

Simulations of the Higgs boson measurements at the PLC

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Outline

- Introduction
- SM Higgs boson production
- Heavy MSSM Higgs bosons
- CP properties of the Higgs boson
- PLC complementarity

Introduction

Simulations of the Higgs boson production at the PLC

Studies by P.Nieźurawski, A.F.Żarnecki, M.Krawczyk (NŻK):

$\mathcal{H} \rightarrow b\bar{b}$ decay channel

- Standard Model

“The SM Higgs boson production $\gamma\gamma \rightarrow h \rightarrow b\bar{b}$ at the photon collider at TESLA,”
Acta Phys. Polon. B **34** (2003) 177 [arXiv:hep-ph/0208234].

- MSSM

“Extended analysis of the MSSM Higgs boson production at the photon collider,”
Proceedings of LCWS 2005 [arXiv:hep-ph/0507006];

“LHC wedge at the PLC: Observability of $\gamma\gamma \rightarrow A, H \rightarrow b\bar{b}$,”
Acta Phys. Polon. B **37** (2006) 1187.

⇒ see also P. Nieźurawski, “Higgs-boson production at the photon collider at TESLA,”
arXiv:hep-ph/0503295 (PhD Thesis).

Introduction

Simulations of the Higgs boson production at the PLC

Studies by P.Nieżurawski, A.F.Żarnecki, M.Krawczyk (NŻK):

$\mathcal{H} \rightarrow WW/ZZ$ decay channels

- **Standard Model**

“Study of the Higgs-boson decays into $W^+ W^-$ and $Z Z$ at the photon collider,”
JHEP **0211** (2002) 034 [arXiv:hep-ph/0207294].

- **2HDM**

“Determination of the Higgs-boson couplings and CP properties in the SM-like two Higgs doublet model,” JHEP **0502** (2005) 041 [arXiv:hep-ph/0403138].

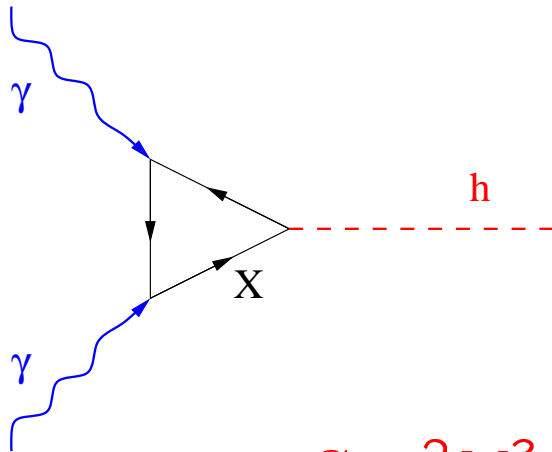
- **Generic model**

“Model-independent determination of CP violation from angular distributions in Higgs boson decays to $W W$ and $Z Z$ at the Photon Collider,”
Acta Phys. Polon. B **36** (2005) 833 [arXiv:hep-ph/0410291].

$$\gamma\gamma \rightarrow h$$

Width

Two-photon width of the Higgs boson $\Gamma_{\gamma\gamma}$ is sensitive to **all** massive and charged **particles** in the loop:



$$\Gamma(h \rightarrow \gamma\gamma) = \frac{G_F \alpha^2 M_h^3}{128 \sqrt{2} \pi^3} \cdot |\mathcal{A}|^2$$

where:

$$\mathcal{A} = A_W(M_W) + \sum_f N_c Q_f^2 A_f(M_f) + \dots$$

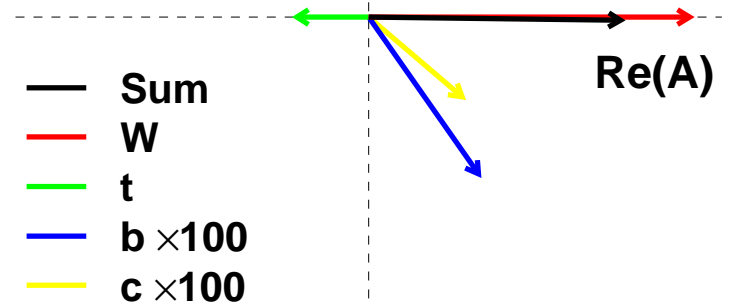
two-photon amplitude

Amplitude

For Higgs boson mass $M_h < 2 M_W$ amplitude \mathcal{A} is **real**

$M_h = 120. \text{ GeV}$

$\text{Im}(\mathcal{A})$



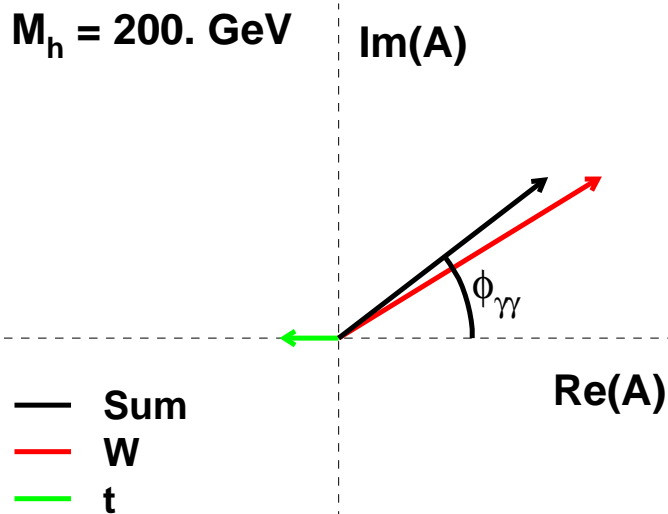
contribution to $\text{Im}(\mathcal{A})$ from light fermions - very tiny

contribution from new heavy particles - real

$$\gamma\gamma \rightarrow h$$

Phase

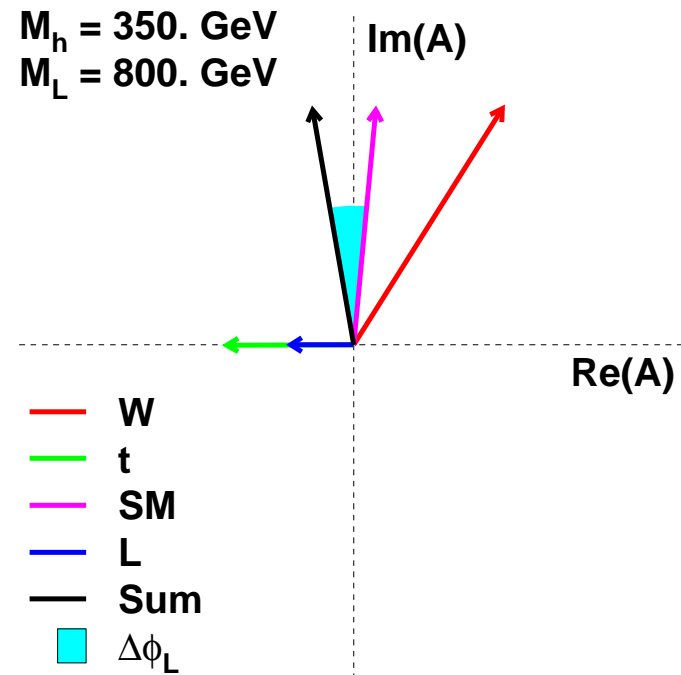
However, for $M_h > 2 M_W$
 W contribution is **complex**



$\Rightarrow \mathcal{A} = |\mathcal{A}| \cdot e^{i\phi}$ phase $\phi_{\gamma\gamma} \neq 0$
 $\Gamma_{\gamma\gamma} \sim |\mathcal{A}|^2 = \text{Im}(\mathcal{A})^2 + \text{Re}(\mathcal{A})^2$

New particles

Will contribute to both $\Gamma_{\gamma\gamma}$ and $\phi_{\gamma\gamma}$.
 For new charged lepton:



for $M_h \sim 350 \text{ GeV}$
 \Rightarrow amplitude mostly **imaginary**: $\text{Re}(\mathcal{A}) \sim 0$
 $\Rightarrow \Gamma_{\gamma\gamma}$ little sensitive to new particles !

$$\gamma\gamma \rightarrow h$$

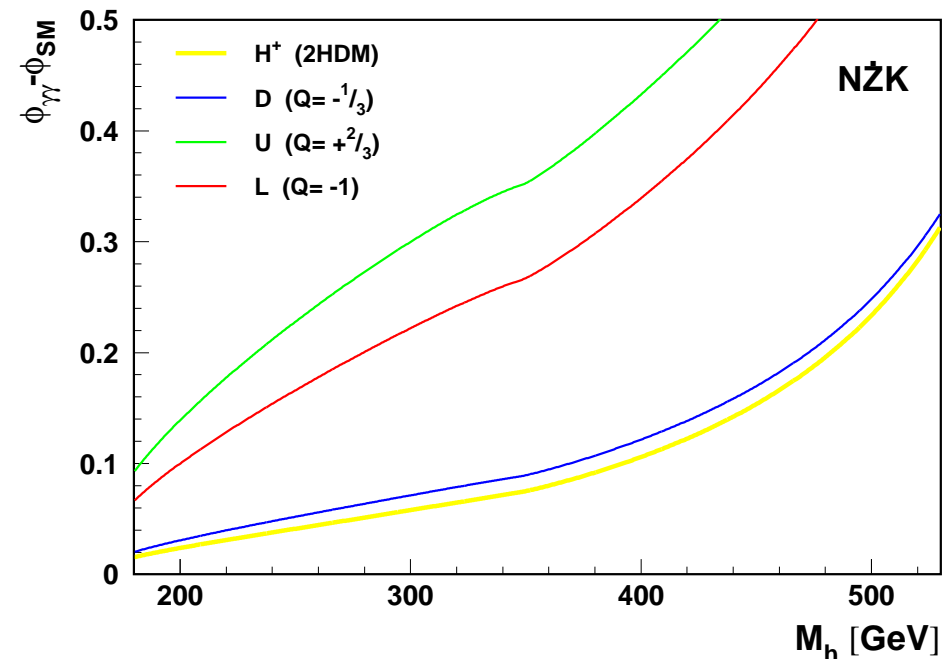
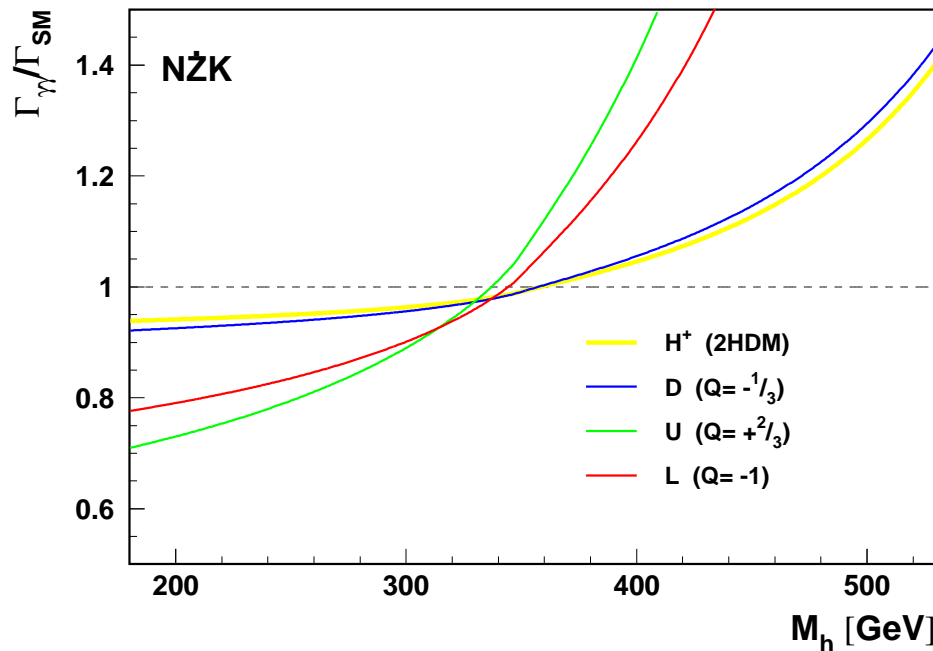
Expected deviations from SM

Contribution to $\Gamma_{\gamma\gamma}$ and $\phi_{\gamma\gamma}$ from new heavy charged particles with mass ~ 800 GeV
 For SM with fourth-generation fermions and for SM-like 2HDM (II) type A
 (additional contribution due to charged Higgs boson)

width

hep-ph/0207294

phase



⇒ little effect on $\Gamma_{\gamma\gamma}$ for $M_h \sim 350$ GeV

⇒ significant change in $\phi_{\gamma\gamma}$

Can we measure it !?

Measuring the two-photon decay width of Intermediate-mass Higgs at a photon-photon collider

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1-3-1 Kagamiyama, Higashi-Hiroshima, 739, Japan

I. Watanabe

Akita Junior College,[†]

46-1 Morisawa, Sakura, Shimokitate, Akita, 010 Japan

(November 22, 2007)

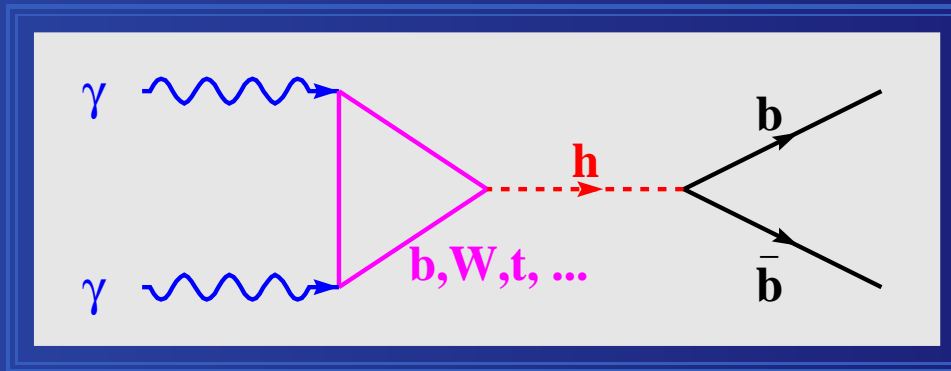
Abstract

Feasibility of a measurement of the partial decay width of the intermediate-mass Higgs boson into two photons at a photon-photon collider

$$m_h \approx 120 \text{ GeV}$$

Process: $\gamma + \gamma \rightarrow h \rightarrow b + \bar{b}$

$$J_z = 0$$



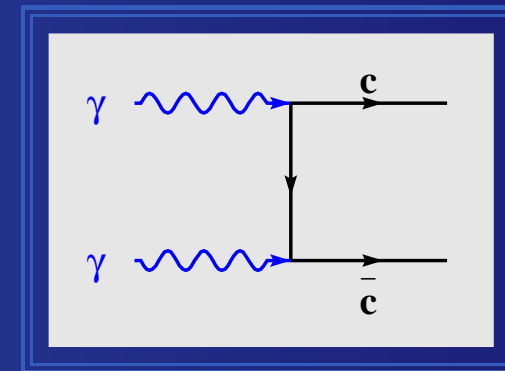
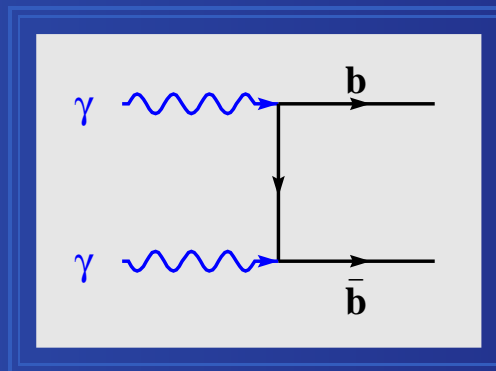
“Hard” background:

- $\gamma + \gamma \rightarrow b + \bar{b}$

- $\gamma + \gamma \rightarrow c + \bar{c}$

$$\sigma \propto Q_q^4$$

$$\sigma^{LO}(|J_z| = 2) \gg \sigma^{LO}(J_z = 0)$$



Other background:

- Resolved photon(s) interactions $\gamma + \gamma \rightarrow X + Q + \bar{Q}$

- Overlaying events

(high intensity of photon-beams in the low-energy part of the spectrum)

$$\gamma\gamma \longrightarrow \textit{higgs} \longrightarrow b\bar{b}$$

Analysis of precision $\sigma(\gamma\gamma \rightarrow \textit{higgs} \rightarrow b\bar{b})$ measurement includes:

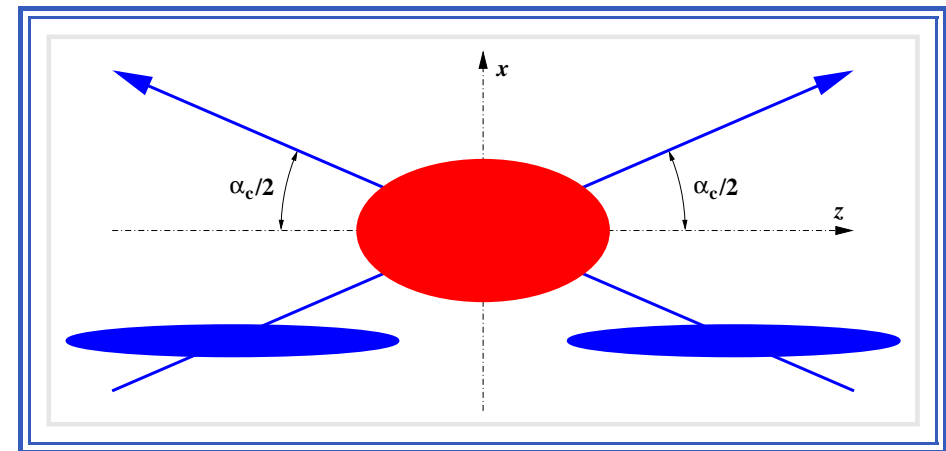
- realistic $\gamma\gamma$ -spectra
- Background: NLO $Q\bar{Q}(g)$ (G. Jikia)
- overlying events $\gamma\gamma \rightarrow \textit{hadrons}$ (OE)
- Detector performance: SIMDET 4.01
- b -tagging (ZVTOP-B-Hadron-Tagger)
- Signal: HDECAY, PYTHIA (PS)

Special treatment:

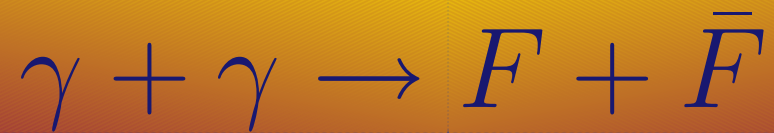
- crossing angle
- primary vertex distribution
- $\gamma\gamma \rightarrow W^+W^-$ background

Crab-crossing of beams

$$\alpha_c = 34 \text{ mrad}$$



	Bunch	Primary vertex
σ_x	140 nm	$3.6 \mu\text{m}$
σ_y	15 nm	11 nm
σ_z	0.3 mm	0.2 mm



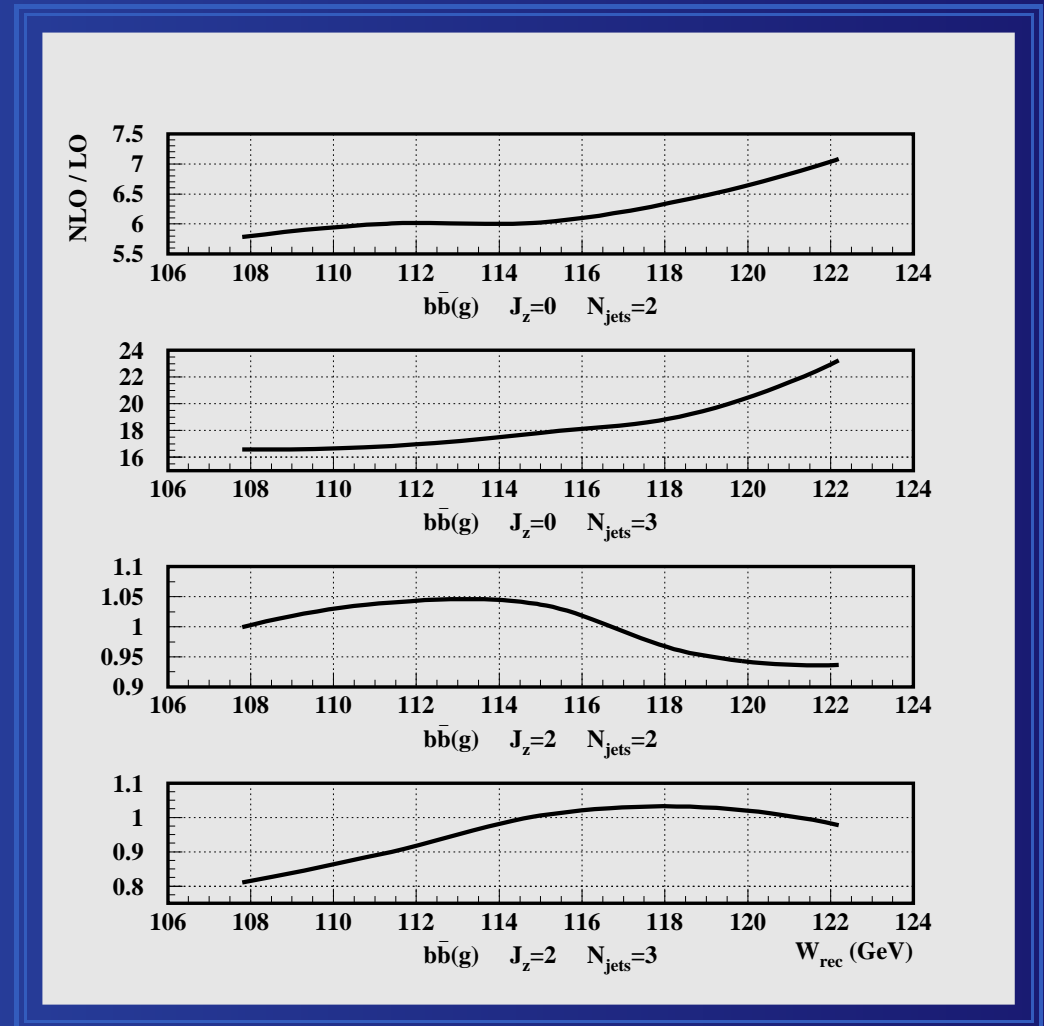
- NLO cross section for massless fermions

$$\Rightarrow \sigma \propto \frac{\alpha^2 \alpha_s}{s}$$

$$\frac{d\sigma}{dE_g}(J_z = 2) \propto \frac{1}{E_g}$$

$$\sigma(J_z = 0) \propto E_g^3$$

- NLO cross section for massive fermions

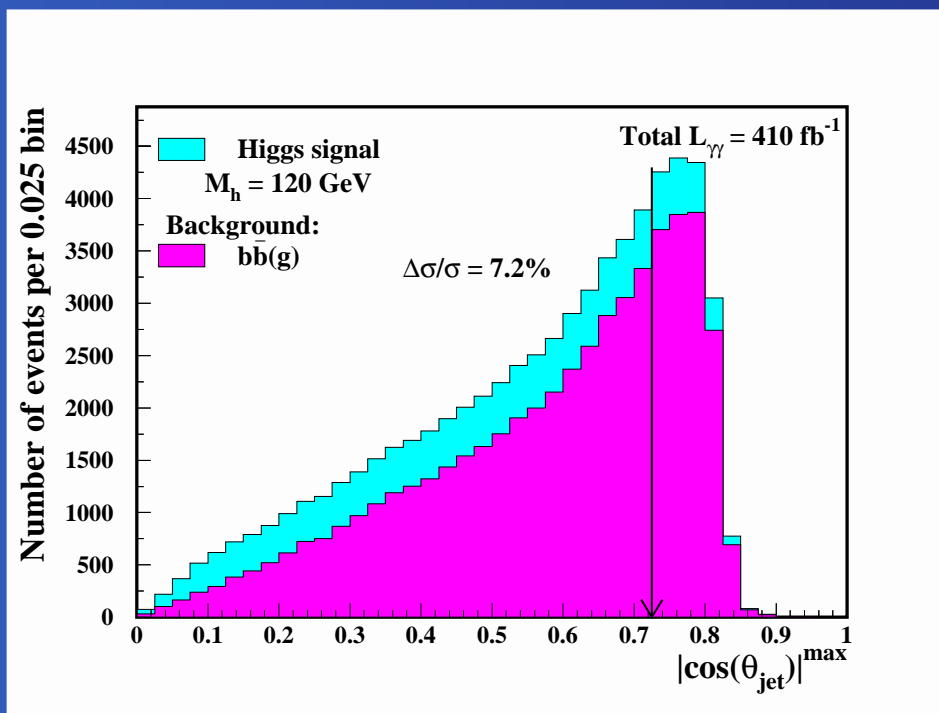


Cuts

Cuts optimized by minimizing:

