Statistical analysis of experimental data Markov Chains

Aleksander Filip Żarnecki



OF WARSAW

Lecture 14 January 23, 2025

Statictical analysis 14



## Markov Chains

2 Markov Chain Monte Carlo

3 Application to parameter fitting



## Applications

- **F**w

(lecture 05)

Described procedure can be used not only to calculate integrals of one-dimensional functions, it is much more general... How to calculate volume of a given shape?

Standard procedure:

scan all dimensions using dense point grid and sum cells with centers inside the volume



#### Monte Carlo integration:

Generate random points in the considered space and count points inside the volume

## **General case**

Examples presented considered the special case: input random variables had uniform distribution and "test function" was binary (returning 0 or 1).

In the general case we want to determine an expectation value of a function  $h(\mathbf{x})$  of random variable vector  $\mathbf{x}$  described by  $f(\mathbf{x})$  pdf:

$$\mu_h \equiv \mathbb{E}_f[h(\mathbf{x})] = \int d\mathbf{x} h(\mathbf{x}) f(\mathbf{x})$$

Monte Carlo determination of  $\mu_h$  assumes we can generate random variables from  $f(\mathbf{x})$ . We can then calculate:

$$\mu_{MC} = \lim_{N \to \infty} \frac{1}{N} \sum_{i} h(\mathbf{x}_{i})$$

where  $\mathbf{x}_i$ , i = 1, ..., N are random (input) variables generated from  $f(\mathbf{x})$ 





## Weighted Monte Carlo

General method for generating random points in multi-dimensional space using acceptance-rejection technique can have very low efficiency, if probability distribution function  $f(\mathbf{x})$  varies a lot, eg. has sharp peaks.

Assume we know how to generate random numbers from g(x). We can then apply the following procedure:

- generate  $\mathbf{x}_i$  distributed according to  $g(\mathbf{x})$
- accept all generated value  $\mathbf{x}_i$ ,

but consider them with additional weight:  $w_i = f(\mathbf{x})/g(\mathbf{x})$ 

For example, when calculating the expectation value of  $h(\mathbf{x})$ :

$$\mu_{MC} \rightarrow \mu_{wMC} = \frac{\sum_{i} w_{i} h(\mathbf{x}_{i})}{\sum_{i} w_{i}}$$

"unweighted" samples considered previously correspond to  $w_i \equiv 1$ 

## Weighted Monte Carlo



When using weighted Monte Carlo "events", number of events has to be replaced by sum of weights:

$$V \rightarrow N_w = \sum_i w_i$$

Variance of the sum of weights:

$$\mathbb{V}(N_w) = \sum_i w_i^2$$

Statistical power of the weighted Monte Carlo sample is equivalent to unweighted sample of:

$$N_{eq} = \frac{N_w^2}{\mathbb{V}(N_w)} = \frac{(\sum_i w_i)^2}{\sum_i w_i^2}$$

For Poisson distributed random number  $\mathbb{V}(N) = N$ 

## **General problem**



Presented above was a simple example of a more general problem: how to estimate parameters of the probability distribution function from the results of the experiment (measurements).

In many cases, parameter value can not be directly extracted from the measurement results.

In the general case, shape of the probability density function for measurement result  $\mathbf{x}$ :

$$\mathbf{x} = (x_1, \ldots, x_n)$$

depends on a set of pdf parameters:

$$\boldsymbol{\lambda} = (\lambda_1, \ldots, \lambda_p)$$

so the probability density should be written as:

 $f(\mathbf{x}; \boldsymbol{\lambda})$ 



## Maximum Likelihood Method

The product:

$$L = \prod_{j=1}^{N} f(\mathbf{x}^{(j)}; \boldsymbol{\lambda})$$

is called a likelihood function.

The most commonly used approach to parameter estimation is the maximum likelihood approach: as the best estimate of the parameter set  $\lambda$  we choose the parameter values for which the likelihood function has a (global) maximum.

