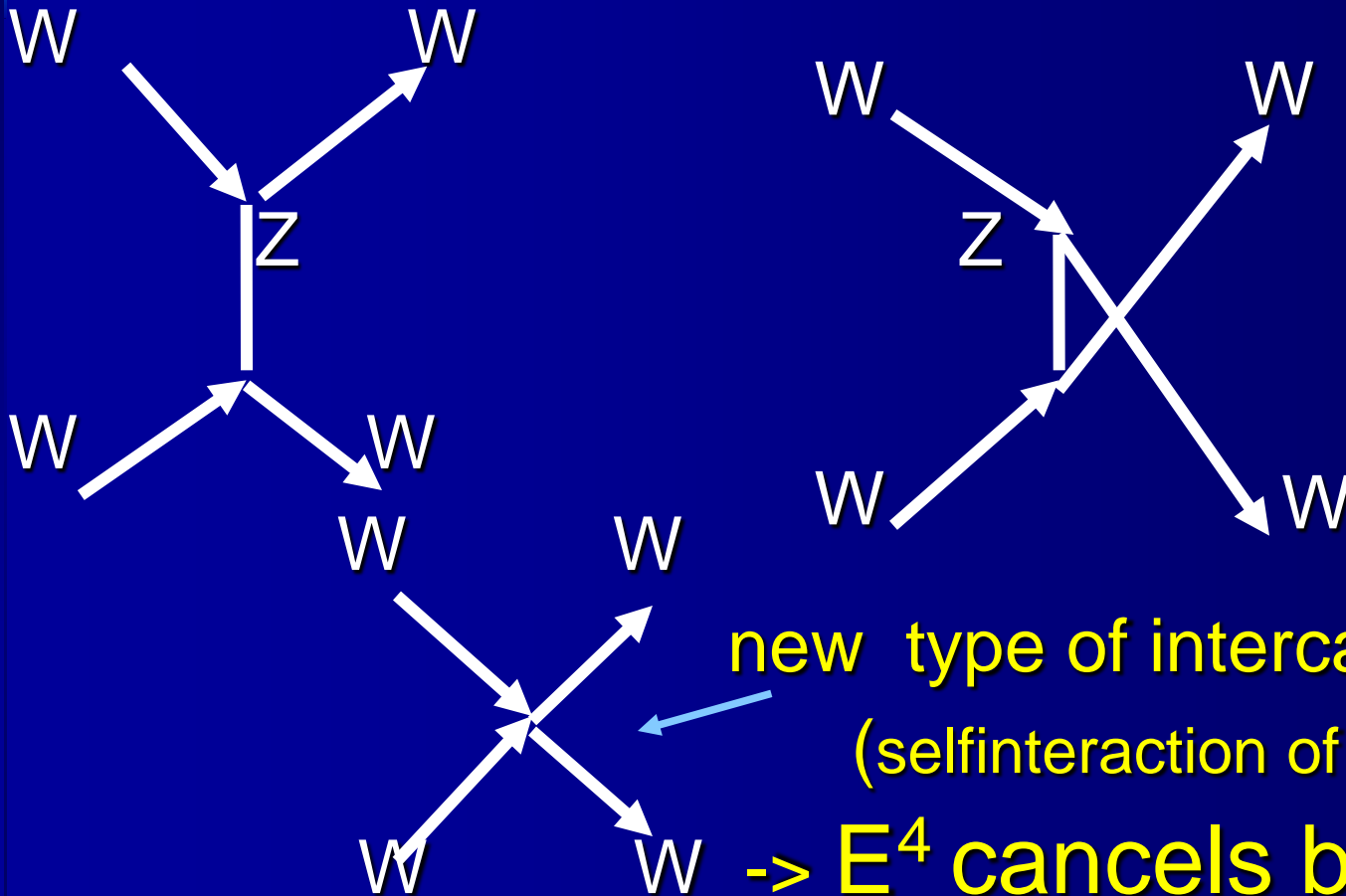
**Price for a good behaviour for E:
it must exist neutral particle Z with
a define interaction (coupling) !**

It was found – this is a Z boson !