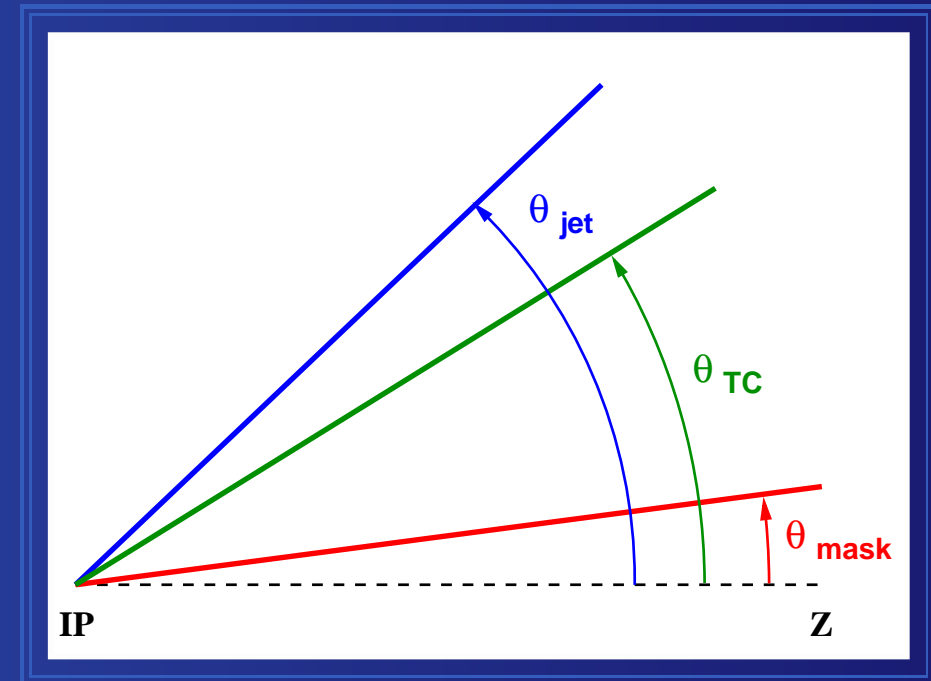
$$\frac{\Delta\sigma(\gamma\gamma \rightarrow h \rightarrow b\bar{b})}{\sigma(\gamma\gamma \rightarrow h \rightarrow b\bar{b})} = \frac{\sqrt{\mu_S + \mu_B}}{\mu_S},$$

For example:



Maximal value of $|\cos\theta_{jet}|$
over all jets in the event

All angular cuts



Detector mask

Particles on Pythia level: $\cos\theta_{mask} \approx 0.99$

OE suppression

Tracks & clusters: $\cos\theta_{TC} = 0.85$

$\gamma\gamma \rightarrow Q\bar{Q}(g)$ suppression

Jets: $|\cos\theta_{jet}|^{\max} = 0.725$



higgs-tagging at $M_h = 120 \text{ GeV}$

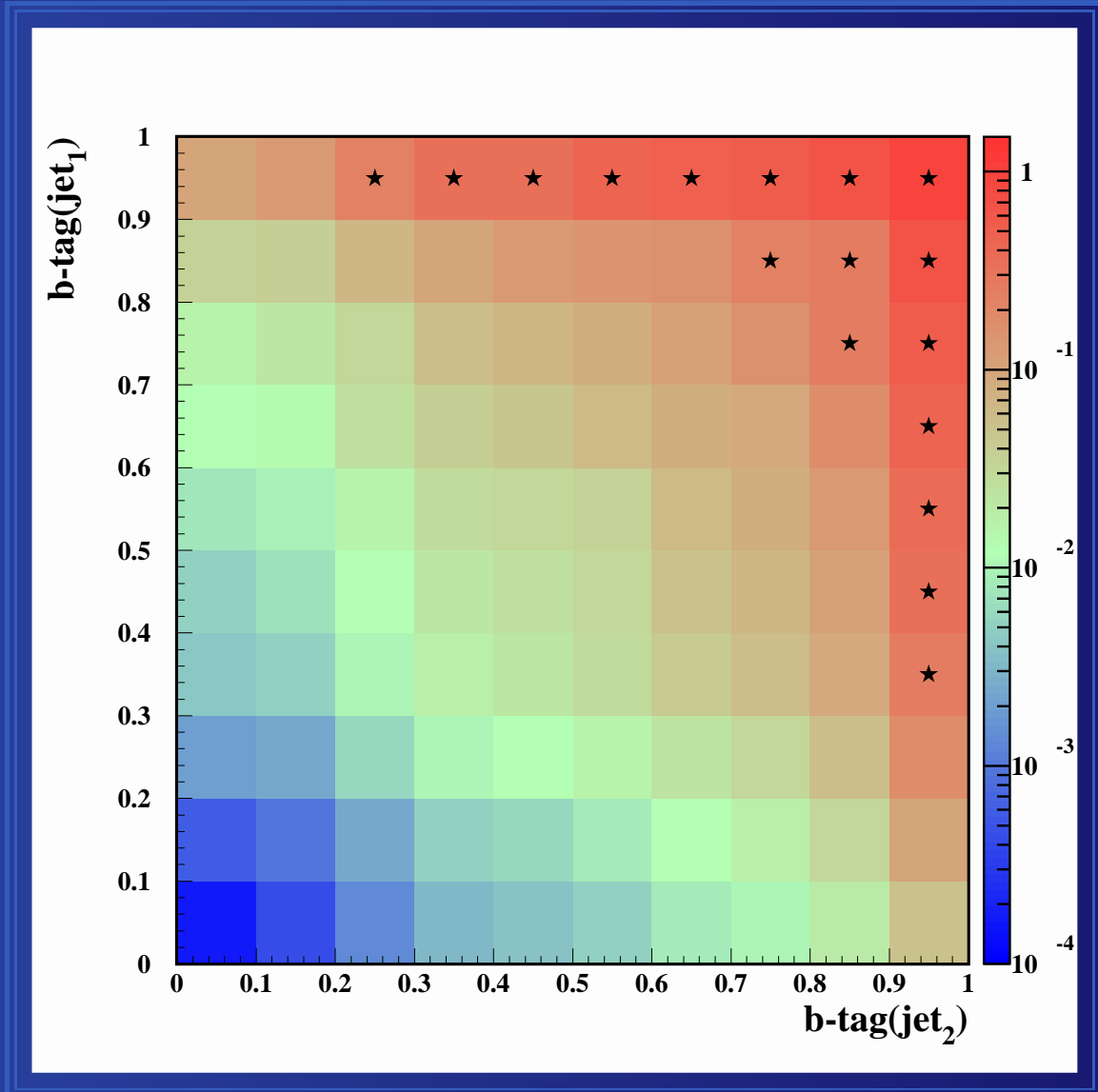
higgs-tagging: a cut on the ratio
of $\gamma\gamma \rightarrow h \rightarrow b\bar{b}$
to $\gamma\gamma \rightarrow b\bar{b}(g), c\bar{c}(g), q\bar{q}$ ($q = u, d, s$)
events

$\Rightarrow \varepsilon_h = 58\%$
 $\varepsilon_{bb} = 50\%$
 $\varepsilon_{cc} = 2.2\%$
 $\varepsilon_{uds} = 0.16\%$

Without OE

$\Rightarrow \varepsilon_h = 71\%$
 $\varepsilon_{bb} = 64\%$
 $\varepsilon_{cc} = 2.9\%$
 $\varepsilon_{uds} = 0.11\%$

**Tighter cuts are needed
due to OE contribution**



Reconstruction & Selection

Selection of $b\bar{b}$ events for $M_h = 120$ GeV:

- OE suppression: clusters & tracks with $|\cos \theta_i| > \cos \theta_{TC} = 0.85$ ignored
- $W_{rec} > 1.2 W_{\gamma\gamma}^{\min}$
- **Jets:** Durham algorithm, $y_{cut} = 0.02$
- $N_{jets} = 2, 3$
- for each jet: $|\cos \theta_{jet}| < 0.725$
- $|P_z|/E < 0.1$

Rejection of W^+W^- events (for $M_h = 150, 160$ GeV):

- for each jet: $M_{jet} < 70$ GeV
- energy below θ_{TC} : $E_{TC} < 90$ GeV
- for each jet: $N_{trk} \geq 4$

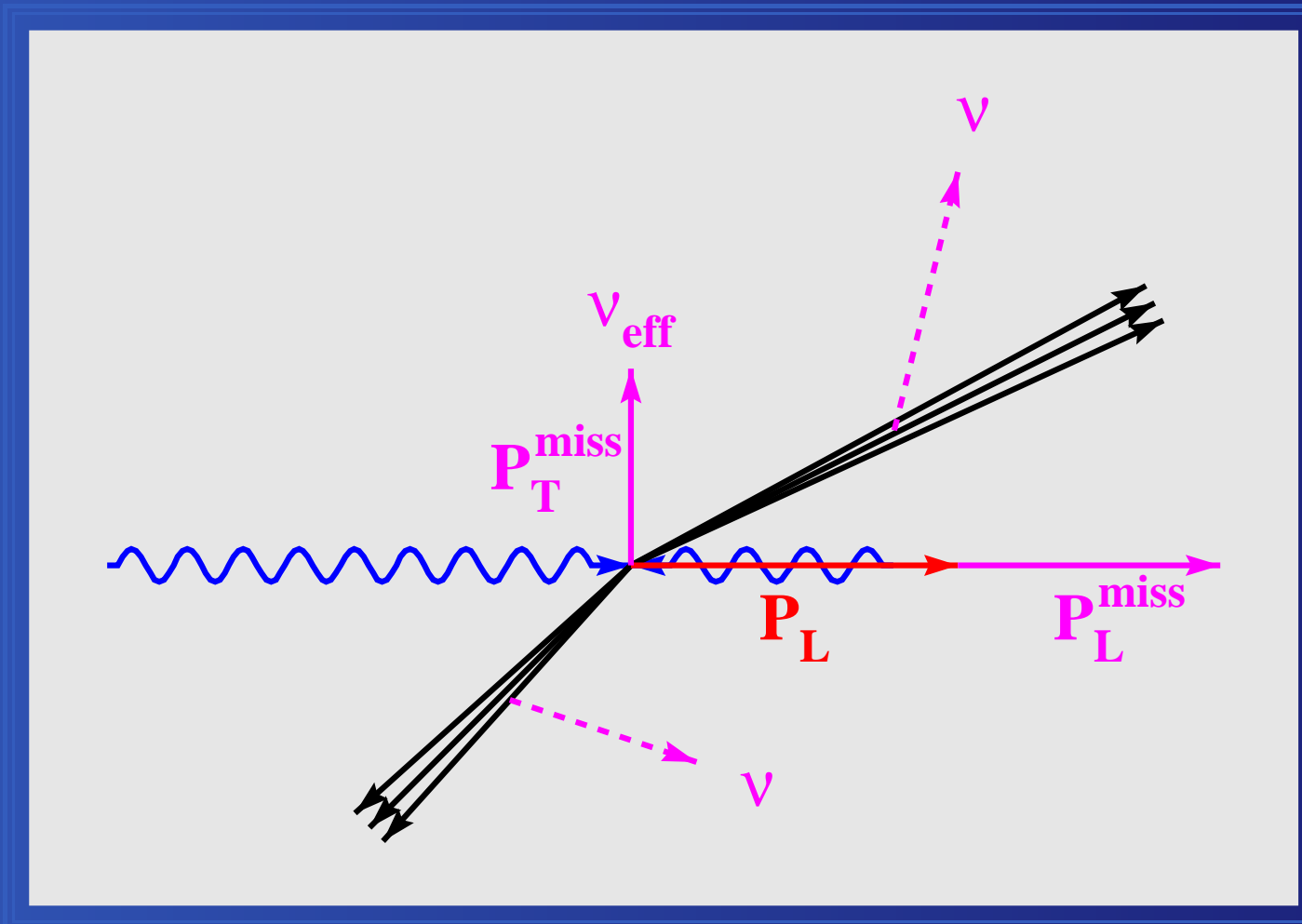
b-tagging: ZVTOP-B-HADRON-TAGGER (T. Kuhl)

Correction for crossing angle: jets boosted with $\beta = -\sin(\alpha_c/2)$



Missing P_T

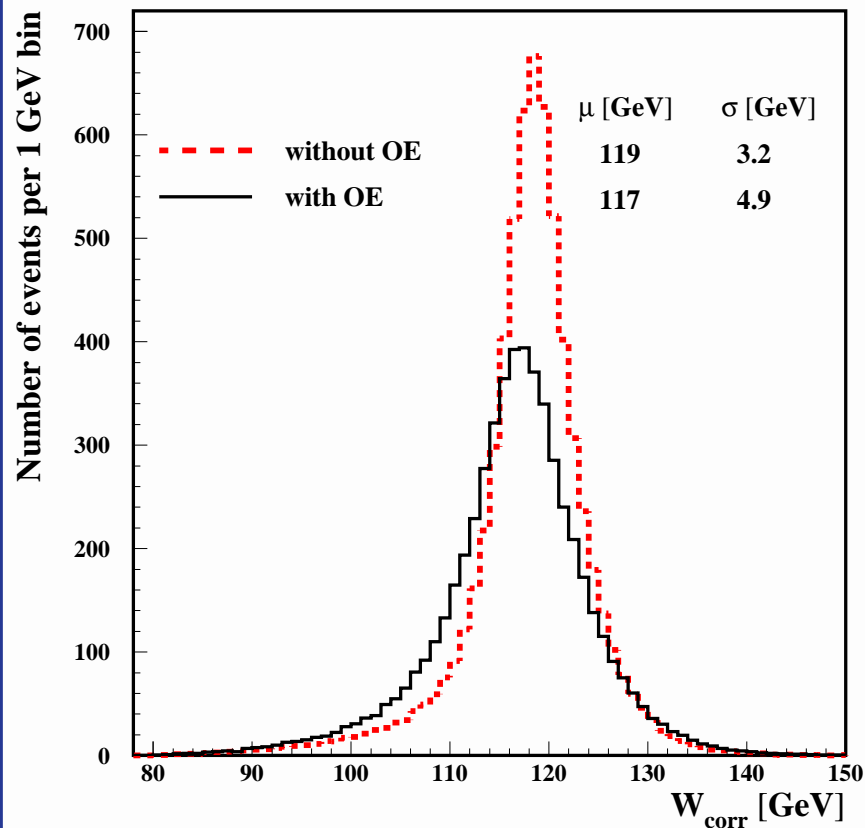
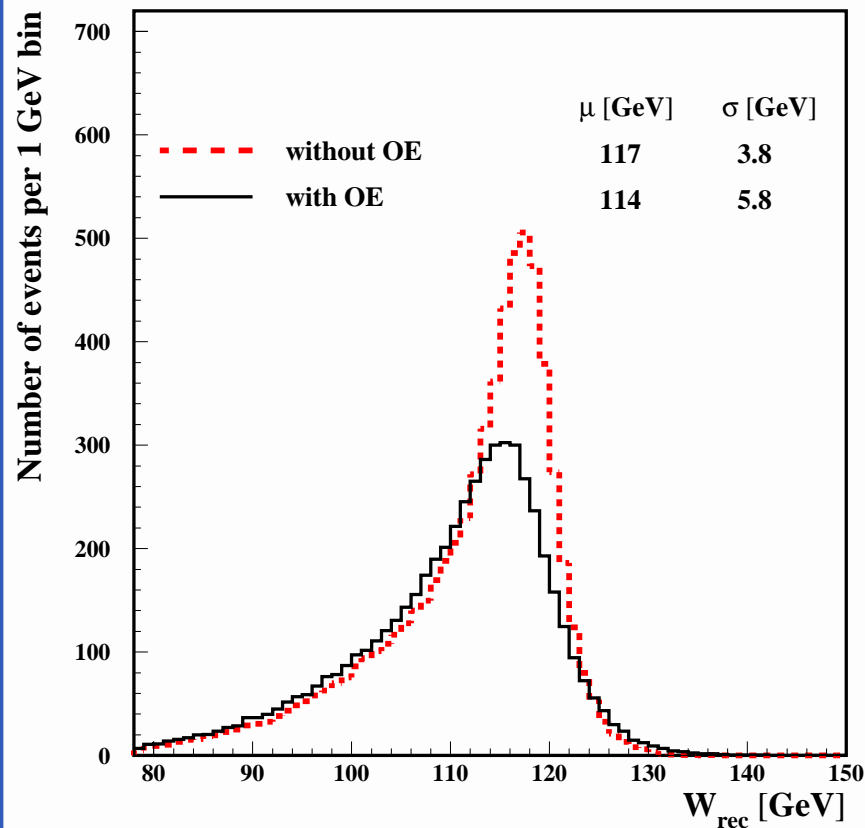
Neutrinos from semileptonic decays of D - and B -mesons.



$$W_{\text{corr}} \equiv \sqrt{W_{\text{rec}}^2 + 2P_T(E_{\text{vis}} + P_T)}$$



SM, $M_h = 120$ GeV



Without OE: 6450 events
With OE : 5530 events

$$W_{corr} \equiv \sqrt{W_{rec}^2 + 2P_T(E + P_T)}$$

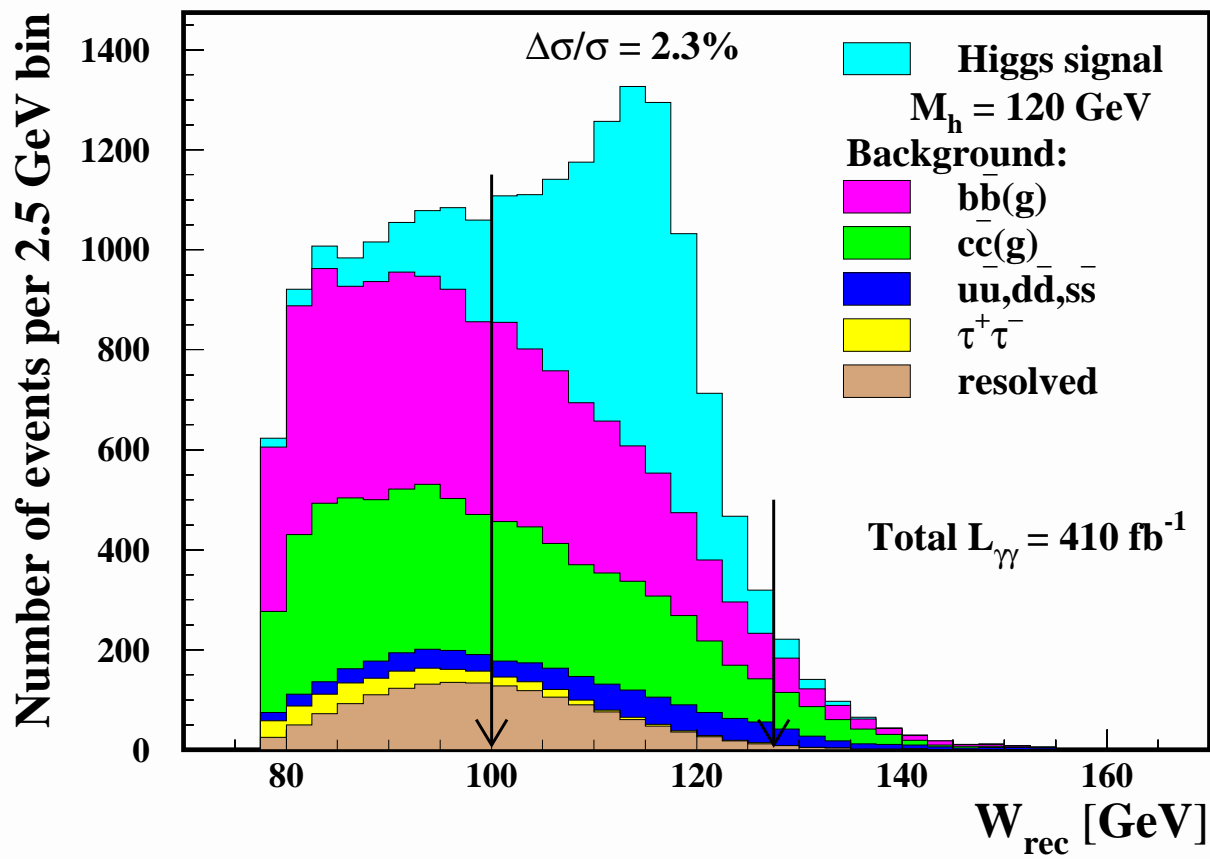
Acta Phys. Pol. B34 177 2003, hep-ph/0208234

Gaussian fit from $\mu - 1.3\sigma$ to $\mu + 1.3\sigma$.