Frequently used is also log-likelihood function

$$\ell = \ln L = \sum_{j=1}^{N} \ln f(\mathbf{x}^{(j)}; \boldsymbol{\lambda})$$

we can look for maximum value of  $\ell$  or minimum of  $-2~\ell~=~-2\ln L$ 





2 Markov Chain Monte Carlo

3 Application to parameter fitting

4 Final exam

### **General concept**



## (Bonamente)

Markov Chain is a stochastic process where we consider the sequence of measurements (random variables)  $X^{(t)}$ . Measurements at fixed time intervals are a frequent case...

Outcome of the measurement (also called "state" of the chain) has to belong to the defined "state space". It is our sample space...

However, the probability density for different states is not given a priori! Instead, probability of the subsequent state (measurement at t + 1) depends only on the current state of the system:

 $P(X^{(t+1)}) = P(X^{(t+1)}|x^{(t)})$ 

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Probability can change in time, but it depends only on the current state of the chain, and not on any state of its earlier history!

This "short memory" property is known as the "Markovian property". As for particle decays!



Consider two boxes with a total of N balls.

The state of the system can be defined by a number *n* of balls which are placed in the first box,  $0 \le n \le N$ . The state space of the system has N + 1 elements.



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- select a box at random,
- move one ball from the selected box (if not empty) in the other one.



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This chain can be presented in terms of the transition probabilities:

$$p(n^{(t+1)}) = \begin{cases} \frac{1}{2} & \text{for } n^{(t+1)} = n^{(t)} \pm 1 \text{ and } n^{(t)} \neq 0 \text{ and } n^{(t)} \neq N \\ 1 & n^{(t+1)} = n^{(t)} \pm 1 \text{ and } (n^{(t)} = 0 \text{ or } n^{(t)} = N) \\ 0 & n^{(t+1)} \neq n^{(t)} \pm 1 \end{cases}$$



14\_Simple.ipynb COPPEN IN COLAB

Result of the algorithm implementation





14\_Simple.ipynb

Result of the algorithm implementation



Looks like symmetry violation?...





Result of the algorithm implementation



Large time scales for count fluctuations  $\Rightarrow$  symmetry restored on longer time scales



14\_Simple.ipynb COPPEN IN COLAB

Result of the algorithm implementation



Fluctuations still visible in ball count distribution...

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Result of the algorithm implementation



Decrease with the chain lenght...





Result of the algorithm implementation



Decrease with the chain lenght...

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14\_Simple.ipynb COPPEN IN COLAB

Result for the extended example (4 boxes)



Even starting from even ball distribution, large fluctuations appear very soon...

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14\_Simple.ipynb COPPEN IN COLAB

Result for the extended example (4 boxes)







Result for the extended example (4 boxes)



#### **Ehrenfest chain**



Simple model of diffusion: same case of two boxes with a total of N balls, but different procedure for generating next step.

The state of the system is defined by a number *n* of balls the first box,  $0 \le n \le N$ .



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$$p(n^{(t+1)}) = \begin{cases} \frac{n^{(t)}}{N} & \text{for } n^{(t+1)} = n^{(t)} - 1\\ \frac{N - n^{(t)}}{N} & n^{(t+1)} = n^{(t)} + 1\\ 0 & n^{(t+1)} \neq n^{(t)} \pm 1 \end{cases}$$



(Bonamente)

previously we were selecting a box



#### **Ehrenfest chain**

14\_Ehrenfest.ipynb

#### Result of the algorithm implementation



Looks again like symmetry violation?...



### Ehrenfest chain

14\_Ehrenfest.ipynb

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Symmetry restored on longer time scales...



## Simple example: Ehrenfest chain

14\_Ehrenfest.ipynb

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Decrease with the chain lenght...



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Decrease with the chain lenght...



### Web example

As a student you can go to the bar each Saturday.

And you need to go back home at some time...