→ A success of theory

Scattering $WW \rightarrow WW$

$\sim E^4$ (since $E^4 E^2/E^2$) - even worst...



new type of interaction
(selfinteraction of W)

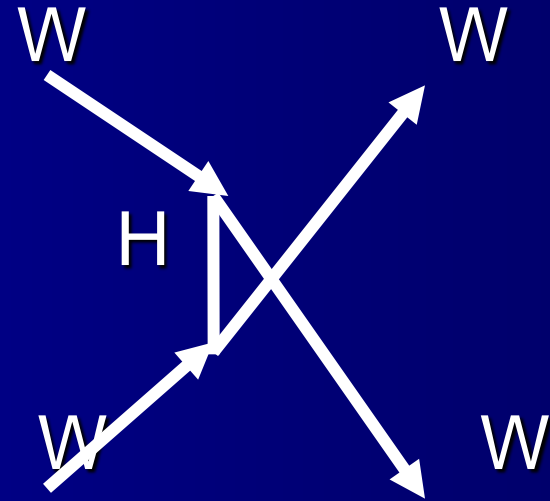
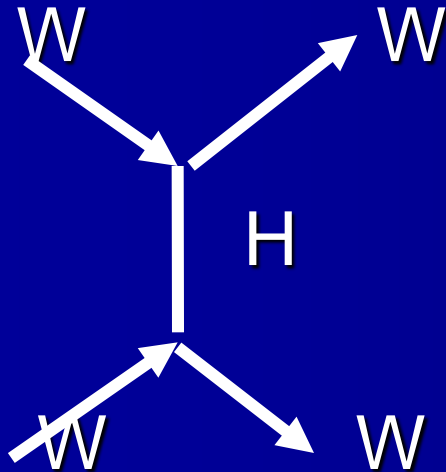
$\rightarrow E^4$ cancels but not E^2

E^2 term, in fact $E^2 M^2$

- In the amplitude - wrong term $E^2 M^2$,
where M - mass of W or Z
- New contribution needed to cancel this bad term \rightarrow the simplest diagram with exchange of spin-0 particle, which couples to W/Z proportionally to mass

Higgs boson needed !

Scattering of W on W: H contribution



Couplings proportional to mass of particles to which H couples, and as a result →
a good high energy behaviour $WW \rightarrow WW$!

Mass generation and consistent description of processes at high energies

- Good high energy behaviour - if H exists and couples to W/Z proportionally to their masses
- Higgs particle is related to the mass generation in SSB

→ properties of H particle related to the mass generation and those needed for a good high energy behaviour $WW \rightarrow WW$

- THE SAME !!!

- Expected relation to gravity (mass..)

The theory ends here - Veltman

Facts and Mysteries in Elementary Particle Physics, 2003

- *„The theory ends here. We need help. Experiments must clear up this mess.”*
→ LHC pp, ILC e^+e^- (PLC $\gamma\gamma$, $e\gamma$)

July 2012 :

LHC - Higgs particle with mass 125 GeV

Questions to lecture 9

- Does fact of decaying of a particle depend on the reference frame?
- Do *a free particle* and *a particle on the mass shell* mean the same?
- Why a free positron can not emit a photon and remain on the mass shell?
- What does it mean – *a virtual particle*?
- Does QED have infinite expressions at intermediate stages of calculation?
- Compare the one- and two –loop QED corrections for the anomalous magnetic moments for muon?
- What is a precision we know (in 2013) the anomalous magnetic moments for muon (experiment minus theory (SM))?
- When the theory is renormalizable?
- Proper high energy behaviour of the probability of EW processes demands existence of spin 0 particle. How such a particle couples to gauge boson Z?
- Is the Higgs boson discovered?