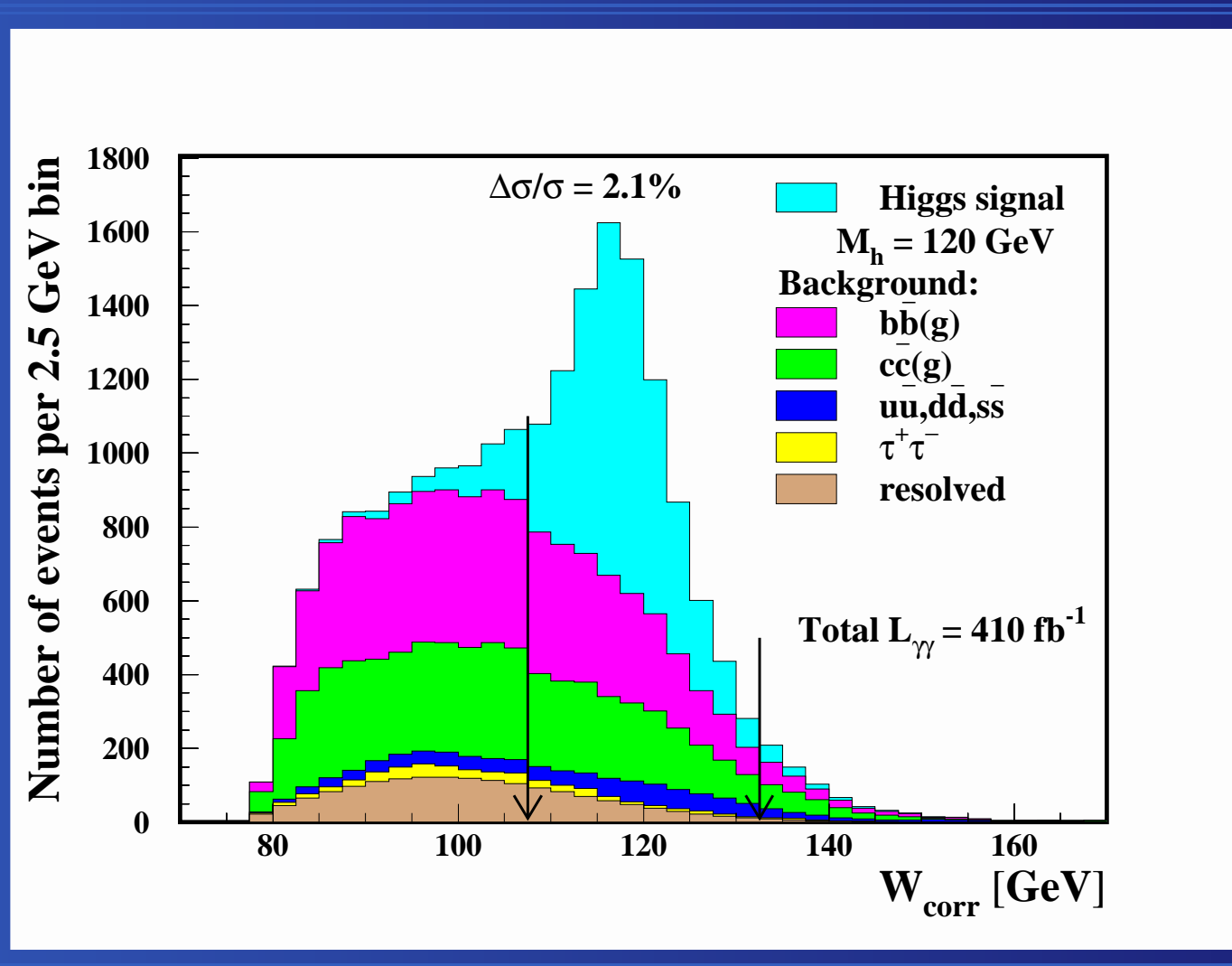
SM, $M_h = 120 \text{ GeV}$

Final results

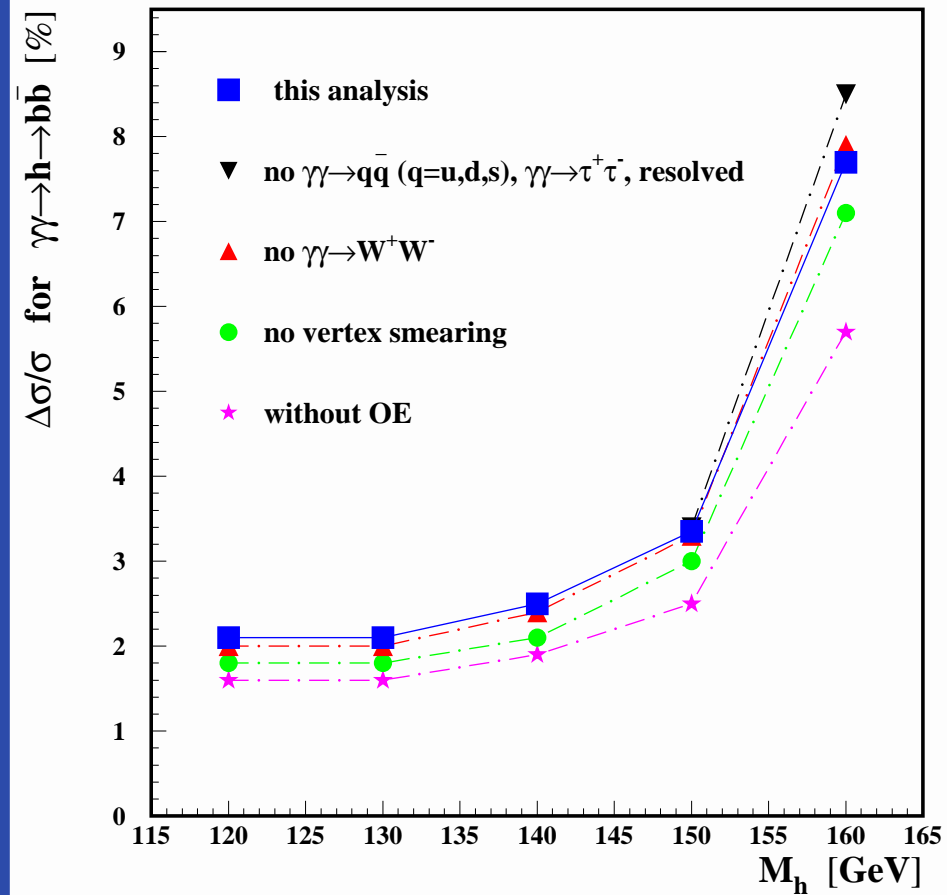


SM, $M_h = 120 \text{ GeV}$

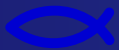
Final results



SM summary



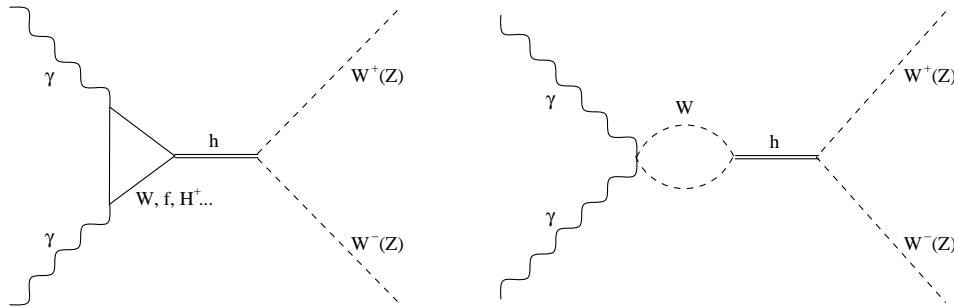
For $M_h = 150, 160$ GeV additional cuts to reduce $\gamma\gamma \rightarrow W^+W^-$.



$$\gamma\gamma \rightarrow (h) \rightarrow W^+W^-, ZZ$$

Higgs signal

For $M_h > 2 M_W$ (phase $\phi_{\gamma\gamma} \neq 0$)
decays to W^+W^- dominate:

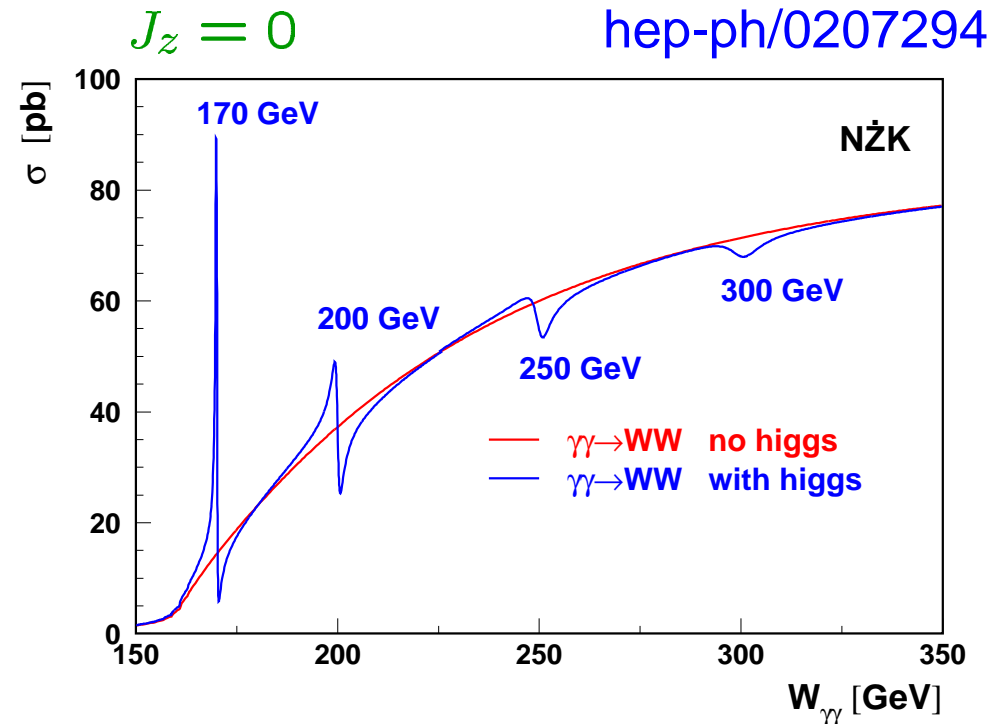


“resonant”

There is a large background from “direct”,
non-resonant production $\gamma\gamma \rightarrow W^+W^-$

Interference

Resonant and direct amplitudes interfere
Large effects expected:



Destructive interference dominates above ~ 200 GeV

G.Belanger, F.Boudjema, Phys.Lett.B288 (1992) 210;
D.A.Morris, et al., Phys. Lett. B323 (1994) 421;
I.F.Ginzburg, I.P.Ivanov, Phys. Lett. B408 (1997) 325.

$$\gamma\gamma \rightarrow (h) \rightarrow W^+W^-, ZZ$$

Phase measurement

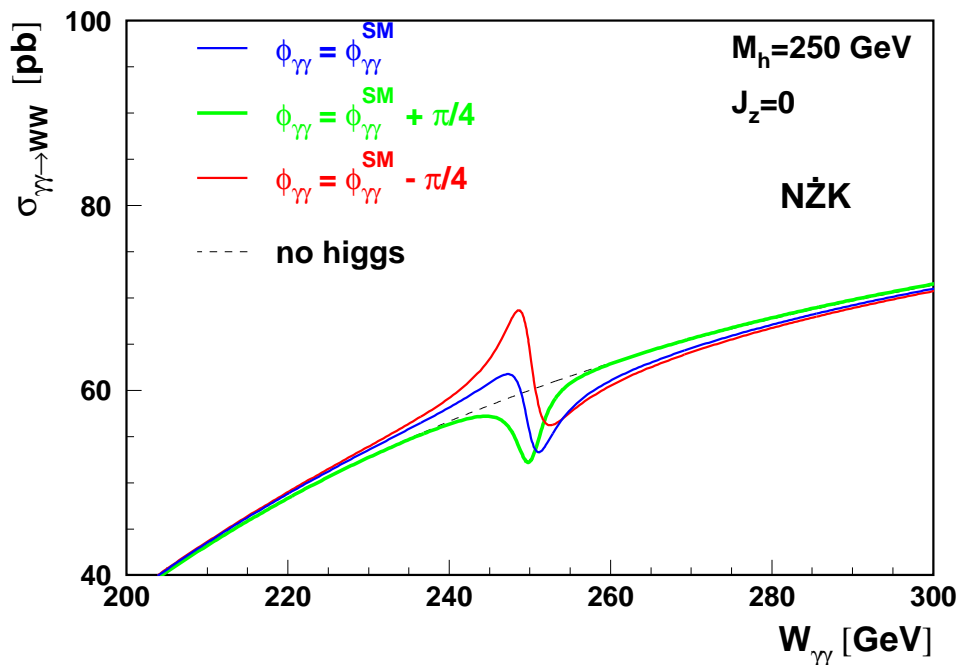
Interference term is sensitive to the phase $\phi_{\gamma\gamma}$ of the $\gamma\gamma \rightarrow h$ amplitude

$$\gamma\gamma \rightarrow ZZ$$

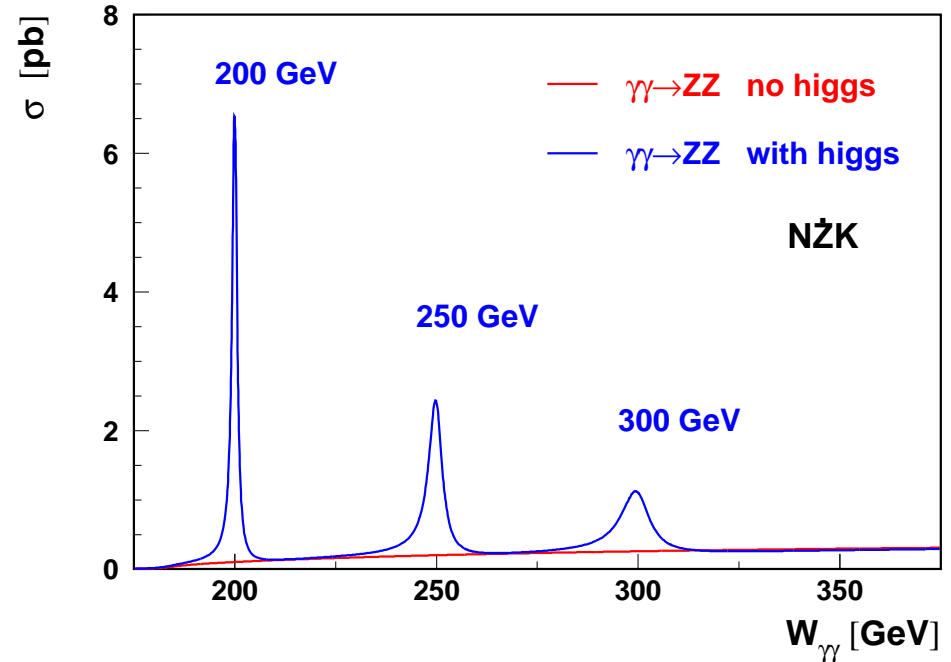
Non-resonant background only at loop level

hep-ph/0207294

$\sigma(\gamma\gamma \rightarrow W^+W^-)$ dependence on $\phi_{\gamma\gamma}$



⇒ can be measured !



⇒ small interference effects

G.J.Gounaris et al., Eur. Phys. J. C13 (2000) 79.

Analysis

Parametrization

“Measured” invariant mass distribution for selected W^+W^- and ZZ events is described by convolution of:

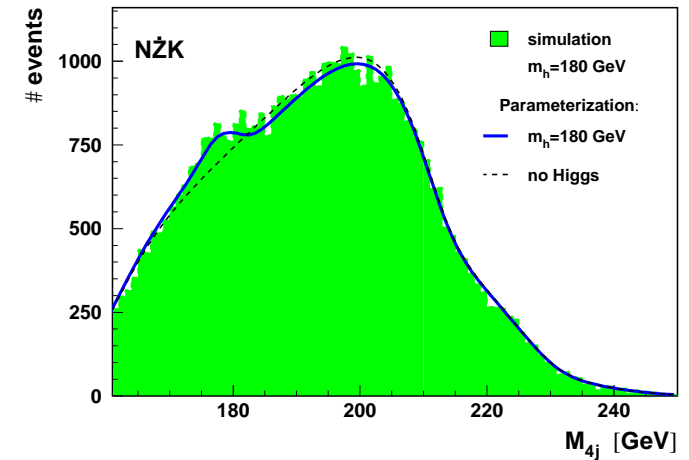
- Analytical luminosity Spectra **CompAZ**
- Cross section formula for signal + background + **interf.**
- Invariant mass resolution parametrized as a function of $W_{\gamma\gamma}$

⇒ mass spectra can be calculated for any $\sqrt{s_{ee}}$ and m_h without time-consuming MC simulation

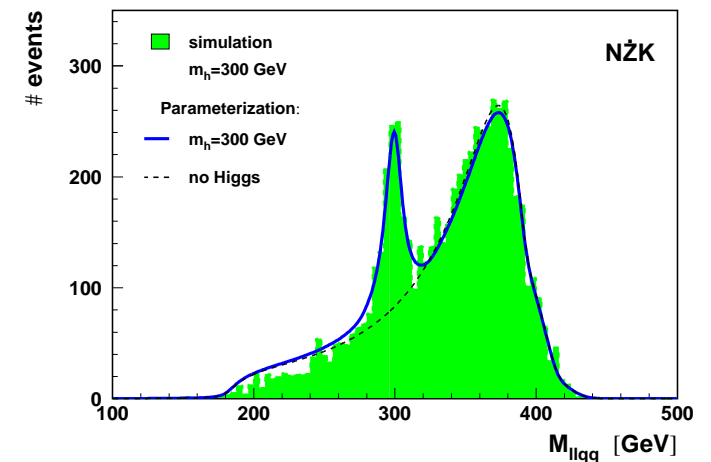
⇒ can be used for fast simulation and fitting

Comparison with full simulation:

W^+W^-



ZZ



SM results

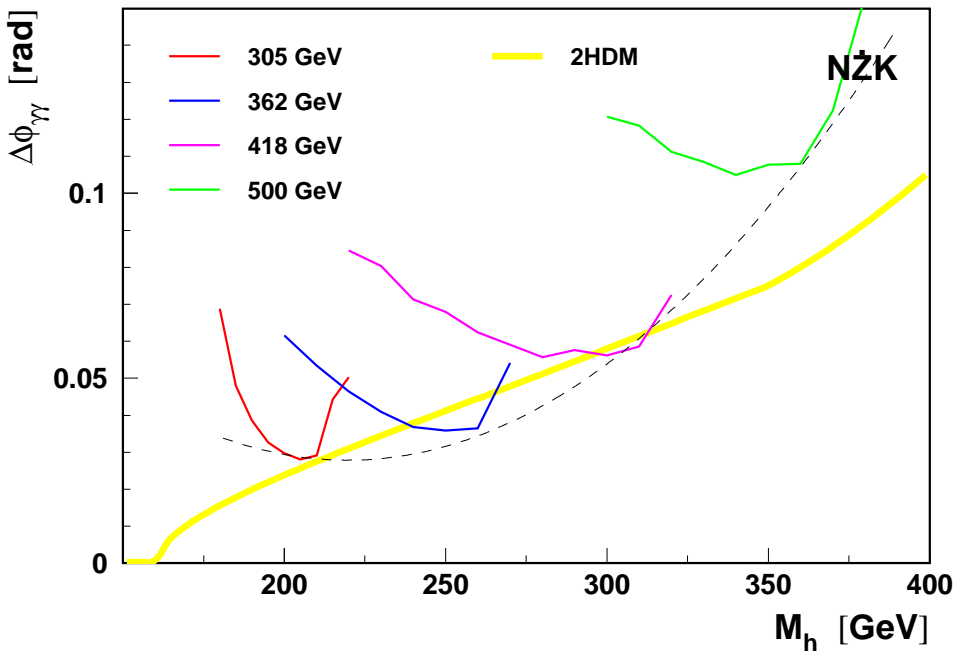
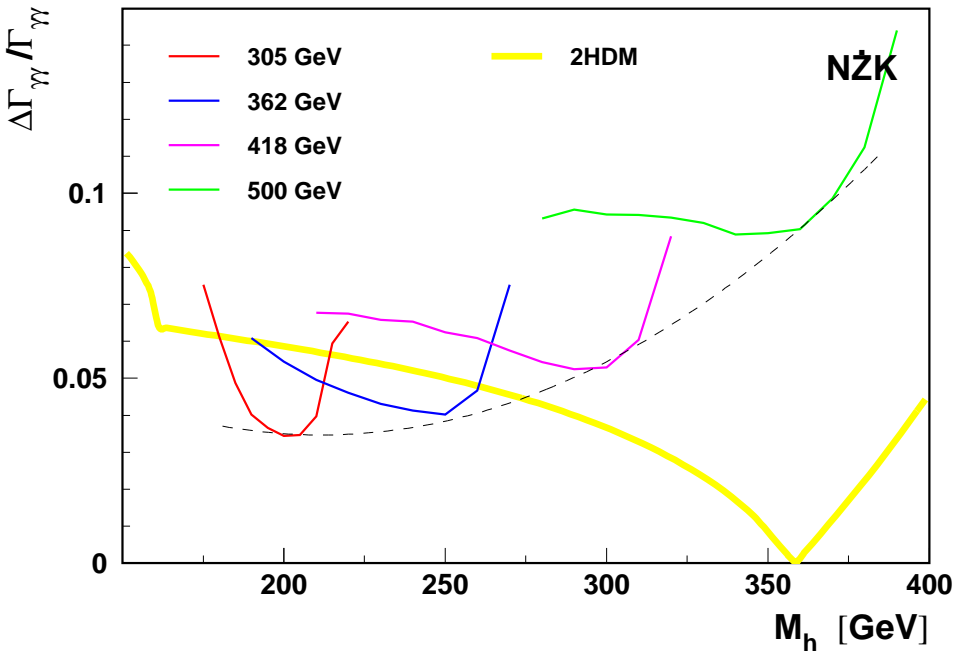
Two parameter fit to W^+W^- and ZZ invariant mass distribution

Expected statistical precision, assuming SM branching ratios (1 PC year):

$\Gamma_{\gamma\gamma}$

hep-ph/0207294

$\phi_{\gamma\gamma}$



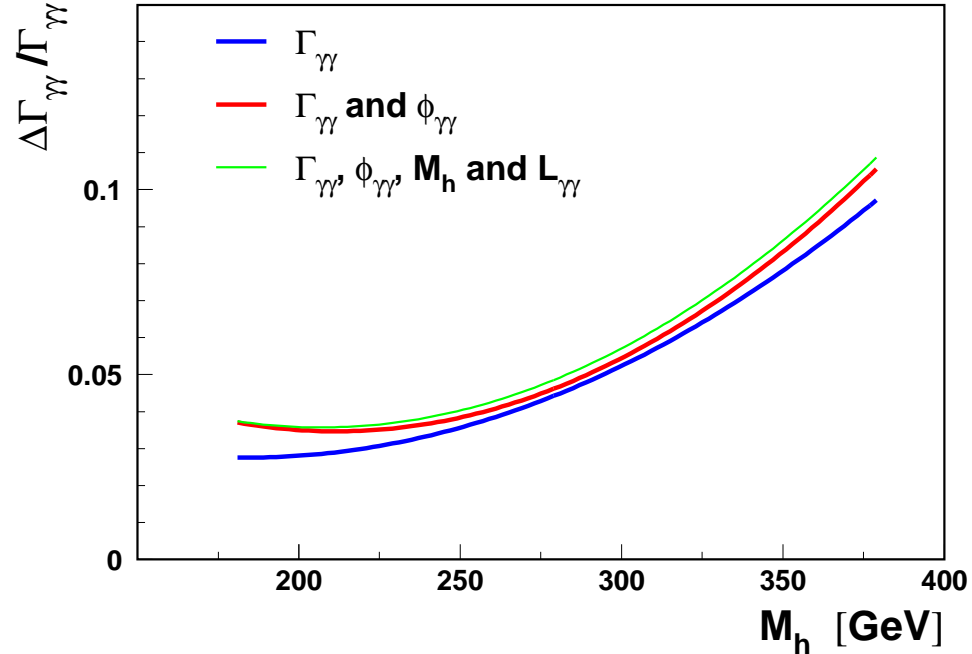
Phase measurement significantly improves our sensitivity to new heavy charged particles e.g. heavy charged Higgs boson of the SM-like 2HDM(II) with $M_{H^\pm} = 800$ GeV at large Higgs boson masses

SM results

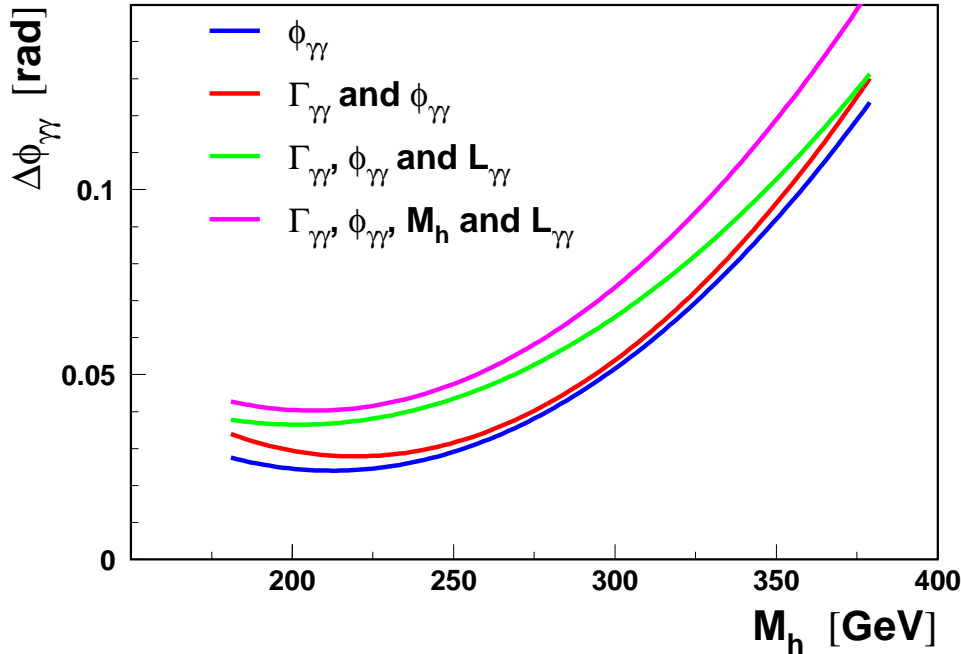
Systematic effects

Statistical precision for different choices of fit parameters:

$\Gamma_{\gamma\gamma}$



$\phi_{\gamma\gamma}$



Precise knowledge of M_h and luminosity not crucial
 \Rightarrow can be constrained by the data itself