#### We can consider the following "chain" of states (shown above):

- you always start from Home going to Bar 1 or Bar 2.
- after each drink in Bar 1 you have three choices: go Back Home, go to Bar 2 and order another drink in Bar 1.
- if you are already in Bar 2, you have only two choices after each round: go Back Home or order another drink (not shown).
- once you get Back Home, you stay there.

#### Piero Paialunga in Towards Data Science



## Web example



14\_TwoBar.ipynb Open in Colab

January 23, 2025

Even if all transition probabilities are known, it is not always possible to obtain statistical properties of the distribution directly...

But one can simulate Markov Chain state sequence many times... Probability of visiting bars:



## Web example

14\_TwoBar2.ipynb COPen in Colab

#### Probability density for the number of drinks:



We can not only estimate the expected number of drinks (which we could also do from the known probabilities), but also the distribution...

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#### **Another example**

The chain in the web example always ended in the single 'Back Home' state. Not very interesting...



Consider an atom irradiated with the laser light tuned to the excitation energy:

- when in ground state, atom has certain probability (per time unit  $\equiv$  simulation step) to get excited
- when in the excited state, atom can radiate photon and go back to the ground state or, with lower probability, radiate softer photon and go to intermediate meta-stable state.
- when in the meta-stable state, probability of radiation (per unit of time) is very low.


## **Another example**

14\_atom.ipynb CO Open in Colab

Example simulation results starting from ground state, 1000 time steps:



Fast oscillations between ground and excited state, longer stays in meta-stable...

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### **Another example**

14\_atom2.ipynb

Example simulation results starting from ground state, 10000 time steps:



System "forgets" about the initial state fast. We can get distributions for different parameters...

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## Another example

14\_atom2.ipynb

Example simulation results starting from ground state, 10000 time steps: After increasing meta-stable state lifetime:



System "forgets" about the initial state fast. We can get distributions for different parameters...

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# Assume that the state space consists of N states: $s_{(1)}, \ldots, s_{(N)}$ . Then, for each state $s_{(i)}$ on can define a set of on-step transition probabilities:

$$p_{ij} = p(X^{(t+1)} = s_{(j)}|X^{(t)} = s_{(i)})$$

We usually require that these probabilities are time-independent (such chain is called time-homogeneous).



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If we now describe state of the system by a N-component vector:

$$(s_{(i)})_j = \delta_{ij}$$
 e.g.  $s_{(1)} = (1, 0, 0, ..., 0)$ 



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then probabilities for different states to proceed after state  $s_{(i)}$  can be written as:

$$\mathbf{p} = \mathbf{s}_{(i)} \cdot \mathbb{T}$$
 where  $\mathbb{T} = (p_{ij})$ 

is the transition matrix

(Bonamente)

Probabilities of states after n time steps are then given by:

$$\mathbf{p}^{(n)} = s_{(i)} \cdot \mathbb{T}^n$$





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Let  $u_k$  denote the probability that the system returns to the initial state  $s_{(i)}$  in exactly k time steps. We can define the total probability for returning to the initial state:

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States can be classified according to this probability:

- if u = 1 state  $s_{(i)}$  is recurrent,
- if u < 1 state  $s_{(i)}$  is transient.

If state is recurrent, it will certainly be observed again (even, if we have to wait very long), and the system will return to this state infinitely often.

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State  $s_{(i)}$  is accessible from the initial state  $s_{(i)}$ , if there is a non-zero probability of reaching this state from the initial state in finite number of time steps:

$$\left(\mathbf{p}^{(m)}\right)_{j} = \left(s_{(i)} \cdot \mathbb{T}^{m}\right)_{j} > 0$$

for some natural number m.



# (Bonamente)



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If a state s_{(j)} is accessible from a recurrent state s_{(i)},
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```

If a Markov chain has a finite number of states and each state is accessible from any other state, then all states are recurrent.



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A state  $s_{(i)}$  is said to be periodic with period T if system can return to this state only at times t divisible by T:

$$(\mathbf{p}^{(t)})_{j} = \begin{cases} p > 0 & \text{for } t\% T = 0 \\ 0 & t\% T ! = 0 \end{cases}$$

All states of irreducible chain share the same period.



(Bonamente)

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All states of irreducible chain share the same period.

A chain is said to be aperiodic, if return to a given state can occur at any time (corresponding to T = 1 in definition above).



(Bonamente)



In most cases, we do not care about the initial system state, we want to calculate the set of probabilities for a system after a large number n of steps:

$$\mathbf{p}^{\infty} \;\; = \;\; \lim_{n o \infty} \mathbf{p}^{(n)}$$

This probabilities are called limiting probabilities.



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This probabilities are called limiting probabilities.

For a irreducible aperiodic Markov Chain with recurrent states, limiting probabilities correspond to the stationary distribution:

$$\pi = \pi \cdot \mathbb{T}$$

and this distribution is unique. Regardless of the starting point of the chain, the same stationary distribution will eventually be reached.



14\_atom3.ipynb Copen in Colab

Evolution of state probabilities for system starting at 'Ground' state at t = 0



Stationary state reached for  $t\sim 1000$ 

Note logarithmic time scale!

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#### this is what we look for in most cases

There are three possible approaches to finding a stationary solution:

• by running multiple Markov Chain instances and looking at final state distribution, simple but time consuming



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- applying the transfer matrix many times, starting for arbitrary initial state vector



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There are three possible approaches to finding a stationary solution:

- by running multiple Markov Chain instances and looking at final state distribution, simple but time consuming
- applying the transfer matrix many times, starting for arbitrary initial state vector
- by looking for analytic solution to the problem:

$$\pi_j = \sum_j \pi_i \ p_{ij}$$
 stationary distribution  
 $\sum_i \pi_i = 1$  normalization constrain  
 $\pi_i \ge 0$ 



#### Herman Scheepers on Towards Data Science

In the analytic approach the problem can be presented as a set of equations:

$$\begin{pmatrix} \mathbb{T}^{\mathsf{T}} - \mathbb{I} \\ \hline 1 & \dots & 1 \end{pmatrix} \cdot \boldsymbol{\pi} = \begin{pmatrix} 0 \\ \vdots \\ 0 \\ \hline 1 \end{pmatrix}$$
$$\mathbb{A} \quad \cdot \boldsymbol{\pi} = \mathbf{b}$$

which are, however, not independent (the problem is over-constrained).



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which are, however, not independent (the problem is over-constrained).

The simple solution is to multiply both sides by  $\mathbb{A}^{\mathsf{T}}$ :

 $\mathbb{A}^{\mathsf{T}}\mathbb{A}\cdot\boldsymbol{\pi} = \mathbb{A}^{\mathsf{T}}\mathbf{b}$ 

which can now be solved with standard linear algebra procedures...

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# **Markov Chains**

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## **General concept**

Fw

arXiv:0905.1629

We introduced Monte Carlo as an alternative method for integrating an arbitrary function.

Arbitrary parameter space can be considered.



h

#### Rejection technique

Generate uniformly distributed random points, select those in the considered parameter space...

## **General concept**

Fw

arXiv:0905.1629

We introduced Monte Carlo as an alternative method for integrating an arbitrary function.

Arbitrary parameter space can be considered.



.

Rejection technique b Generate uniformly distributed random points, select those in the considered parameter space... Efficiency can be low...





14\_mcmc.ipynb







14\_mcmc.ipynb







14\_mcmc.ipynb

Generation of random points from the surface considered in lecture 05



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CO Open in Colab 14\_mcmc.ipynb





# **General concept**

We do not want to reject events!

Random move procedure: subsequent points generated by random variations of previous ones

# dom move procedure: subseq

Markov Chain Monte Carlo procedure

If the new point is outside the considered parameter space, do not reject it, but take the last point again (!)