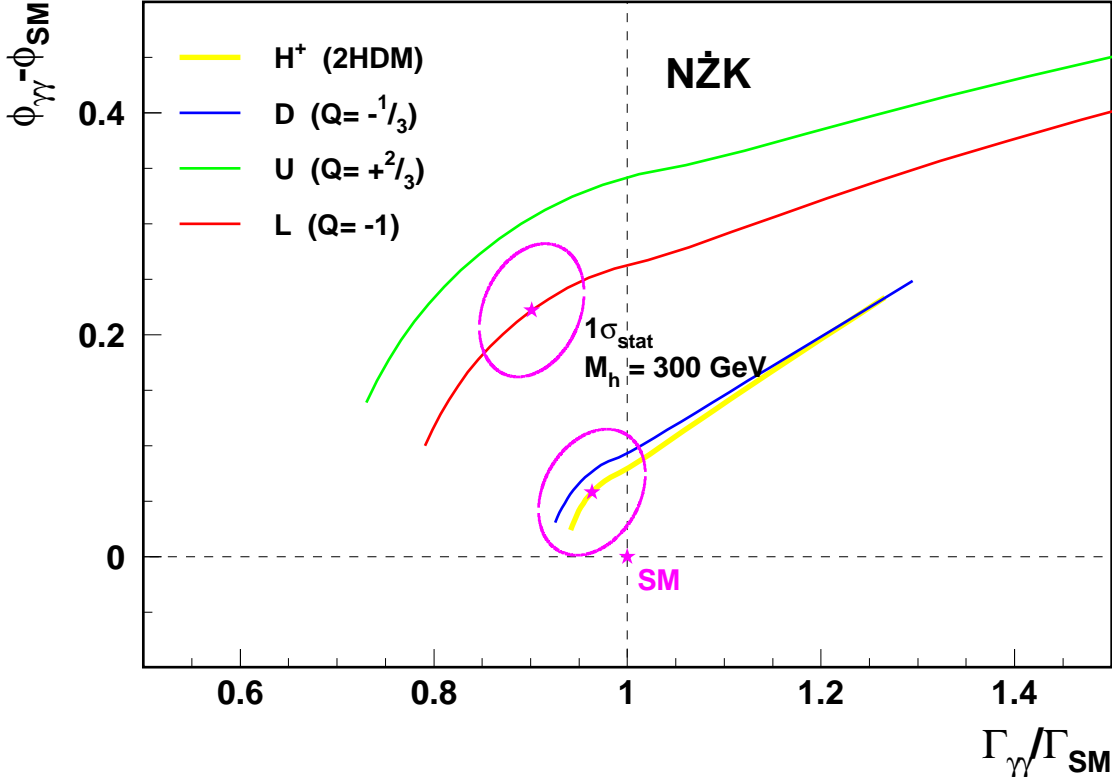
\Rightarrow sensitive to M_h and luminosity uncertainty

SM results

Expected statistical error contours (1σ) in $\phi_{\gamma\gamma} - \Gamma_{\gamma\gamma}$, for $M_h = 300$ GeV:

4th generation lepton
 $M_L = 800$ GeV \Rightarrow

SM-like 2HDM (II) \Rightarrow
 $M_{H^+} = 800$ GeV

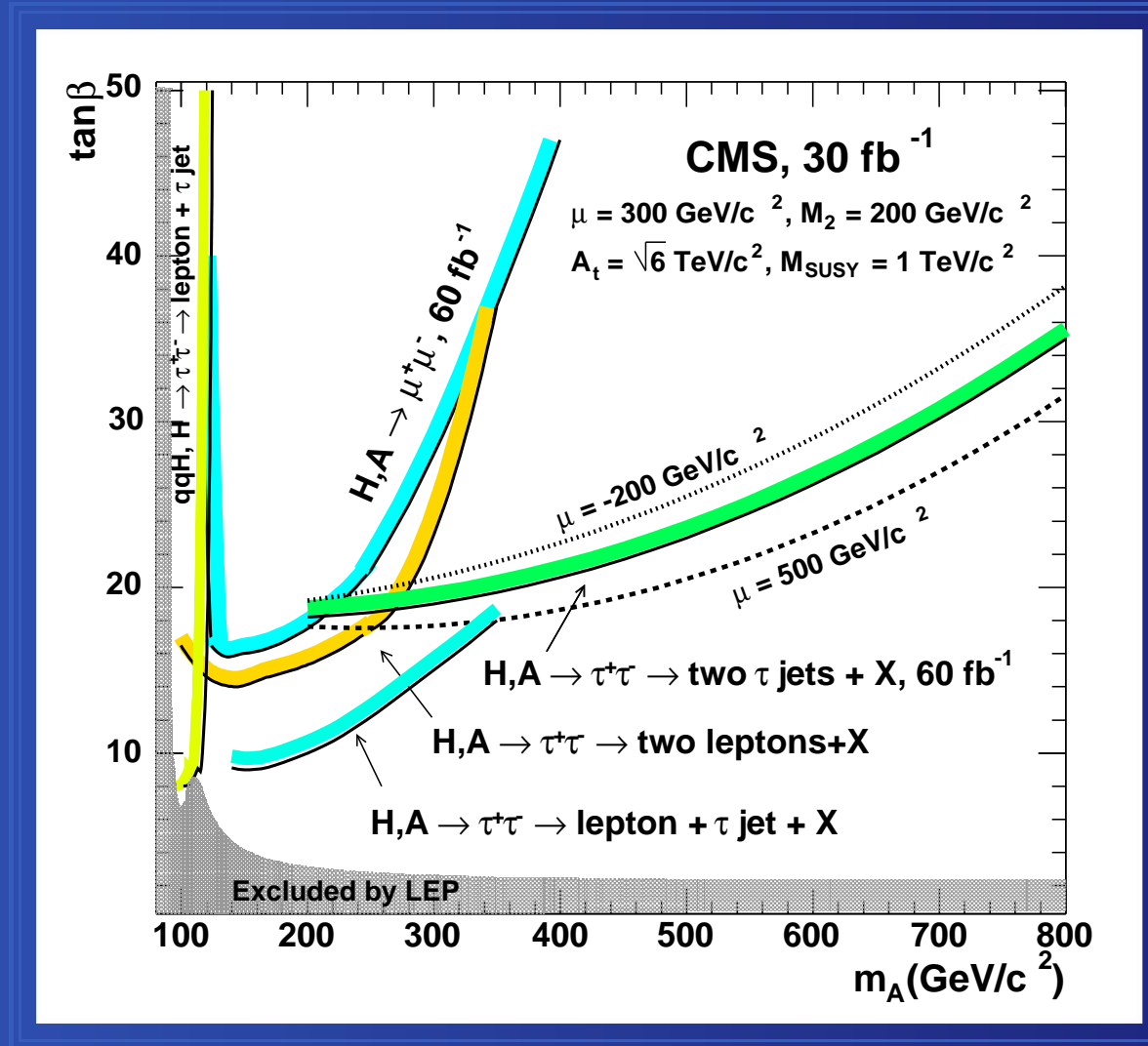


\Rightarrow separation not possible without phase measurement !

Two parameter fit to W^+W^- and ZZ invariant mass distribution; one year of Photon Collider running.

MSSM Study: LHC wedge

LHC wedge: only one light SM-like higgs boson h can be seen

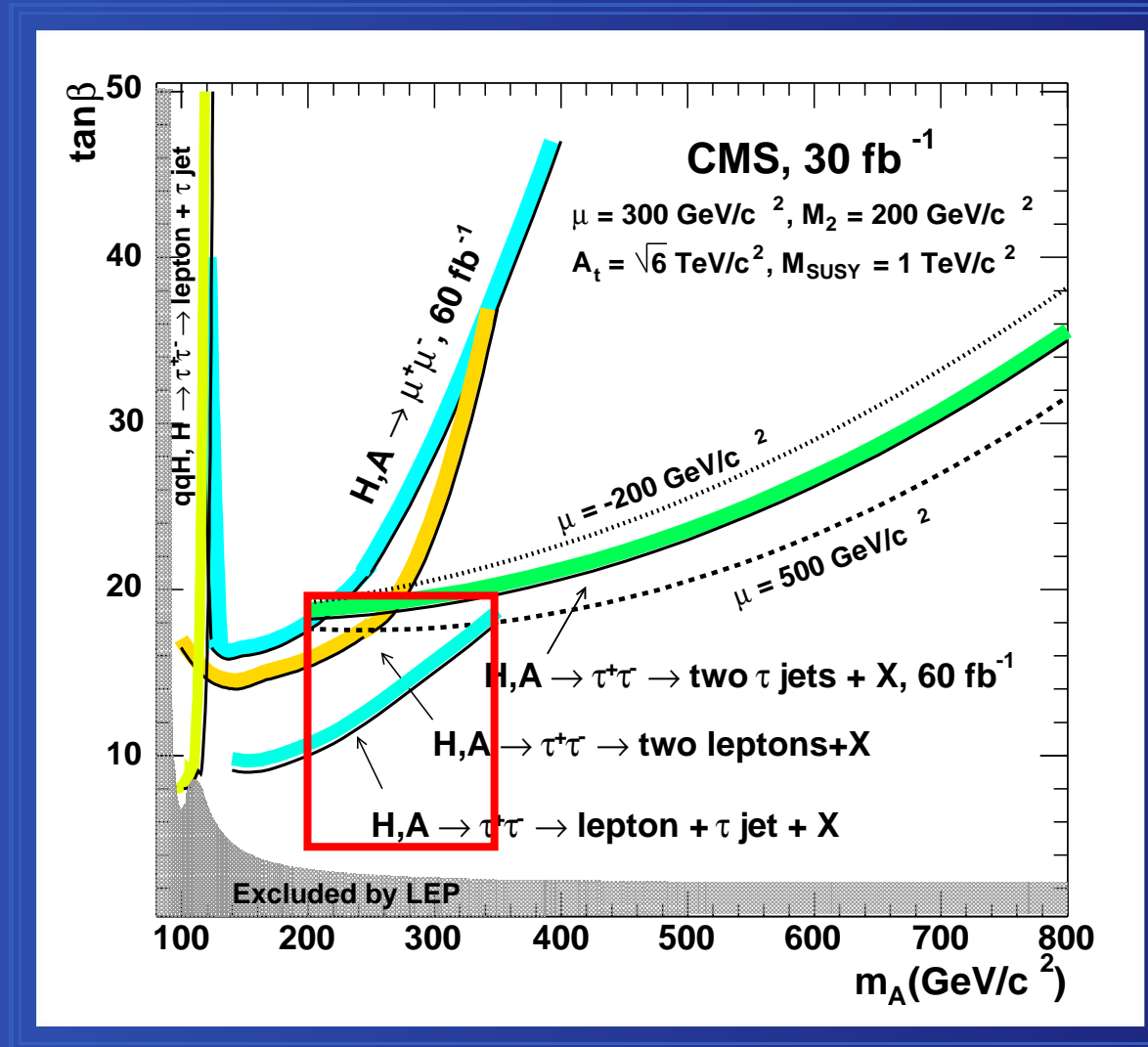


Photon Collider offers unique possibility for precise H/A studies



Choice of MSSM parameters

LHC wedge



From: CMS NOTE 2003/033
 (the same results as in newer CMS CR 2004/058)



Choice of MSSM parameters

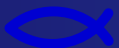
We consider following MSSM parameter sets:

Symbol	μ [GeV]	M_2 [GeV]	$A_{\tilde{f}}$ [GeV]	$M_{\tilde{f}}$ [GeV]
I	200	200	1500	1000
II	-150	200	1500	1000
III	-200	200	1500	1000
IV	300	200	2450	1000

I and III – as in M. Mühlleitner *et al.* but we take higher $A_{\tilde{f}}$ to have M_h above limit

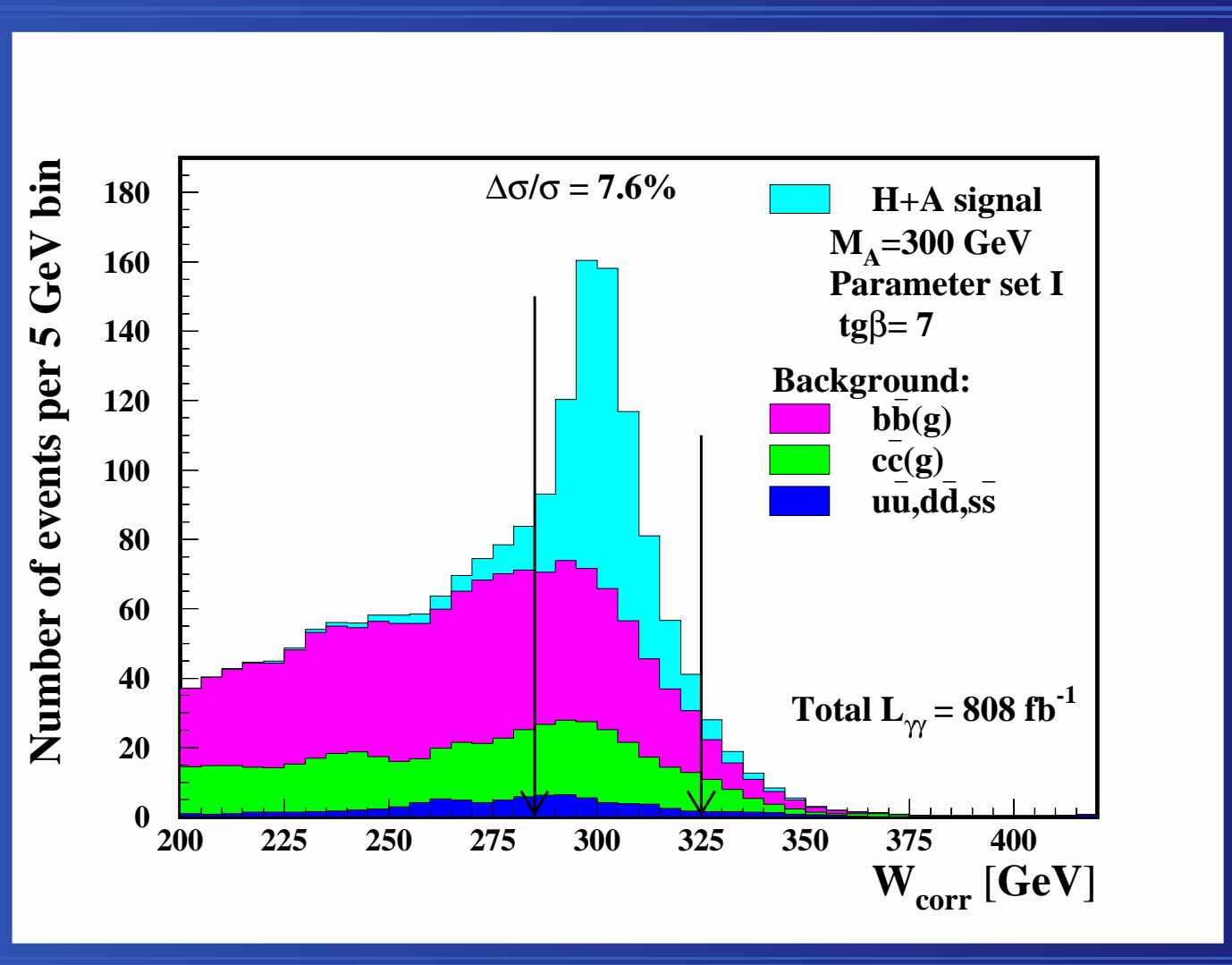
II – an intermediate scenario

IV – as in CMS NOTE 2003/033



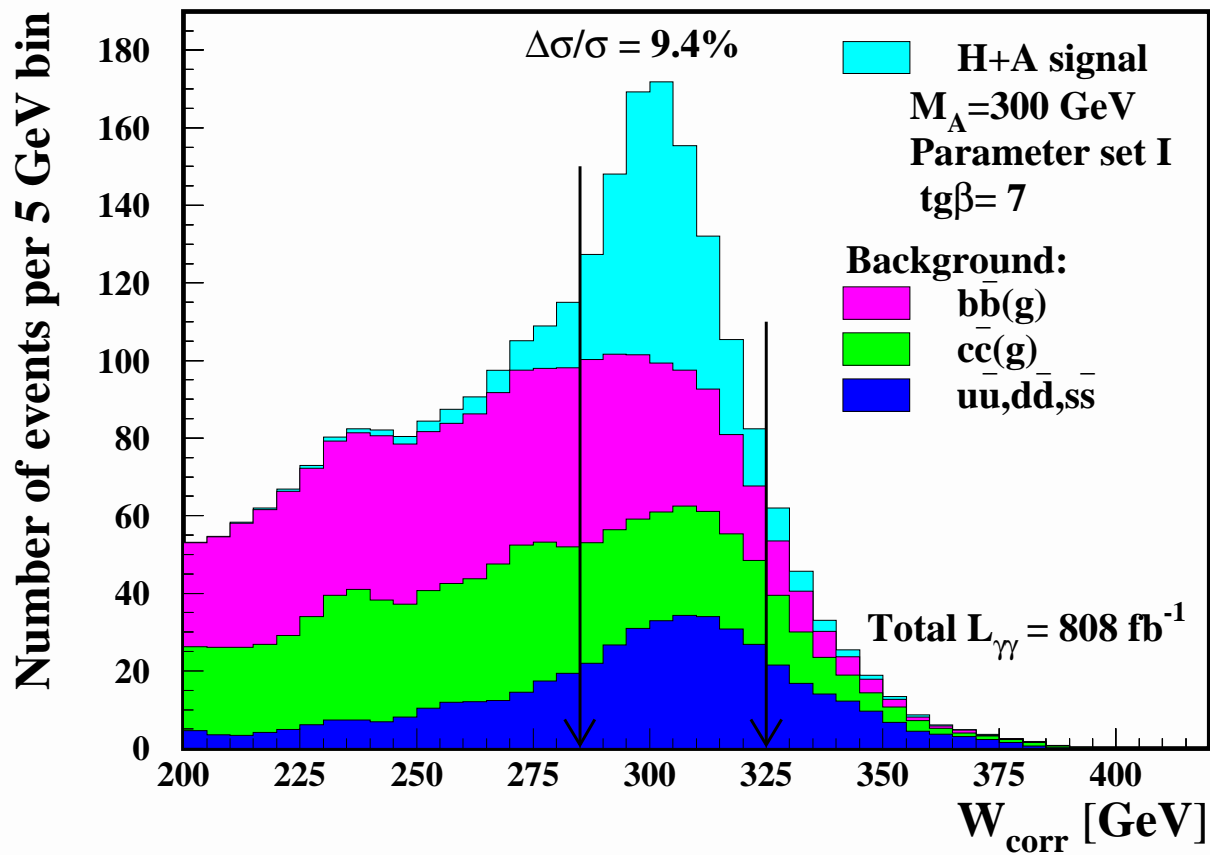
$$M_A = 300 \text{ GeV}$$

Without OE, without $\gamma\gamma \rightarrow W^+W^-$



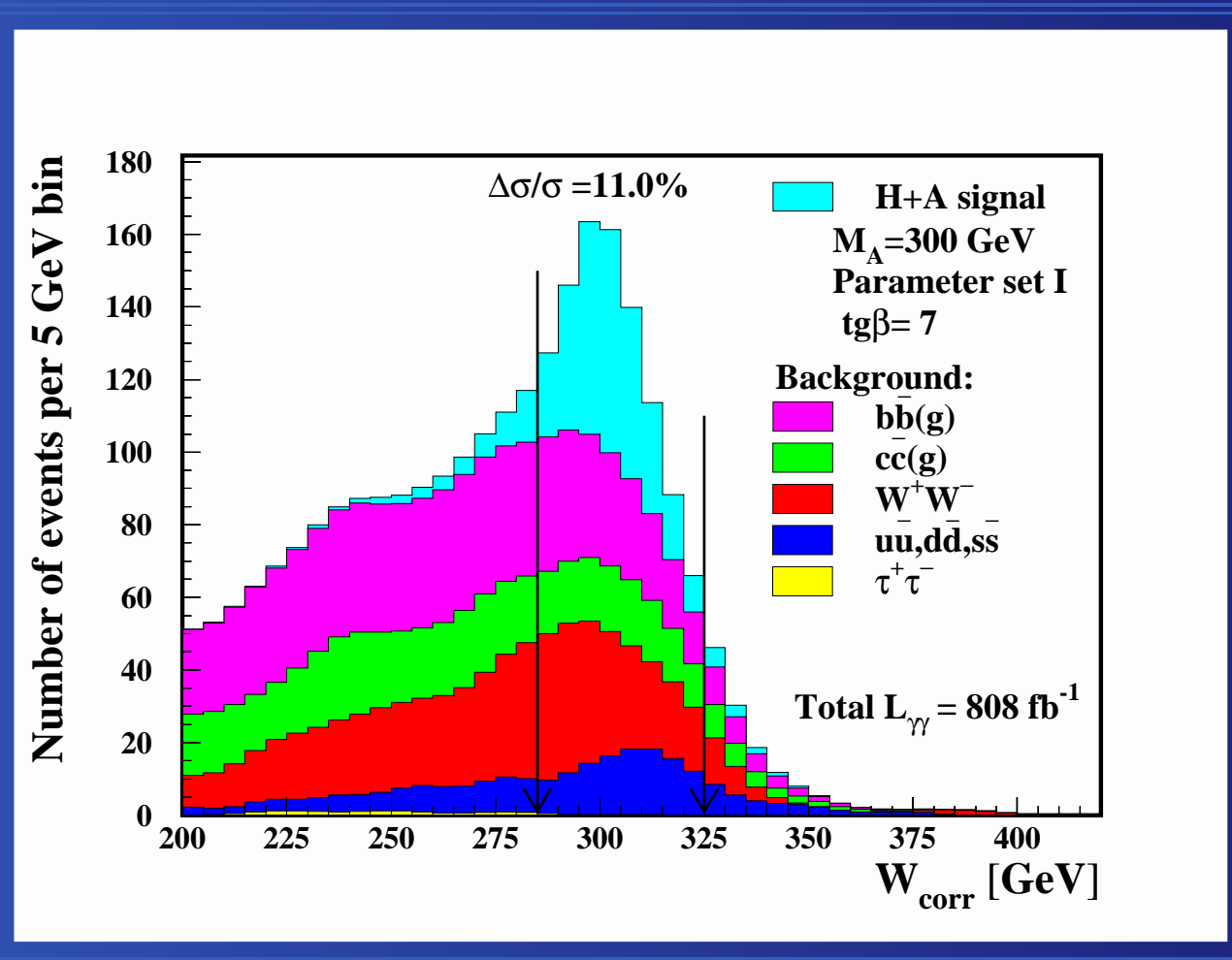
$$M_A = 300 \text{ GeV}$$

With OE, without $\gamma\gamma \rightarrow W^+W^-$



$$M_A = 300 \text{ GeV}$$

Final results



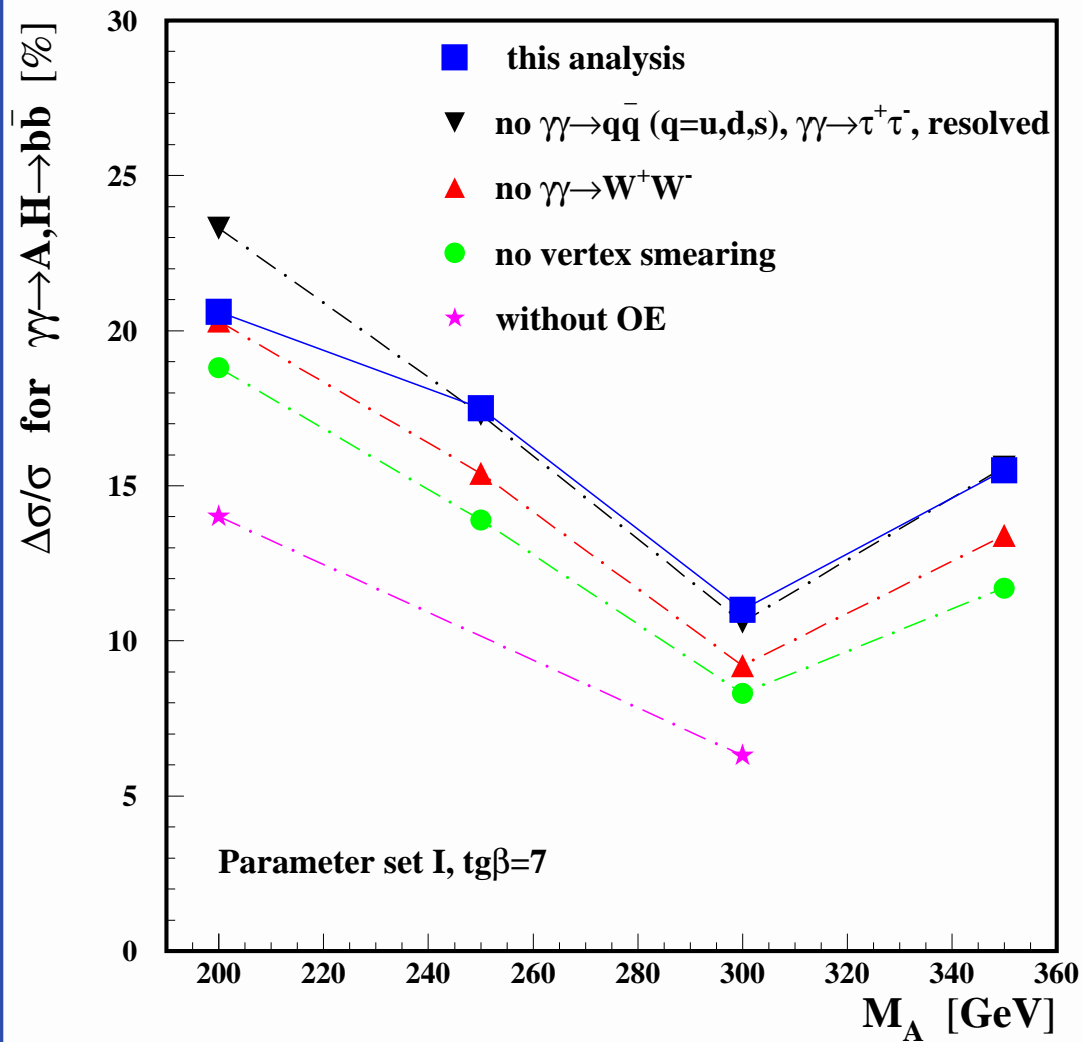
$\gamma\gamma \rightarrow \text{hadrons}$ (resolved) as a separate contribution – inefficient generation