#### Can this procedure work ?





h

# Markov Chain MC example

14\_mcmc.ipynb

Using maximum step size:  $\Delta x = \Delta y = 1$ 





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N=100

# Markov Chain MC example

14\_mcmc.ipynb

Using maximum step size:  $\Delta x = \Delta y = 1$ 






14\_mcmc.ipynb

Using maximum step size:  $\Delta x = \Delta y = 1$ 



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14\_mcmc.ipynb 🚺 Open in Colab

Using maximum step size:  $\Delta x = \Delta y = 1$ 

Markov Chain simulation example N = 100000Single variable distribution  $0.16^{1}$ 1.00 0.5 0.75 0.4 0.8 0.3 0.50 0.2 0.25 0,1 0.0 -1.00 -0.75 -0.50 -0.25 0.00 0.25 0.50 0.75 1.00 ≻ 0.00 190 -0.25 0.8 Count 0.6 -0.50 -0.750.2 -1.000:0 ···1.00 -0.75 -0.50 -0.25 0.00 0.25 0.50 0.75 1.00 · -1.00 -0.75 -0.50 -0.25 0.00 0.25 0.50 0.75 1.00 N=100 000

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CO Open in Colab 14\_mcmc.ipynb

Using maximum step size:  $\Delta x = \Delta y = 1$ 











We can reduce number of "duplicates" by reducing step:  $\Delta x = \Delta y = 0.2$ 



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N = 1000





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# More general case

14\_mcmc2.ipynb

#### Gaussian probability distribution in the considered parameter space



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N = 1000

# More general case

14\_mcmc2.ipynb

#### Gaussian probability distribution in the considered parameter space



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## More general case



#### Gaussian probability distribution in the considered parameter space





# Metropolis–Hastings algorithm

Consider chain described by on-step transition probability  $p(X^{(t+1)}|X^{(t)})$ 

- To generate points distributed according to f(X), for each step t:
  - generate candidate point  $X^{\star}$  from  $p(X^{\star}|X^{(t)})$
  - compute the Metropolis-Hastings ratio:

$$R = \frac{f(X^{\star}) p(X^{(t)}|X^{\star})}{f(X^{(t)}) p(X^{\star}|X^{(t)})}$$

• for the next step take

$$X^{(t+1)} = \begin{cases} X^* & \text{with probability } p^* = \min\{R, 1\} \\ X^{(t)} & \text{otherwise.} & \text{with probability } 1 - p^* \end{cases}$$





(Givens)



14\_mcmc3.ipynb

Using maximum step size:  $\Delta x = \Delta y = 1$ 





14\_mcmc3.ipynb

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14\_mcmc3.ipynb

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14\_mcmc3.ipynb

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14\_mcmc3.ipynb

Using maximum step size:  $\Delta x = \Delta y = 0.05$ 





14\_mcmc3.ipynb

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14\_mcmc3.ipynb

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14\_mcmc3.ipynb

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14\_mcmc3.ipynb

Using maximum step size:  $\Delta x = \Delta y = 0.2$ 





14\_mcmc3.ipynb

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14\_mcmc3.ipynb

Using maximum step size:  $\Delta x = \Delta y = 0.2$ 





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## Markov Chain MC example (2)

14\_mcmc3.ipynb

Using maximum step size:  $\Delta x = \Delta y = 0.2$ 





# **Markov Chains**

- 1 Markov Chains
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- 🕘 Final exam

# **Bayesian approach**

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(lecture 01)

Bayes theorem can be used to generalize the concept of probability.

In particular, one can consider "probability" of given hypothesis H (theoretical model or model parameter, eg. Hubble constant) when taking into known outcome D (data) of the experiment

$$P(H|D) = \frac{P(D|H)}{P(D)} \cdot P(H)$$

There are two problems with this approach:

- H can not be considered an event, sampling space can not be properly defined
- we need to make a subjective assumption about the "prior" P(H) describing our initial belief in hypothesis H

For these reasons I try to avoid it, and do not refer to P(H|D) as "probability". Rather use "degree of belief" for results of the procedure applied to non random events



(lecture 06)

The likelihood function:

 $L(\boldsymbol{\lambda}, \mathbf{x}) = \prod_{j=1}^{N} f(\mathbf{x}^{(j)}; \boldsymbol{\lambda})$ 

describes the probability of a given measurement results  ${\sf x}$  for the selected parameter values  ${\sf \lambda}.$ 

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describes the probability of a given measurement results **x** for the selected parameter values  $\lambda$ .

In the bayesian approach we can refer it to "probability distribution" for the parameters  $\lambda$ :

 $f(\boldsymbol{\lambda}) \sim L(\boldsymbol{\lambda}, \mathbf{x}) \cdot p(\boldsymbol{\lambda})$ 

where  $p(\lambda)$  is the assumed prior distribution for parameters  $\lambda$ . (usually flat)



(lecture 06)

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If we know  $f(\lambda)$ , we can construct Markov Chain in  $\lambda$  space.

With Metropolis–Hastings algorithm, starting from arbitrary  $\lambda^{(0)}$  point, the chain should converge to  $f(\lambda)$  distribution for  $N \to \infty$ .

A.F.Żarnecki

Statictical analysis 14



1000 events were collected in the muon lifetime measurement. Distribution can be described by the formula:

 $\frac{dN}{dt} = \frac{N_{sig}}{\tau} e^{-\frac{t}{\tau}} + \frac{dN_{bg}}{dt}$ with flat background level known to be  $\frac{dN_{bg}}{dt} = 10 \pm \Delta \ \mu s^{-1}$ 





# Histogram can be fitted using iterative $\chi^2$ minimization procedure

(without bg constraint)





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#### Parameter evolution in the Markov Chain



#### Stable distribution obtained already after about 100 iterations

A.F.Żarnecki

Statictical analysis 14

14\_mcfit1.ipynb

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Statictical analysis 14

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Parameter distributions after  $N = 10\ 000$  iterations (skipping first 100)





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Parameter distributions after  $N = 100\ 000$  iterations (skipping first 1000)


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Parameter distributions after  $N = 100\ 000$  iterations (skipping first 1000) Including background level constraint





### Nominal solution from Markov Chain (mean values of parameters)





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But we can also get the probability distribution of the fit results:





14\_mcfit3.ipynb

But we can also get the probability distribution of the fit results:





14\_mcfit3.ipynb COPPEN in Colab

But we can also get the probability distribution of the fit results:







Markov Chain Monte Carlo does not work "out of the box"



### It converges fast with the proper choice of parameter variation steps

A.F.Żarnecki



Markov Chain Monte Carlo does not work "out of the box"



Convergence can be very slow, if parameter steps too small...

A.F.Żarnecki



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Markov Chain Monte Carlo does not work "out of the box"



### Fluctuations significantly increased, if steps are too large...

A.F.Żarnecki

200 events were measured in the electron scattering experiment. The expected distribution corresponds to that of the single slit diffraction:

$$p(x) = C a \cdot \left(\frac{\sin a(x-x_0)}{a(x-x_0)}\right)^2$$

where *a* is the scaling factor and  $x_0$  is the position of the maximum.









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### **Final remarks**

Fw

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However, it is crucial to make sure they converge, before using their output for the analysis. Algorithm tuning may be required...

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Events generated with Markov Chain MC are not independent! One should not use subsequent events together in the analysis (eg. for background estimates)



## **Markov Chains**

- Markov Chains
- 2 Markov Chain Monte Carlo
- 3 Application to parameter fitting



### **Final exam**

As described in the syllabus, assessment will be based on home exercises and the written exam. 50% of points collected from exercises and exam (with same weights) required to pass.

For the written exam, you will have to solve five problems similar to those in homeworks (maybe a little bit more complex, as you get 13 points for each).

Problems will be put on Kampus on Sunday, February 2nd, and you should upload solutions to Kampus (each one as a separate file) within one week, till Sunday, February 9th (23:55).

By uploading the solutions to Kampus you declare that they resulted from your own work and that you have not shared nor discussed them with anyone.