\Rightarrow we estimate number of events in the mass window

$\gamma\gamma \rightarrow q\bar{q}$ ($q = u, d, s$) (unpolarized cross sec.) is overestimated
 but compensates the lack of *resolved* contribution

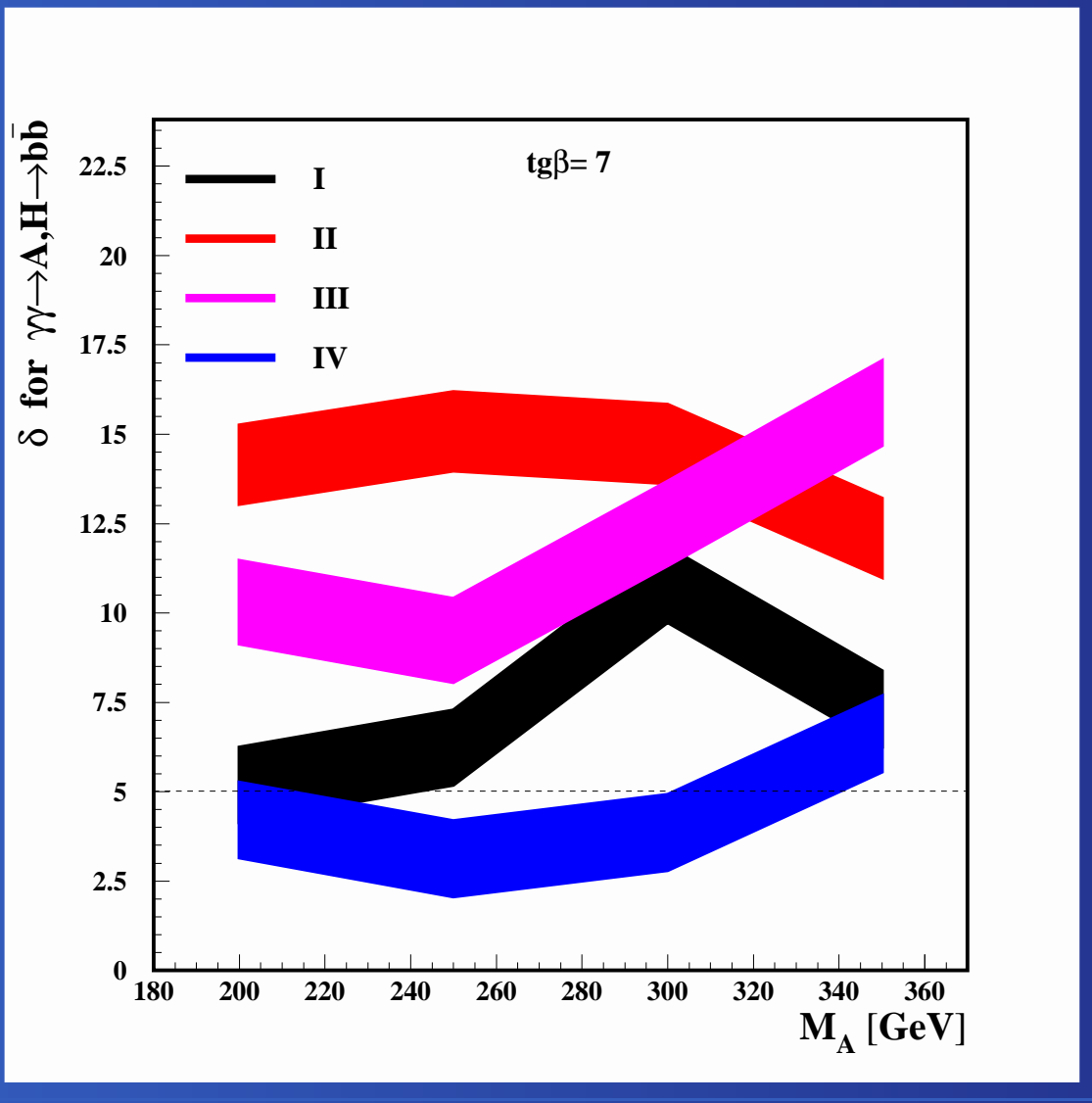


MSSM summary



Photon Collider reach

Significance for $\gamma\gamma \rightarrow A, H \rightarrow b\bar{b}$

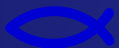


Symbol	μ [GeV]	$A_{\tilde{f}}$ [GeV]
I	200	1500
II	-150	1500
III	-200	1500
IV	300	2450

Significance

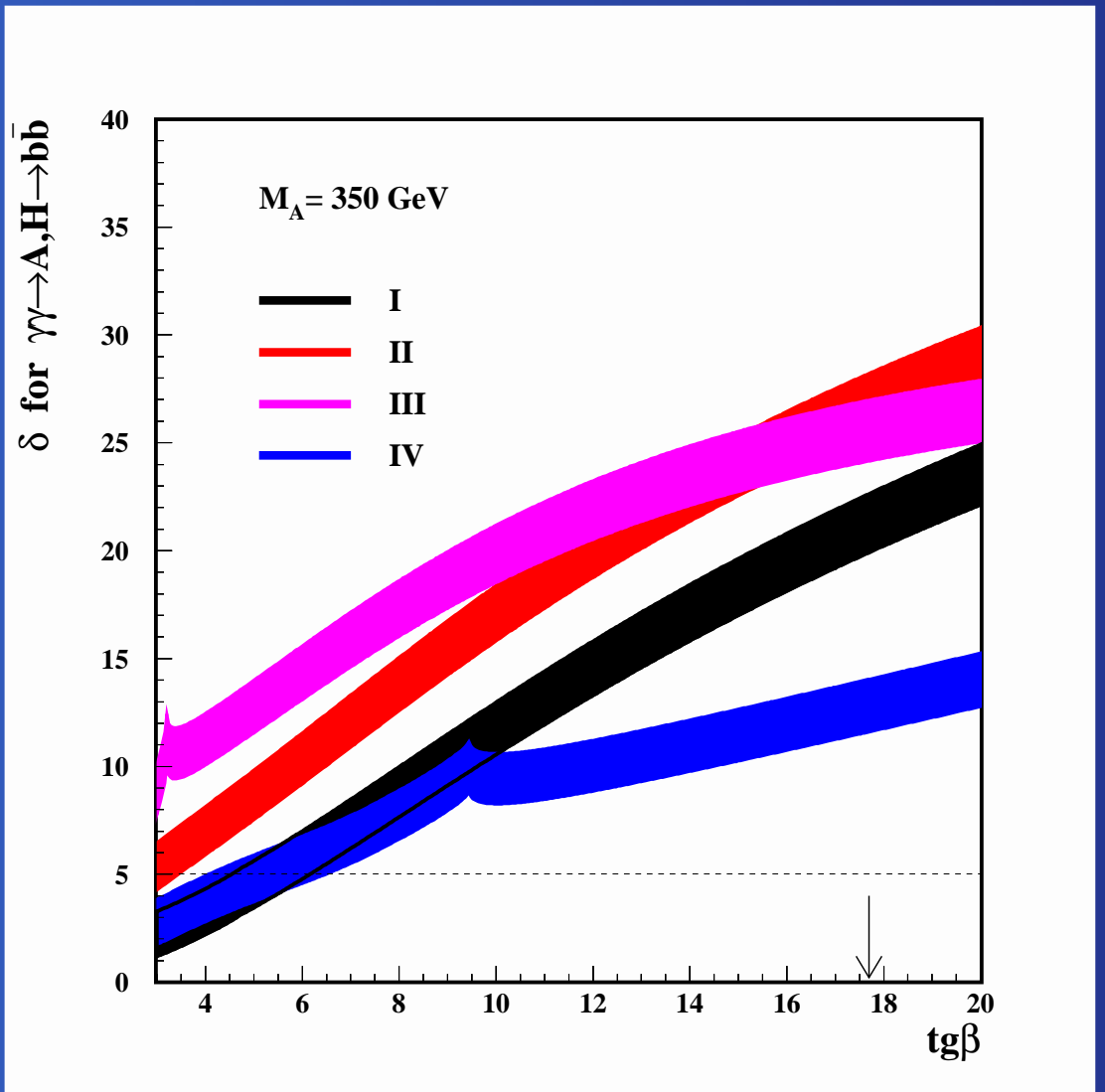
$$\delta = \frac{\mu_S}{\sqrt{\mu_B}} \pm \sqrt{1 + \frac{\mu_S}{\mu_B}}$$

The band widths indicate the level of possible statistical fluctuations



Photon Collider reach

Significance for $\gamma\gamma \rightarrow A, H \rightarrow b\bar{b}$



Symbol	μ [GeV]	$A_{\tilde{f}}$ [GeV]
I	200	1500
II	-150	1500
III	-200	1500
IV	300	2450

Significance

$$\delta = \frac{\mu_S}{\sqrt{\mu_B}} \pm \sqrt{1 + \frac{\mu_S}{\mu_B}}$$

The band widths indicate the level of possible statistical fluctuations

Arrow – lower limit of the LHC discovery region.



MKSZ: Results

Cross sections of $\gamma\gamma \rightarrow A, H \rightarrow b\bar{b}$ and $\gamma\gamma \rightarrow b\bar{b}$ processes

M. Mühlleitner, M. Krämer, M. Spira,
P. Zerwas, Phys. Lett. B 508 (2001) 311.

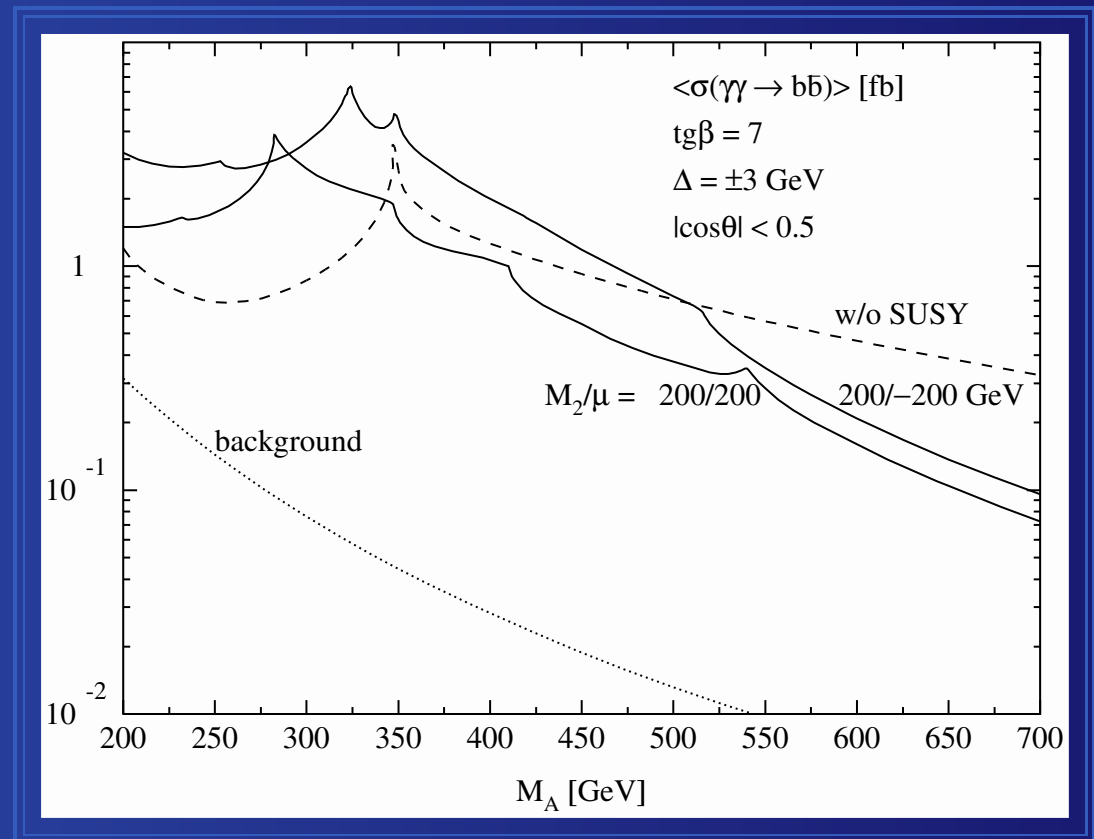
Considered MSSM parameter sets

μ [GeV]	M_2 [GeV]	$A_{\tilde{f}}$ [GeV]
200	200	0
-200	200	0

Also the limit of vanishing SUSY-particle contributions considered.

Results for $M_A = 200-700$ GeV
and for $\tan\beta = 7$

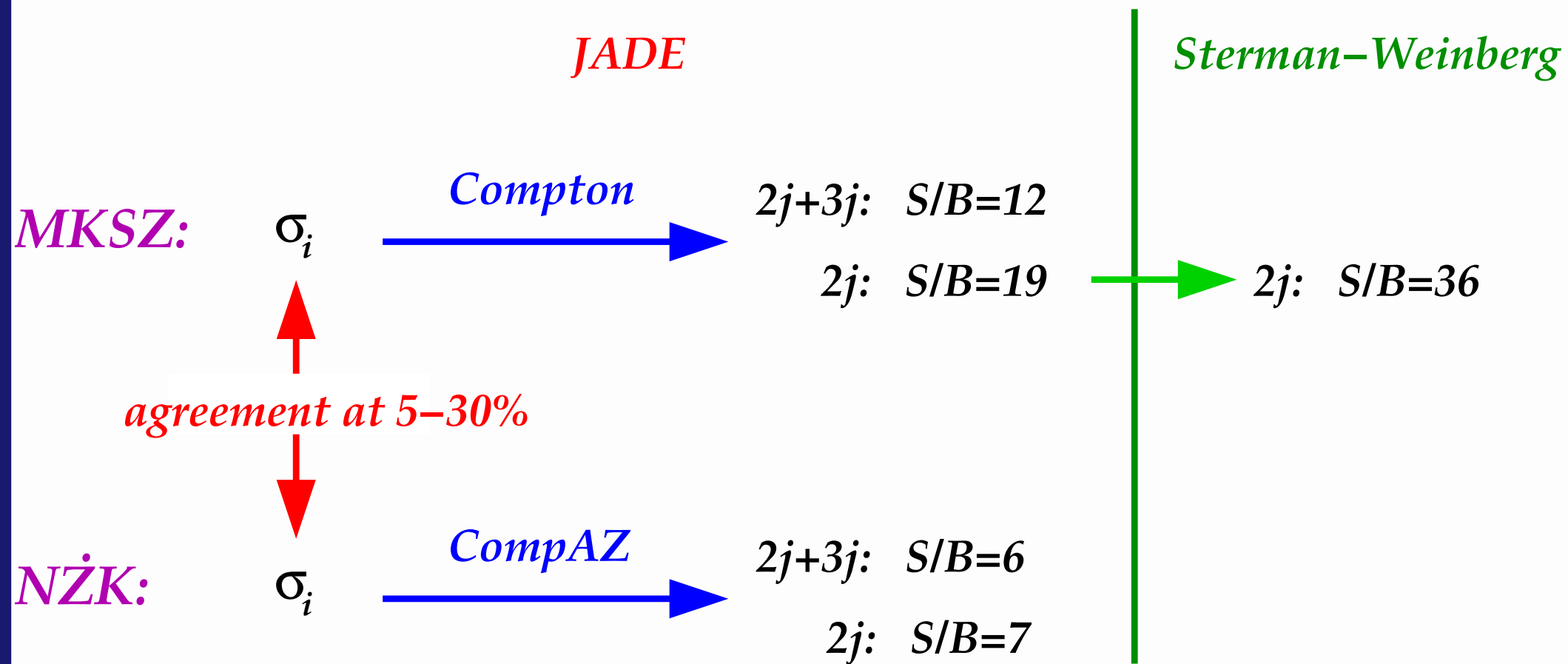
Results for $M_A = 200-700$ GeV



Average cross sections in the invariant mass window ± 3 GeV.

$M_A = 300$ GeV, $\mu = 200$ GeV $\Rightarrow S/B \approx 36$

Summary of comparison



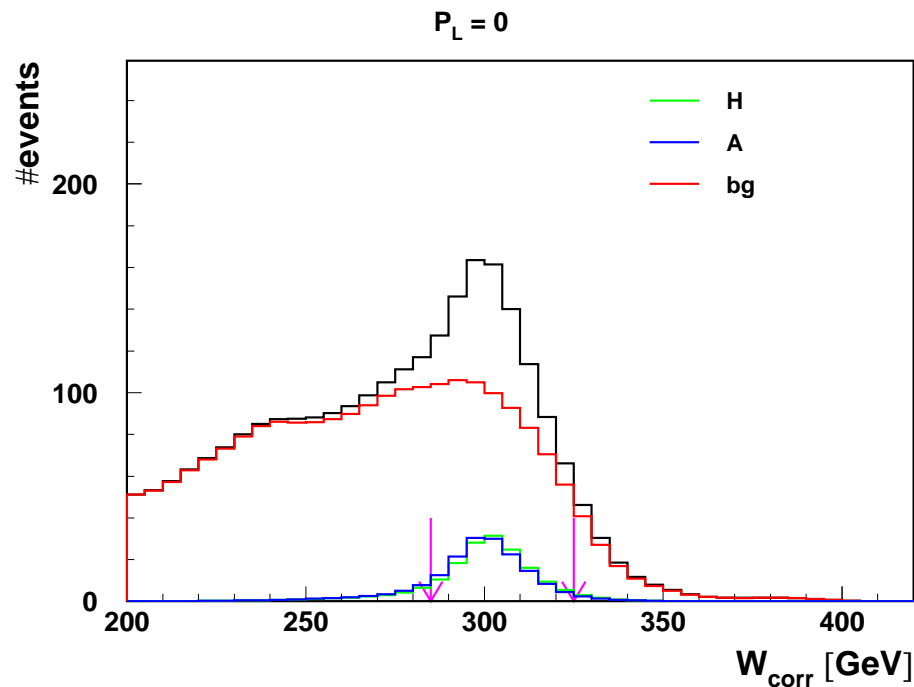
More detailed description: hep-ph/0612369

Higgs CP properties

MSSM results

Results for $M_A = 300$ GeV

Circular laser polarization, $P_C = 100\%$



We can not distinguish between H and A
 \Rightarrow measurement of

$$\sigma_{\text{tot}} = \sigma_H + \sigma_A$$

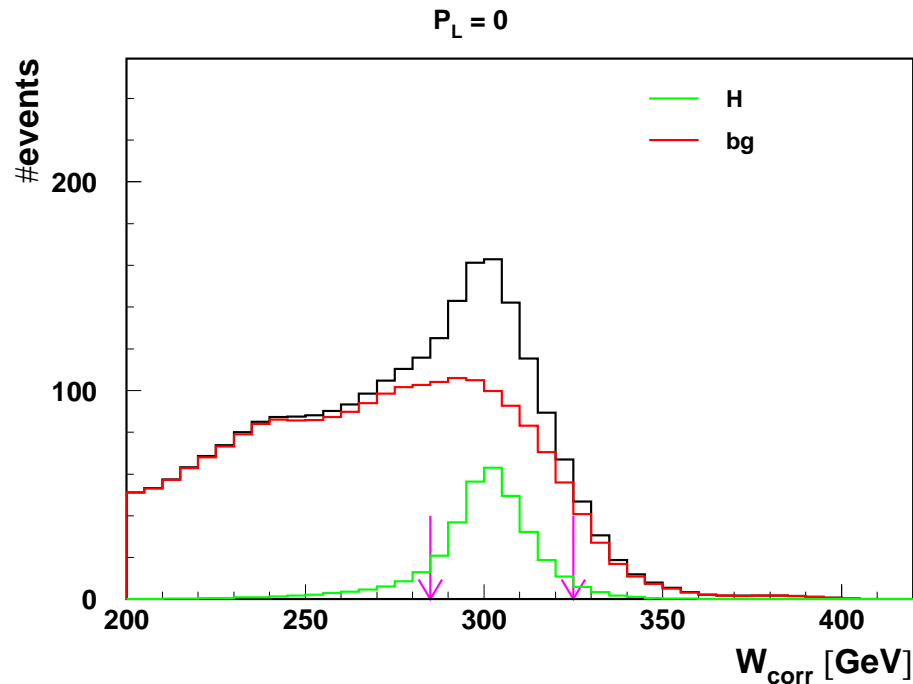
Corrected invariant mass distributions

Higgs CP properties

MSSM results

Results for $M_A = 300$ GeV

Circular laser polarization, $P_C = 100\%$



We can not distinguish between H and A
 \Rightarrow measurement of

$$\sigma_{\text{tot}} = \sigma_H + \sigma_A$$

\Rightarrow Need for linear photon polarization

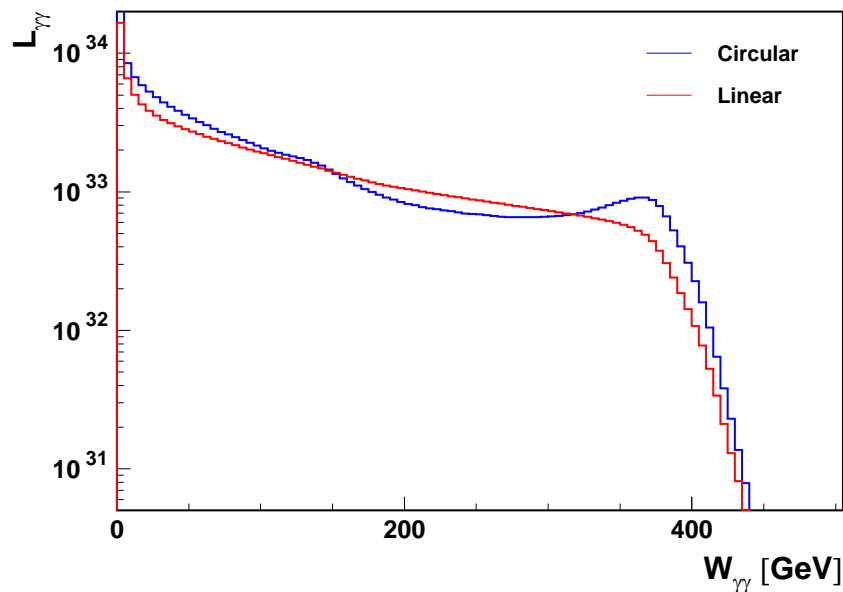
Corrected invariant mass distributions

New results
with all backgrounds included!

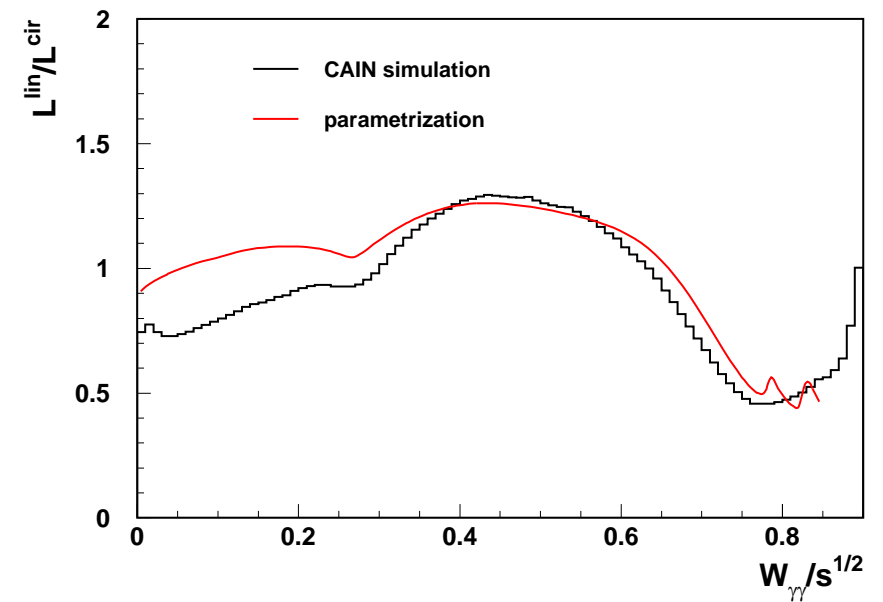
Linear polarization

CAIN simulation

$\gamma\gamma$ luminosity spectra for **circular** and **linear** laser beam polarization



Ratio of $\gamma\gamma$ luminosities

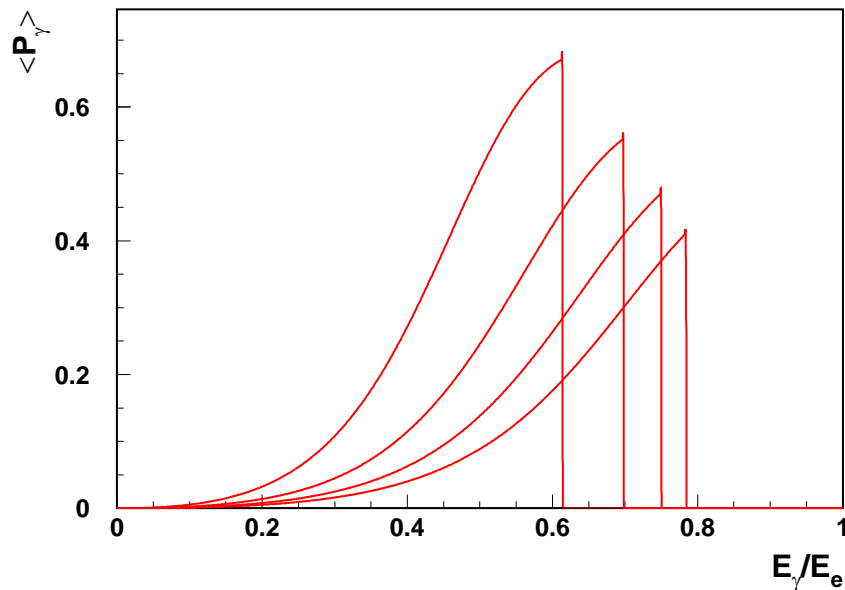


CompAZ gives proper description of the spectra modification

Linear polarization

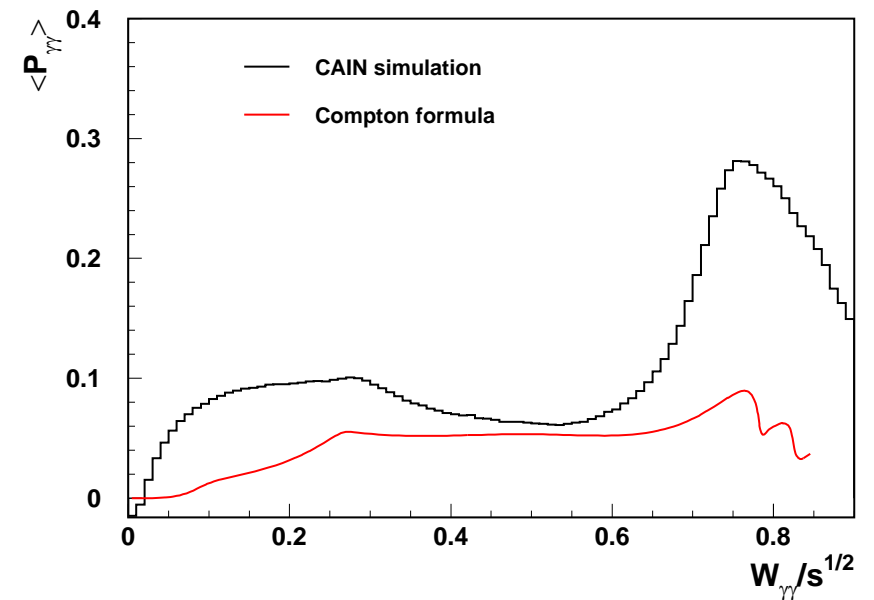
CAIN simulation

Expected photon polarization
from LO Compton process



for $E_e = 100, 150, 200$ and 250 GeV

$\gamma\gamma$ polarization from CAIN for 250 GeV



CompAZ fails to describe polarization !

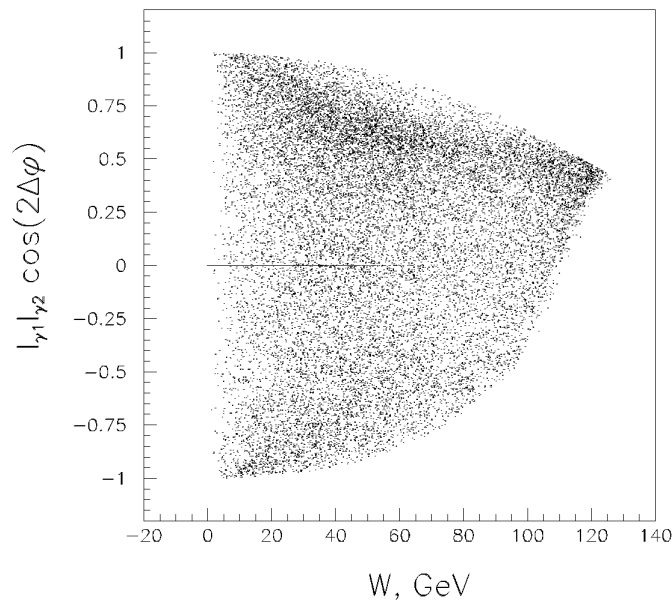
Linear polarization

Angular correlations

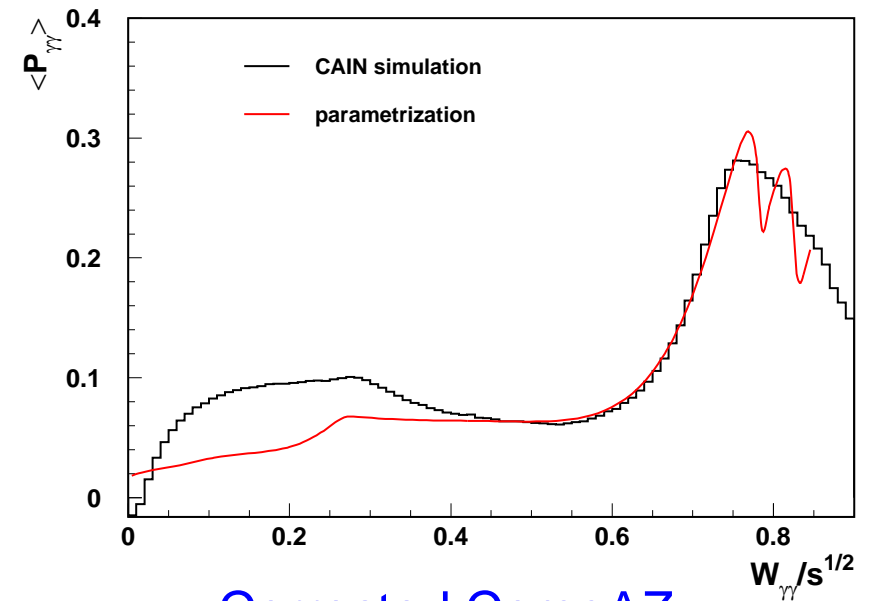
As pointed out by V.Telnov (“Nontrivial effects in linear polarization at photon colliders”, ECFA workshop, Montpellier, November 2003) there are large correlations between photon polarization and scattering direction. In collision of two very thin beams:

$$\langle P_{\gamma_1} P_{\gamma_2} \rangle \gg \langle P_{\gamma_1} \rangle \cdot \langle P_{\gamma_2} \rangle$$

V.Telnov



Average $\gamma\gamma$ polarization from CAIN

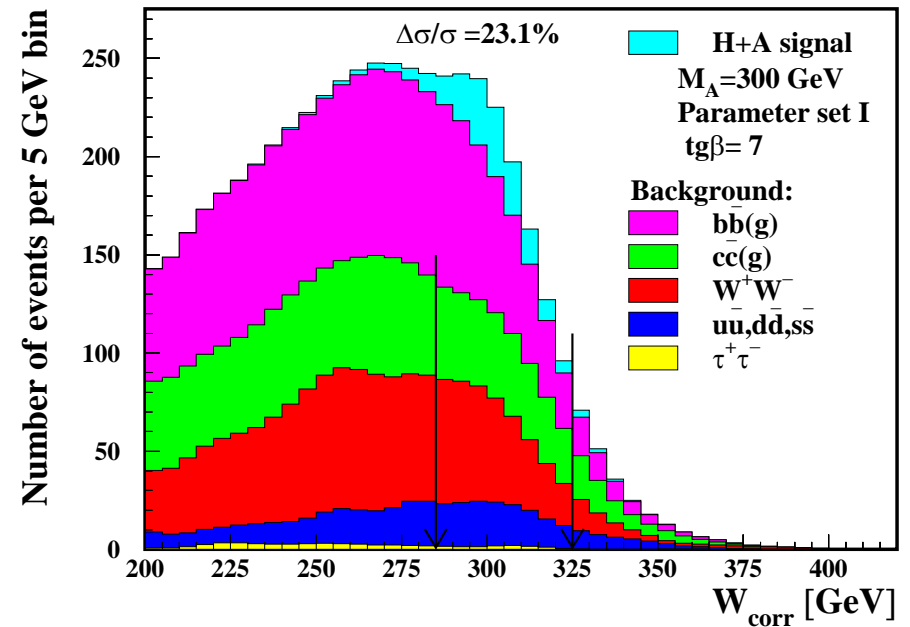
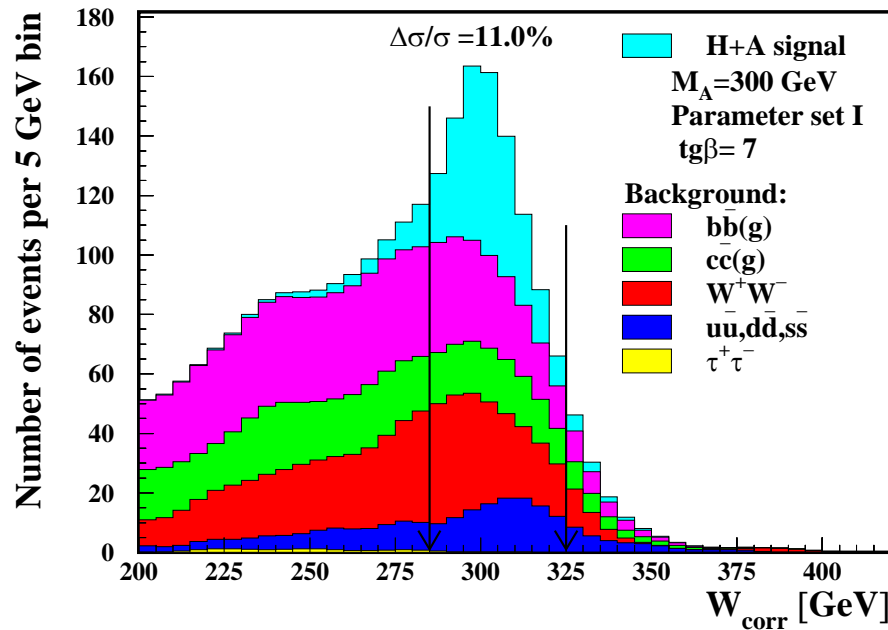


Corrected CompAZ

Results

$M_A = 300$ GeV

Linear laser polarization, $P_L = 100\%$



Lower luminosity at M_A , lower $J_z = 0$ contribution \Rightarrow signal down by 40%

Higher $J_z = 2$ contribution \Rightarrow no background suppression \Rightarrow background up by 90%

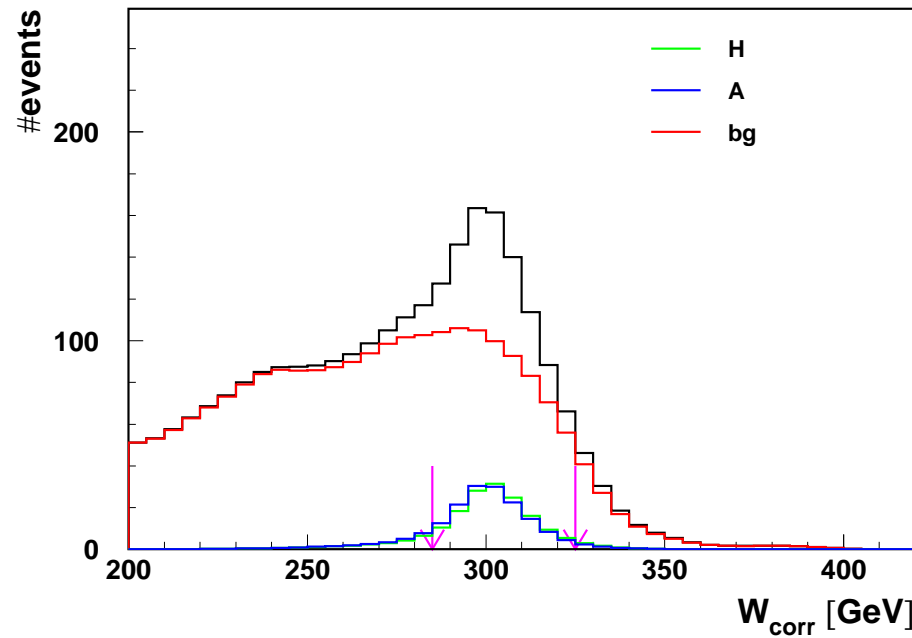
Selection cuts differ !!!

Results

$$\underline{M_A = 300 \text{ GeV}}$$

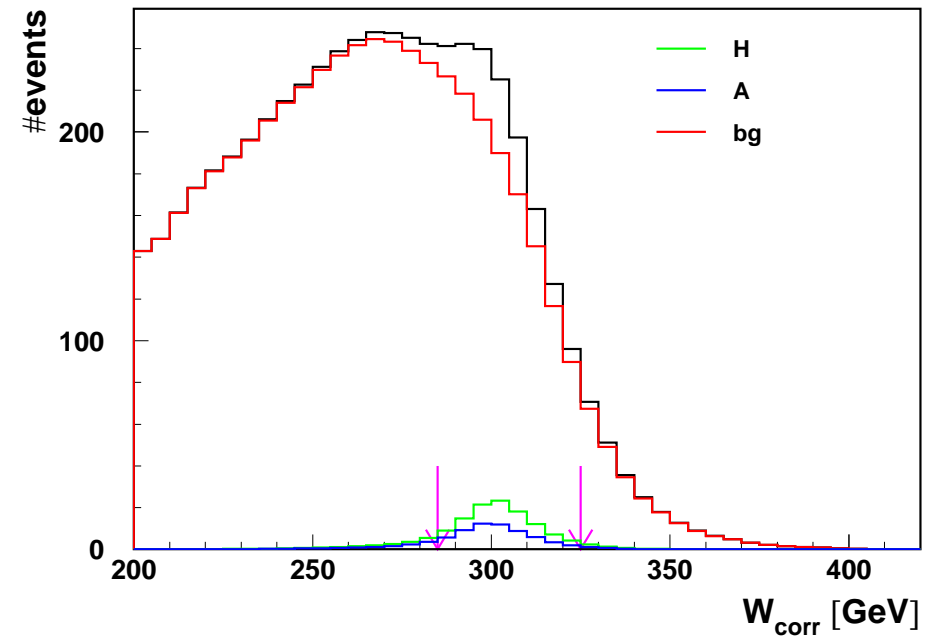
Measurements start to be sensitive to the Higgs boson(s) **CP properties**.

Circular laser polarization, $P_C = 100\%$
 $P_L = 0$



$$f \approx 0.5$$

Linear laser polarization, $P_L = 100\%$
 $P_C = 0$



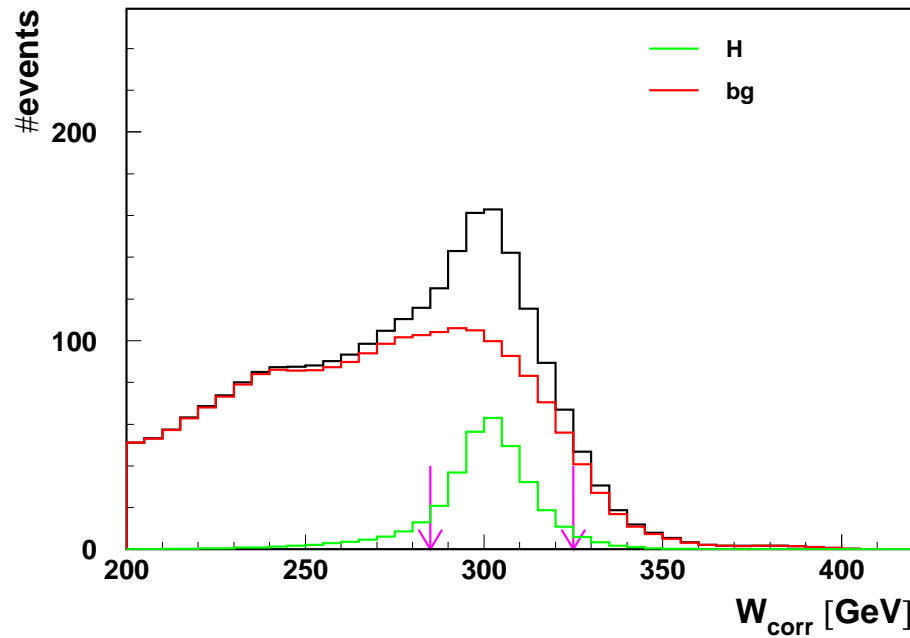
$$\sigma_H = f \cdot \sigma_o, \quad \sigma_A = (1 - f) \cdot \sigma_o \quad \text{where } \sigma_o \equiv \sigma_{H+A}(P_L = 0)$$

Results

$$\underline{M_A = 300 \text{ GeV}}$$

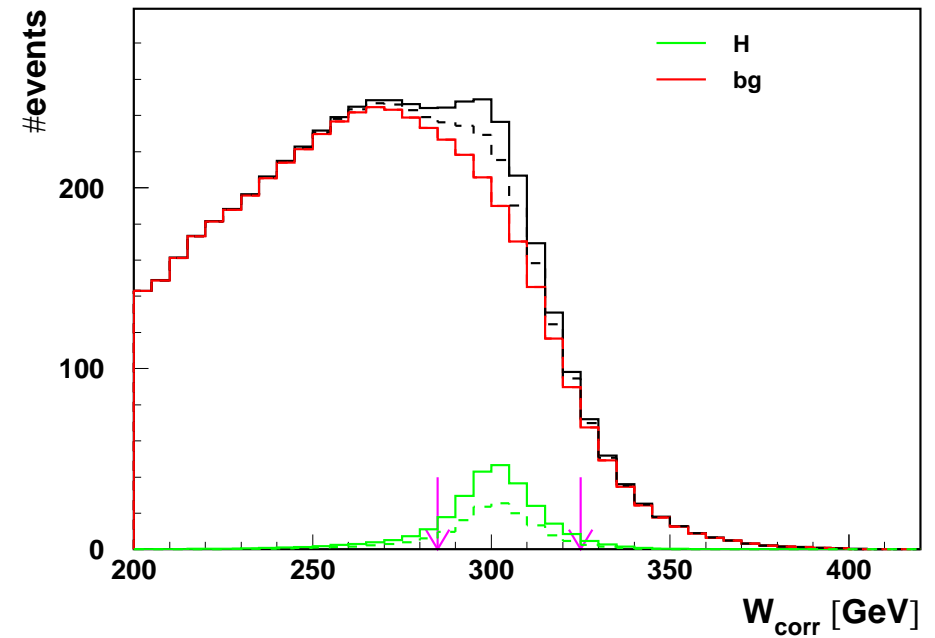
Measurements start to be sensitive to the Higgs boson(s) **CP properties**.

Circular laser polarization, $P_C = 100\%$
 $P_L = 0$



$$f = 1$$

Linear laser polarization, $P_L = 100\%$
 $P_L = 1$



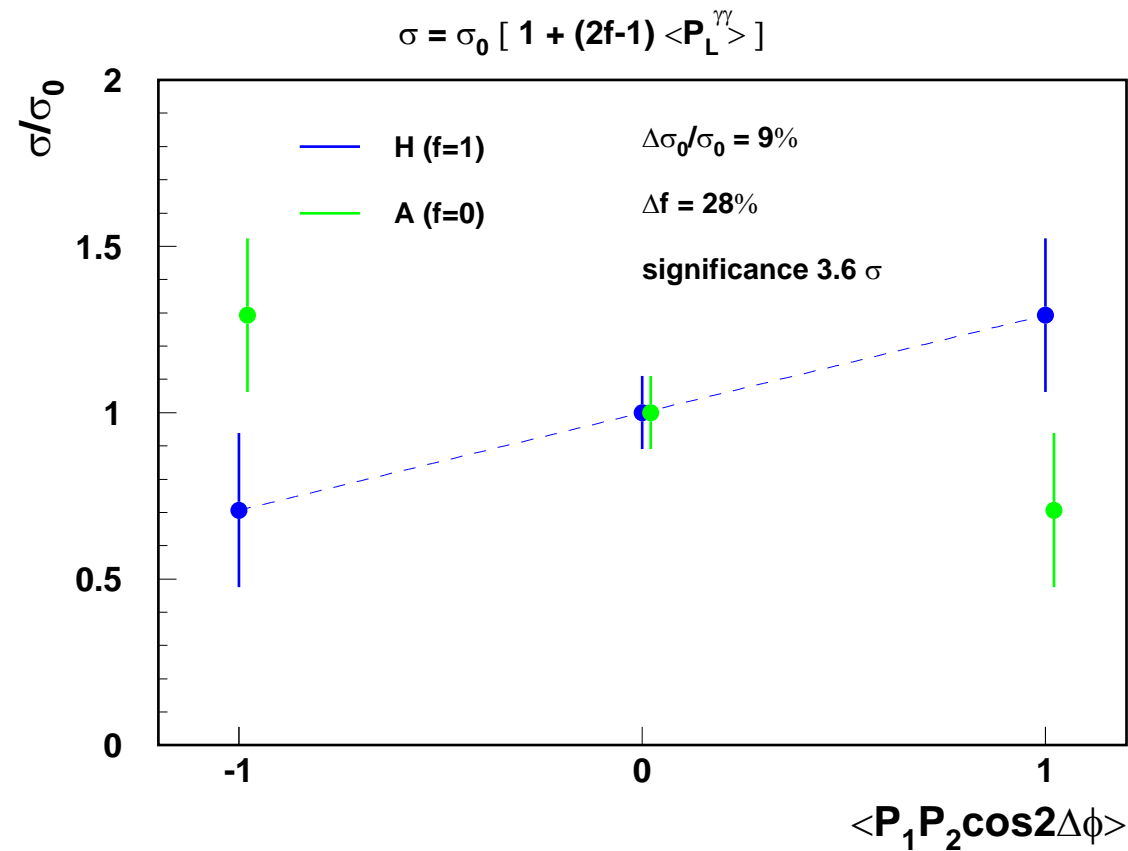
parallel (solid), perpendicular (dashed)

$$\sigma_H = f \cdot \sigma_o, \quad \sigma_A = (1 - f) \cdot \sigma_o \quad \text{where } \sigma_o \equiv \sigma_{H+A}(P_L = 0)$$

Results

$M_A = 300$ GeV

Results expected after 3×1 years of PLC running



σ_0 corresponding to MSSM parameter set I

Higgs CP

Generic model

Model with a generic tensor couplings of a Higgs boson \mathcal{H} , to ZZ and W^+W^- :

$$g_{\mathcal{H}ZZ} = ig \frac{M_Z}{\cos \theta_W} \left(\lambda_H \cdot g^{\mu\nu} + \lambda_A \cdot \varepsilon^{\mu\nu\rho\sigma} \frac{(p_1 + p_2)_\rho (p_1 - p_2)_\sigma}{M_Z^2} \right)$$
$$g_{\mathcal{H}WW} = ig M_W \left(\lambda_H \cdot g^{\mu\nu} + \lambda_A \cdot \varepsilon^{\mu\nu\rho\sigma} \frac{(p_1 + p_2)_\rho (p_1 - p_2)_\sigma}{M_W^2} \right)$$

with: $\lambda_H = \lambda \cdot \cos \Phi_{HA}$ $\lambda_A = \lambda \cdot \sin \Phi_{HA}$

Standard Model (scalar) couplings are reproduced for $\Phi_{HA} = 0$ ($\lambda_H = 1$ and $\lambda_A = 0$).

Pseudoscalar Higgs boson corresponds to $\lambda_H = 0$ and $\Phi_{HA} = \frac{\pi}{2}$ $\lambda_A = 1$.

We consider **small CP violation** (small deviations from SM), i.e. $|\Phi_{HA}| \ll 1$

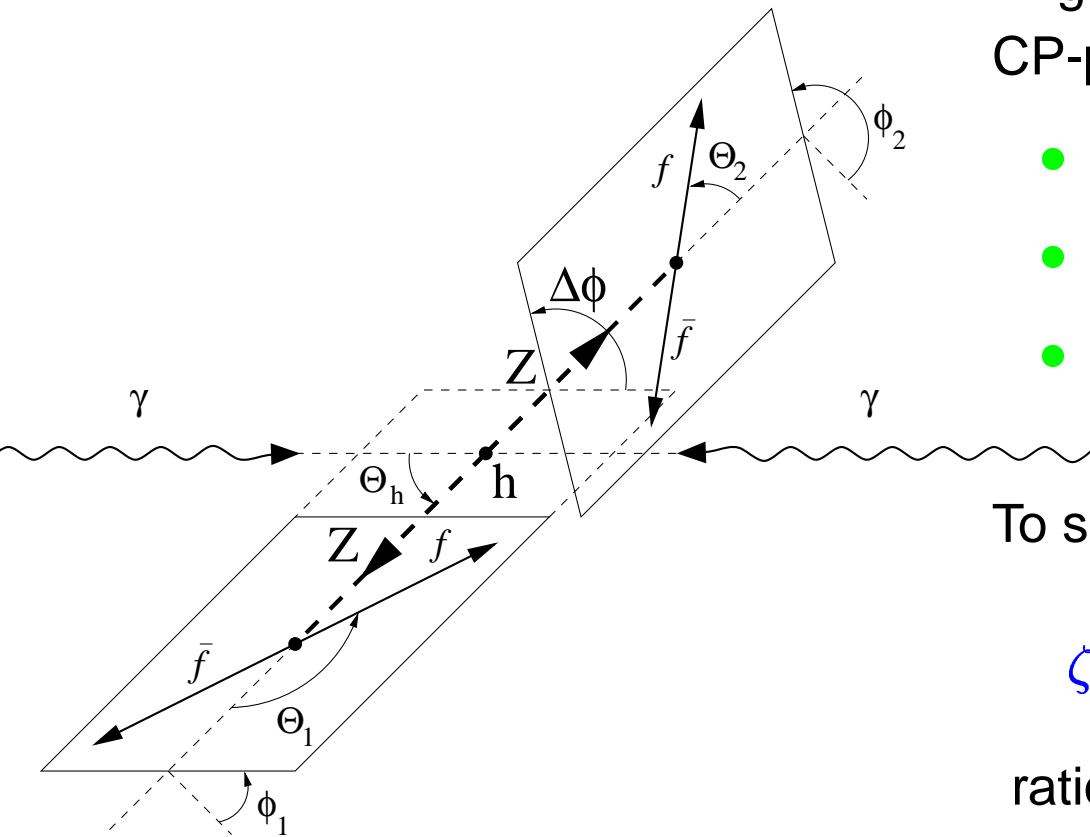
Model: S.Y. Choi, D.J. Miller, M.M. Mühlleitner and P.M. Zerwas, hep-ph/0210077;
D.J. Miller, S.Y. Choi, B. Eberle, M.M. Mühlleitner and P.M. Zerwas, Phys. Lett. B505 (2001) 149;
D.J. Miller, *Spin and Parity in the HZZ vertex*, ECFA/DESY meeting, Prague, November 2002.

Higgs CP from $\mathcal{H} \rightarrow \tau^+\tau^-$: K. Desch, A. Imhof, Z. Was, M. Worek, hep-ph/0307331;
K. Desch, Z. Was, M. Worek, Eur.Phys.J.C29 (2003) 491, hep-ph/0302046.

Higgs CP from $\mathcal{H} \rightarrow t\bar{t}$: E. Asakawa, K. Hagiwara, hep-ph/0305323.

Generic model

Angular distributions



Angular variables used in the analysis of higgs CP-properties:

- higgs decay angle angle Θ_h
- polar angles Θ_1 and Θ_2
- angle between two Z/W decay planes,

$$\Delta\phi = \phi_2 - \phi_1$$

To simplify the analysis, we introduce

$$\zeta = \frac{\sin^2 \Theta_1 \cdot \sin^2 \Theta_2}{(1 + \cos^2 \Theta_1) \cdot (1 + \cos^2 \Theta_2)}$$

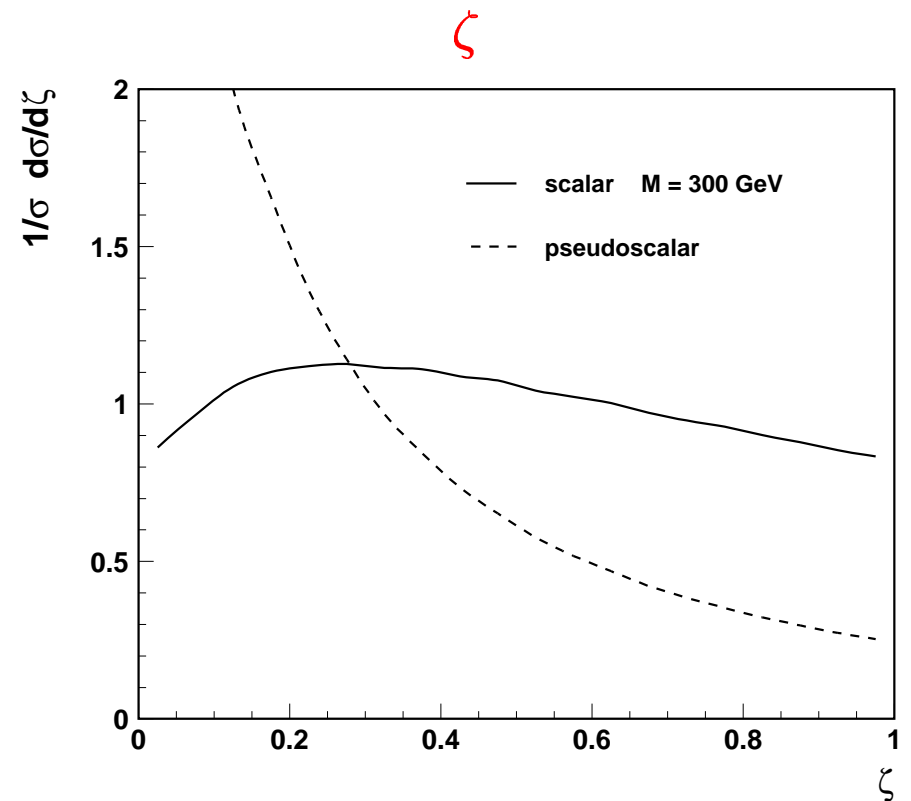
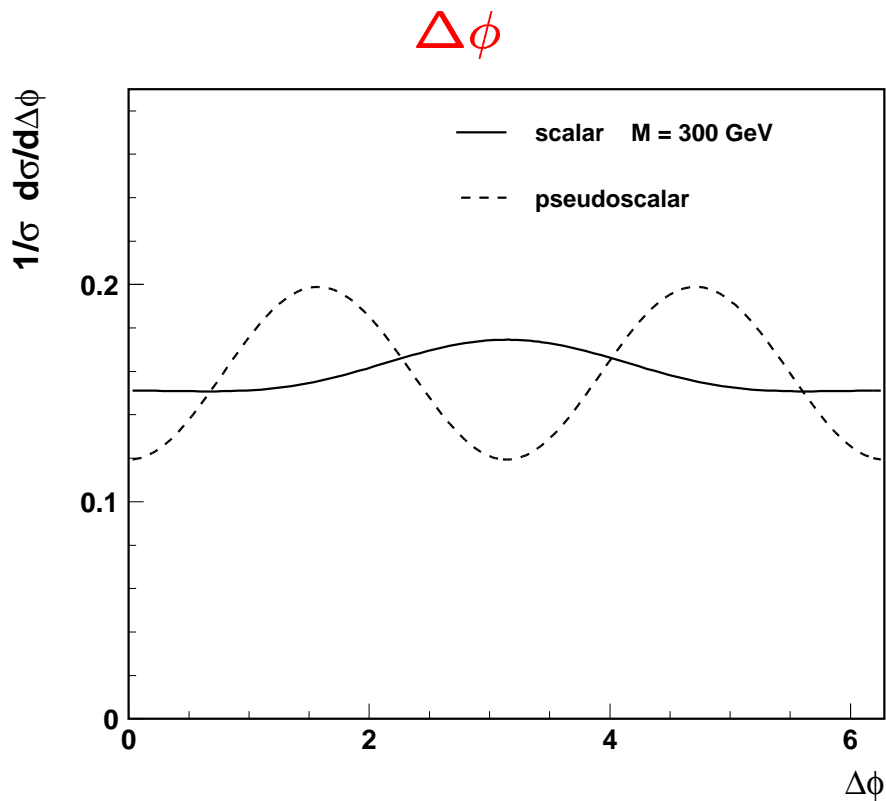
ratio of the distributions expected for a scalar and a pseudoscalar higgs (for $M_h \gg M_Z$).

All polar angles are calculated in the rest frame of the decaying particle.

Generic model

Angular distributions

Normalized angular distributions expected for **scalar** and **pseudoscalar** higgs, for $\mathcal{H} \rightarrow ZZ \rightarrow l^+l^-jj$ $M_{\mathcal{H}} = 300$ GeV.



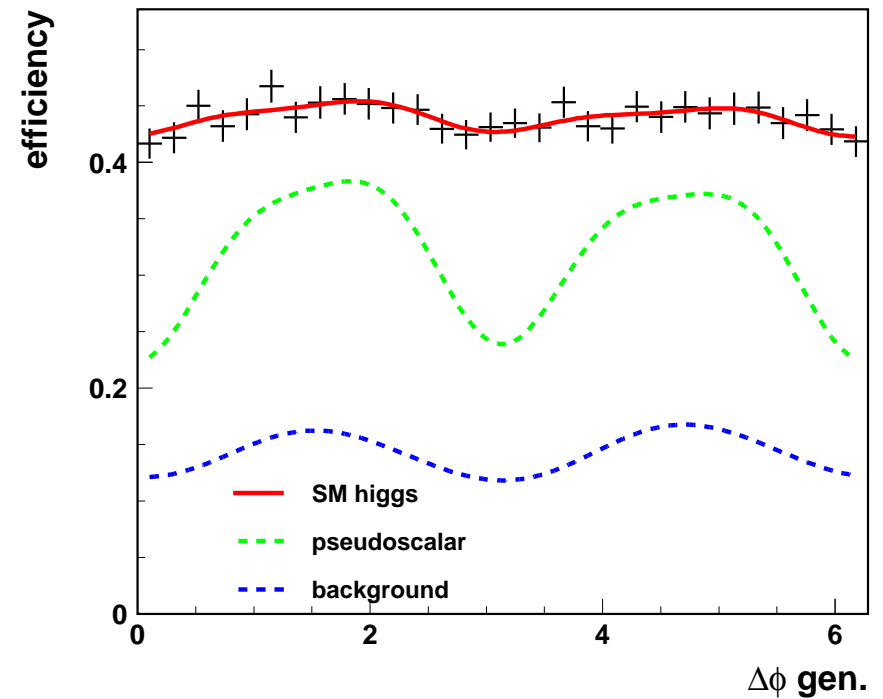
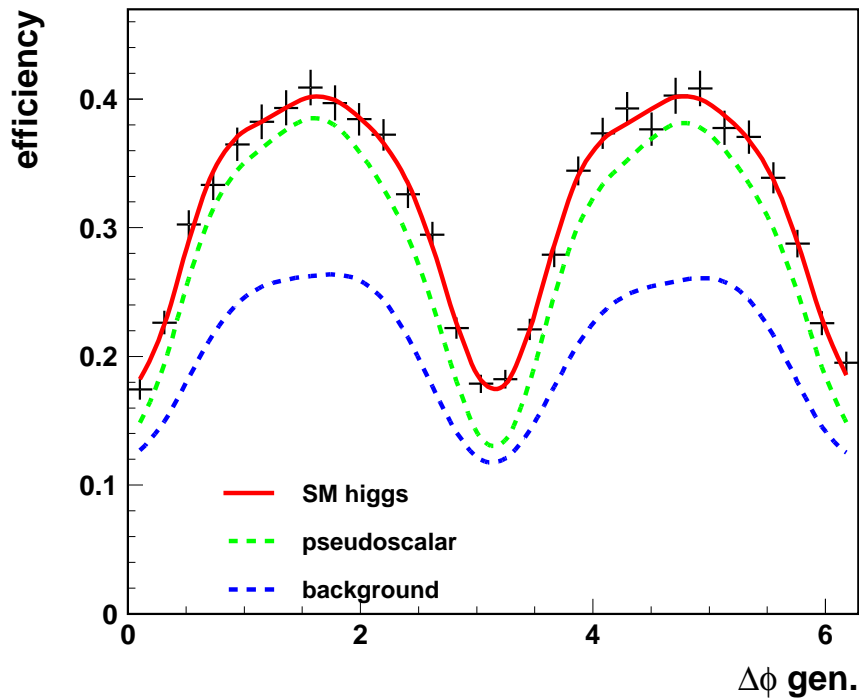
Both distributions clearly distinguish between decays of scalar and pseudoscalar higgs.

Generic model

Nonuniformity of selection efficiency in $\Delta\phi$ largest for small m_h

$m_h = 200$ GeV, $\sqrt{s_{ee}} = 305$ GeV

$m_h = 300$ GeV, $\sqrt{s_{ee}} = 418$ GeV

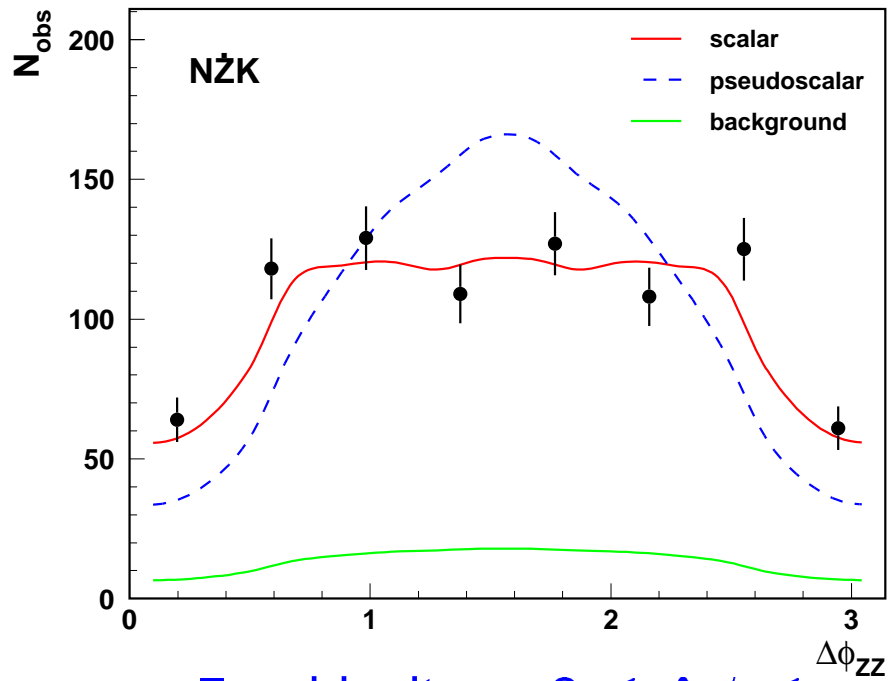


Effect much stronger for background events and pseudoscalar higgs
due to different $\cos\theta_{j,l}$ distribution

Generic model

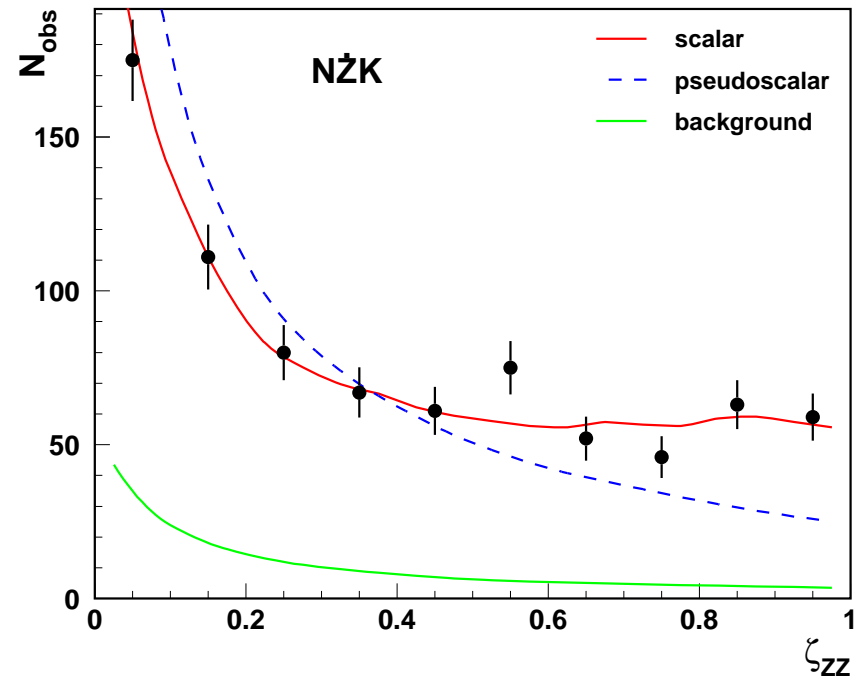
Measured $\Delta\phi$ and ζ distributions for $h \rightarrow ZZ \rightarrow q\bar{q}l^+l^-$ $m_h = 200$ GeV
 after 1 year of PC running at $\sqrt{s_{ee}}=305$ GeV, $\mathcal{L} = 610$ fb $^{-1}$
 $\Rightarrow \sim 675$ reconstructed SM higgs events expected + 145 ZZ background events

Measured $\Delta\phi_{ZZ}$ distribution:



$q \leftrightarrow \bar{q}$ ambiguity $\Rightarrow 0 \leq \Delta\phi \leq \pi$

Measured ζ_{ZZ} distribution:



pseudoscalar normalized to the same number of events

Generic model

Sensitivity

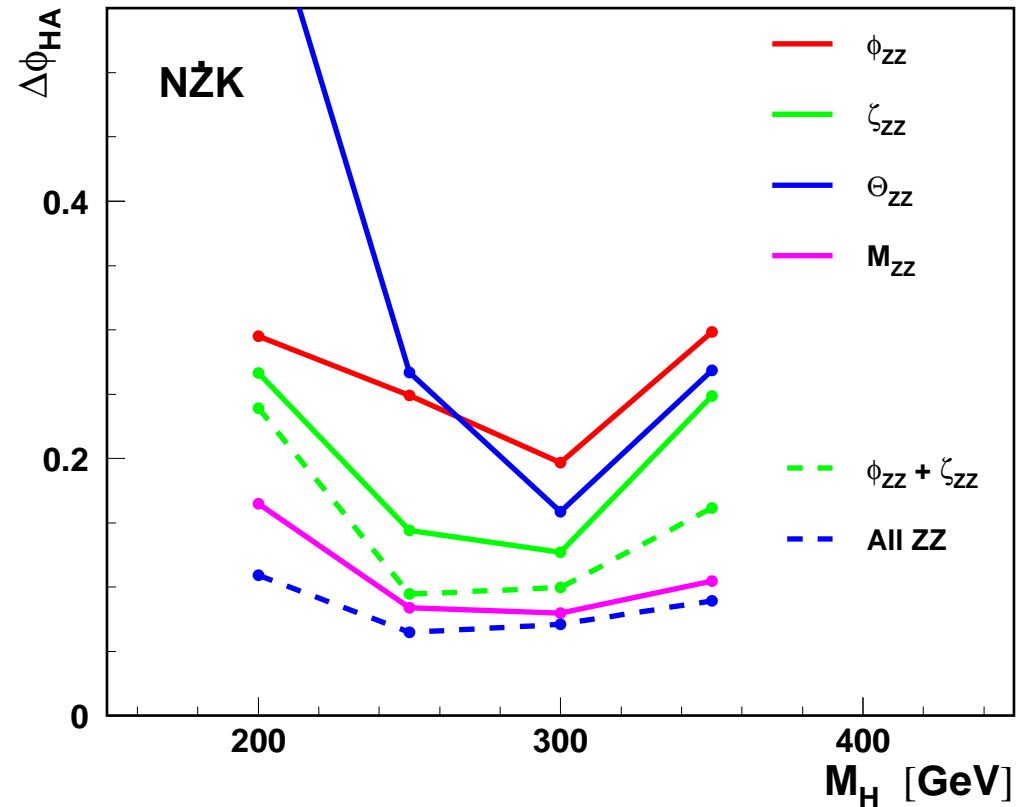
Statistical error on Φ_{HA}
from fits to different distributions \Rightarrow

Fits of two parameters:

Φ_{HA} + normalization

We assume here:

$$\begin{aligned}\Gamma_{\gamma\gamma} &= \Gamma_{\gamma\gamma}^{SM} \\ \phi_{\gamma\gamma} &= \phi_{\gamma\gamma}^{SM} \\ \lambda &= \lambda^{SM} \equiv 1\end{aligned}$$



In the **general case** We can not assume that $\Gamma_{\gamma\gamma}$, $\phi_{\gamma\gamma}$ and λ are the same as in the SM
 \Rightarrow fit all distributions simultaneously to constrain all parameters

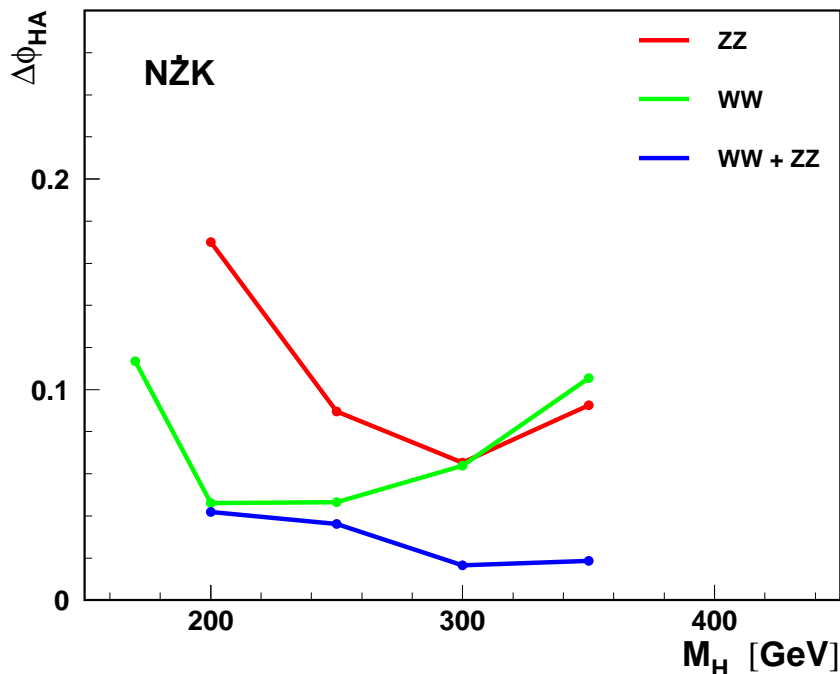
Generic model

Results

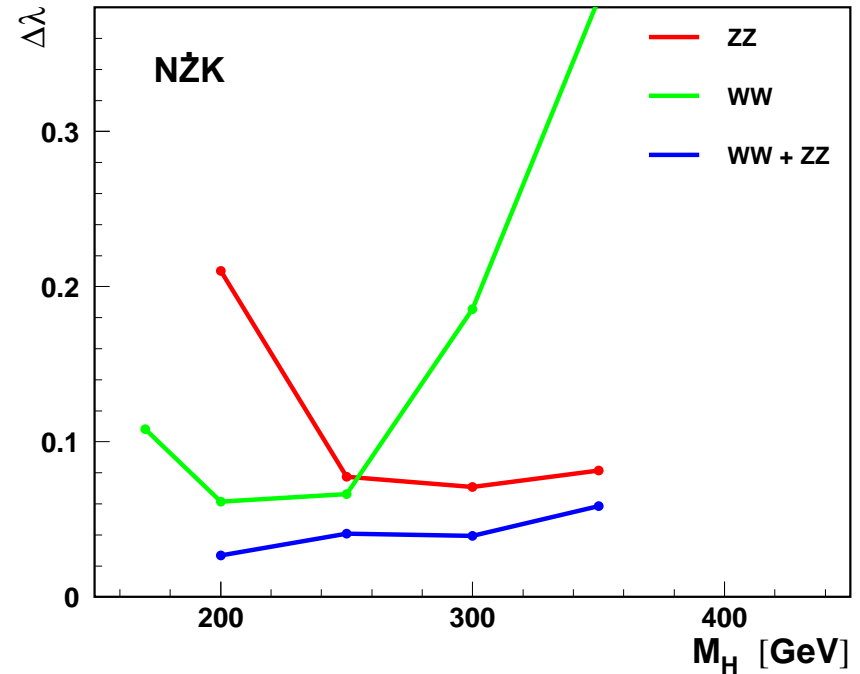
Combined measurement for W^+W^- and ZZ decay channels from simultaneously fit of $\Gamma_{\gamma\gamma}$, $\phi_{\gamma\gamma}$, λ and Φ_{HA} to all considered distributions

Measurement error for Higgs-boson couplings to vector bosons:

CP phase Φ_{HA}



Value λ



assuming SM-like couplings: $\lambda = 1$, $\Phi_{HA} = 0$

$W^+W^- \Rightarrow$ higher statistics, but huge background \Rightarrow large systematic uncertainties

CP conserving 2HDM (II)

Higgs boson couplings

Scalar Higgs bosons h and H
with basic couplings (relative to SM):

$$\chi_x = g_{\mathcal{H}xx} / g_{\mathcal{H}xx}^{SM} \quad \mathcal{H} = h, H, A$$

	h	H	A
χ_u	$\frac{\cos \alpha}{\sin \beta}$	$\frac{\sin \alpha}{\sin \beta}$	$-i \gamma_5 \frac{1}{\tan \beta}$
χ_d	$-\frac{\sin \alpha}{\cos \beta}$	$\frac{\cos \alpha}{\cos \beta}$	$-i \gamma_5 \tan \beta$
χ_V	$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$	0

For charged Higgs boson couplings
(loop contribution to $\Gamma_{\gamma\gamma}$) we set

$$M_{H^\pm} = 800 \text{ GeV} \quad \mu = 0$$

Higgs couplings are related by
“patter relation”

$$(\chi_V - \chi_d)(\chi_u - \chi_V) + \chi_V^2 = 1$$

I. F. Ginzburg, M. Krawczyk and P. Osland,
hep-ph/0101331

Instead of angles α and β use couplings
 χ_V and χ_u to parametrize cross sections

$$0 \leq \chi_V \leq 1$$

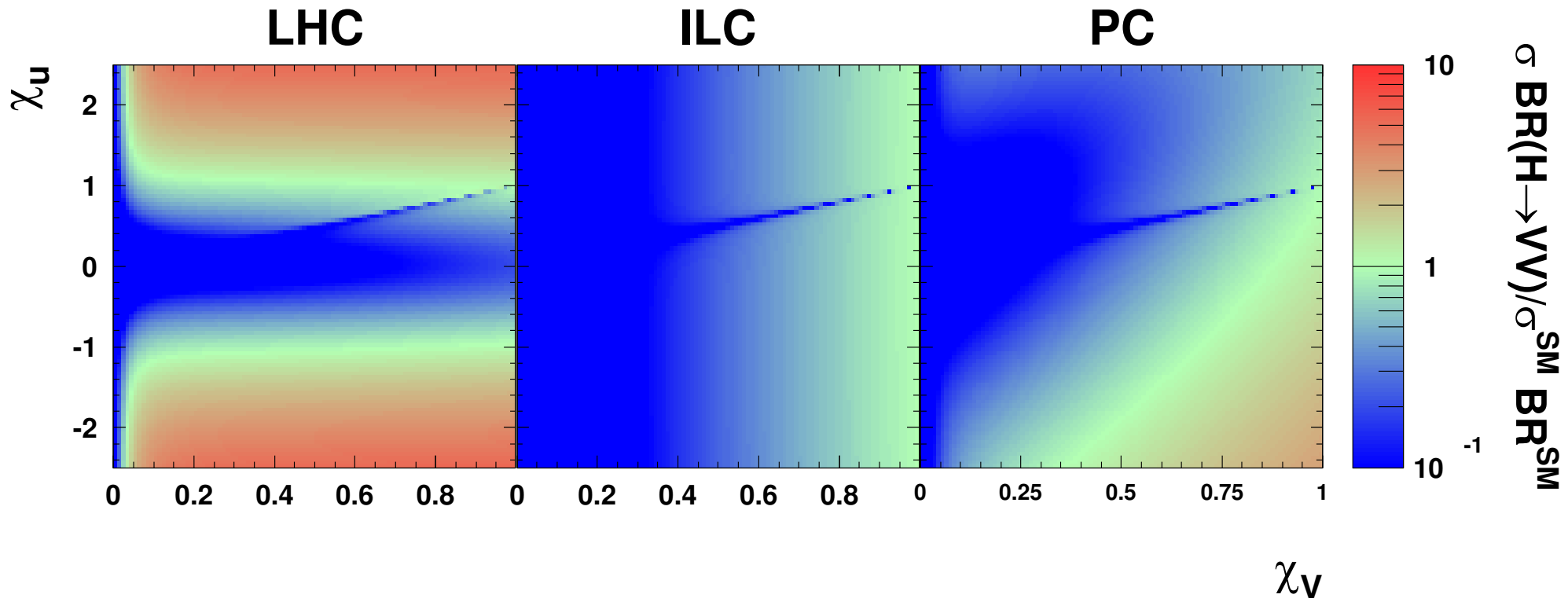
If we neglect H decays to h and A
(small) cross sections and BRs calculated for H are also valid for h

LHC ⊕ ILC ⊕ PC

Measurements at LHC, ILC and Photon Collider are complementary, being sensitive to different combinations of Higgs-boson couplings

Cross sections × BR relative to SM

$M_H = 250 \text{ GeV}$



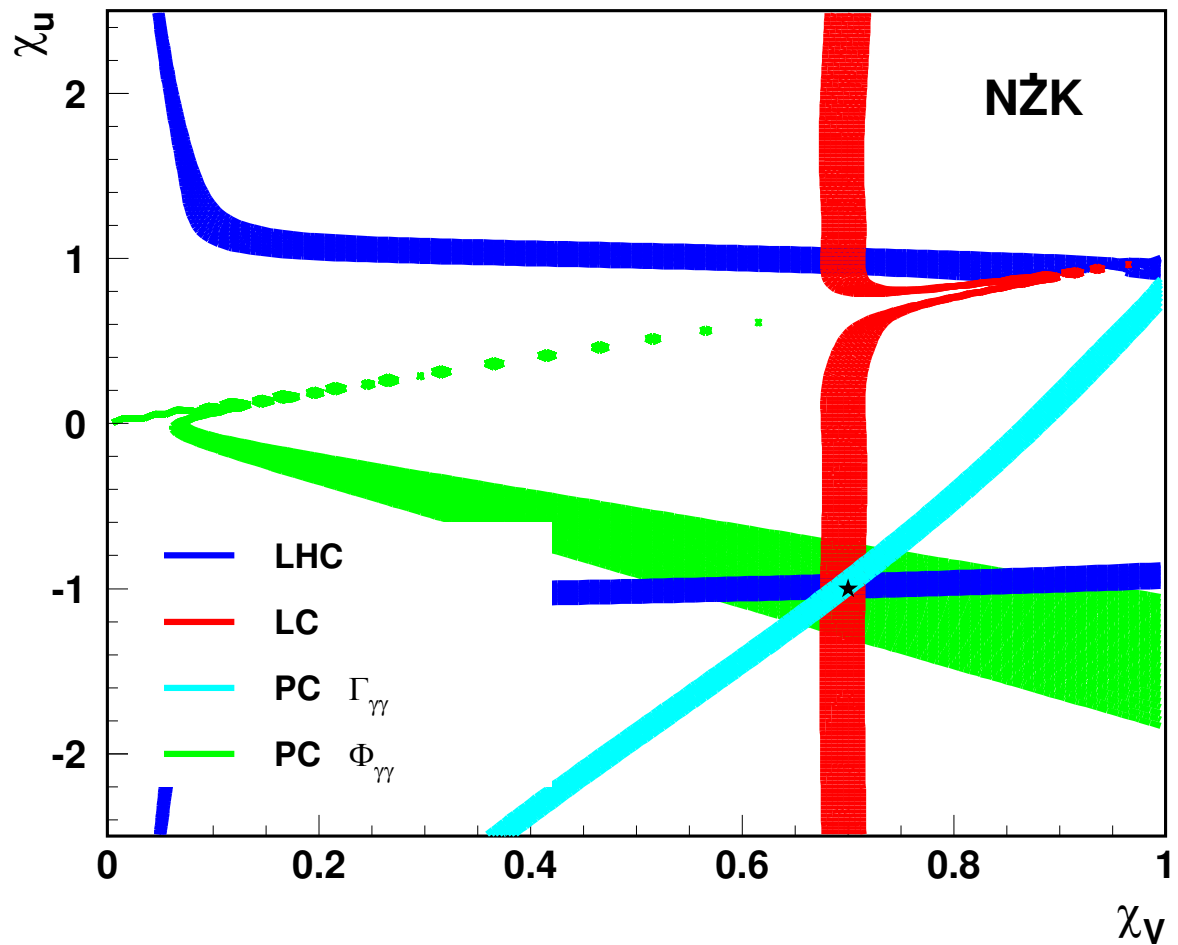
LHC ⊕ ILC ⊕ PC

Allowed coupling values (1σ) from **cross section** measurements at **LHC**, **ILC** and **PC**, and the phase measurement at **PC**.

Consistency of all these measurements verifies the **coupling structure of the model**

statistical errors only

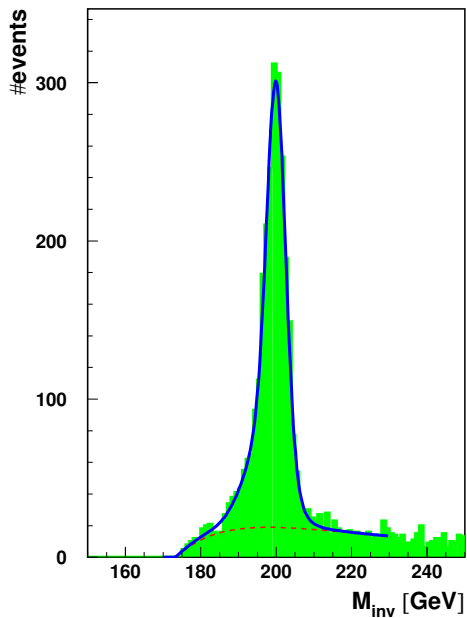
$$\chi_V = 0.7 \quad \chi_u = -1 \quad M_H = 250 \text{ GeV}$$



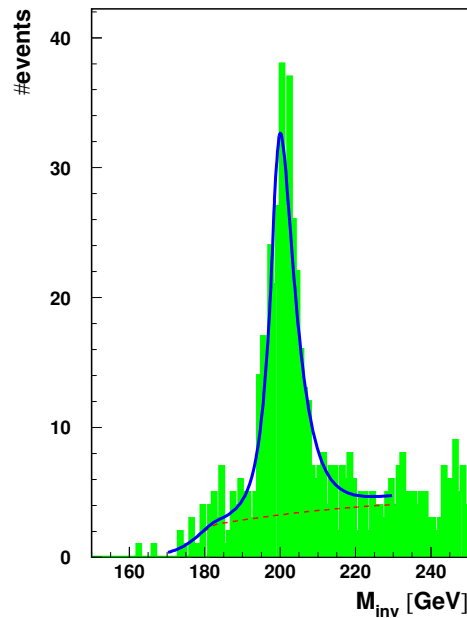
LHC ⊕ ILC ⊕ PC

Combined fit to the expected invariant mass distributions:

LHC

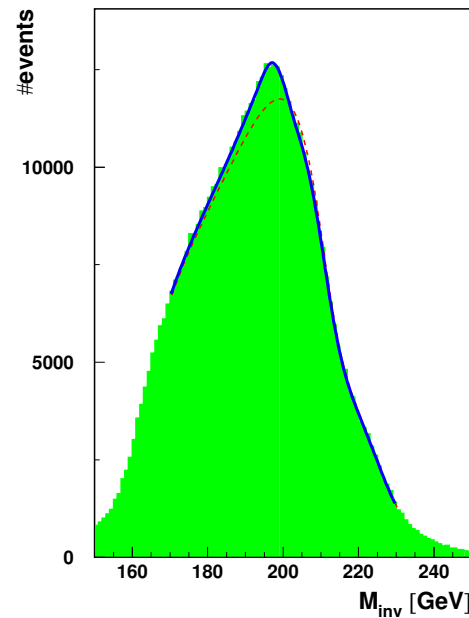


ILC



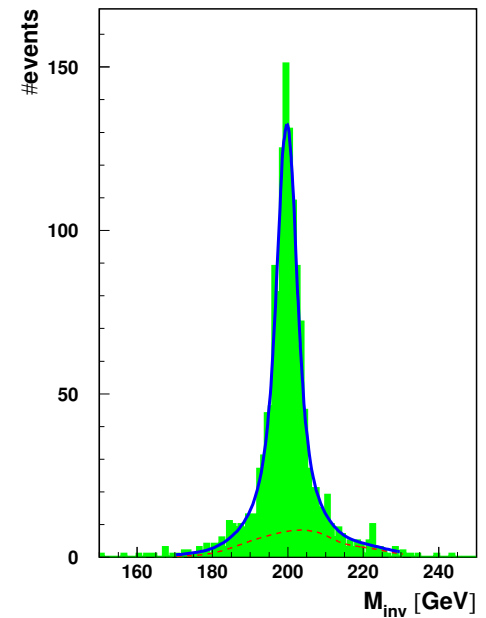
PC

$H \rightarrow WW$



PC

$H \rightarrow ZZ$



12 parameter fit: ● χ_V ● χ_u ● M_H ● Φ_{HA}

+ 8 normalization and $\gamma\gamma$ -spectra shape parameters (systematic uncertainties)

LHC ⊕ ILC ⊕ PC

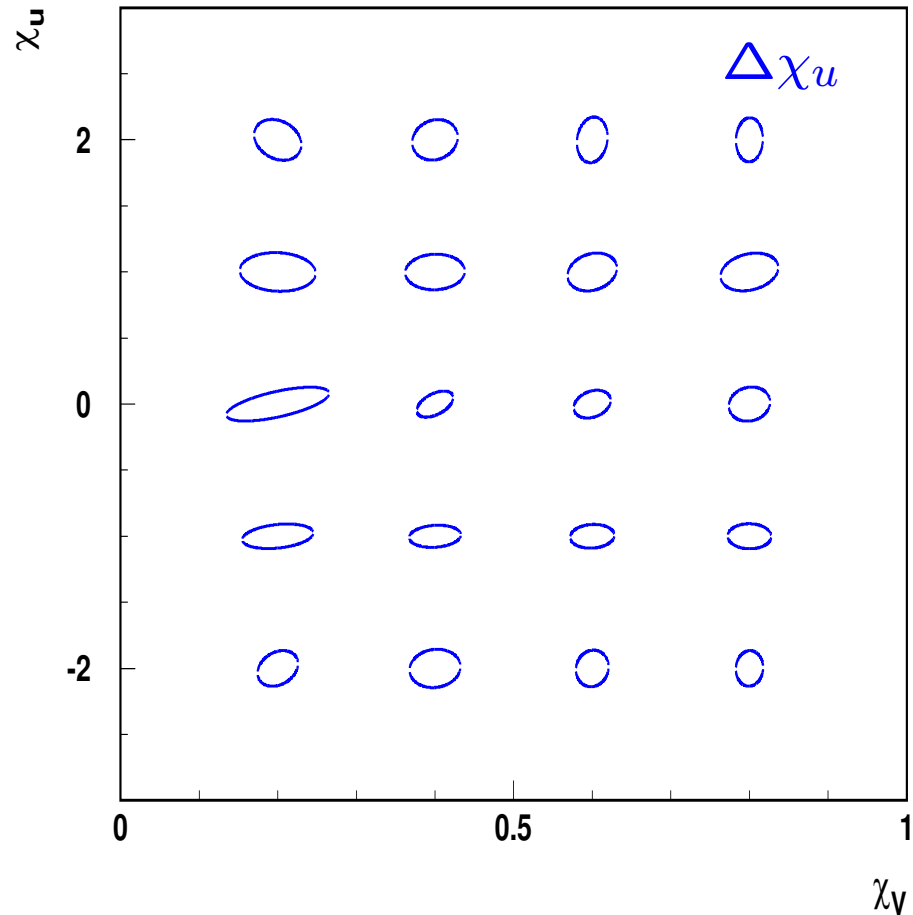
Simultaneous fit to LHC, ILC and PC (W^+W^- and ZZ) invariant mass distributions

1σ (stat.+sys.) contours

H couplings to vector bosons (χ_V) and up fermions (χ_u) for $M_H = 250$ GeV

$$\langle \Delta\chi_V \rangle = 0.03$$

$$\langle \Delta\chi_u \rangle = 0.13$$



LHC ⊕ ILC ⊕ PC

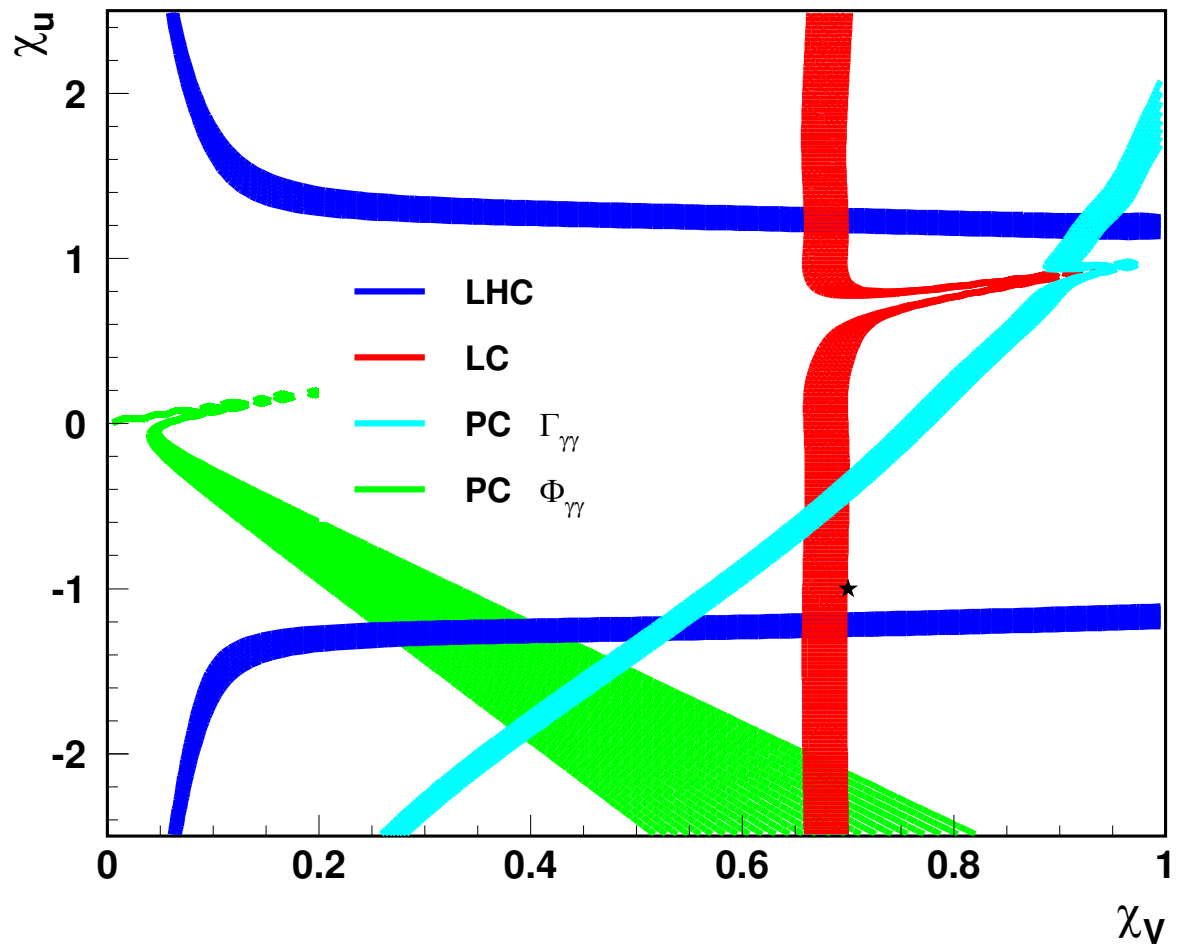
Allowed coupling values from **cross section** measurements at **LHC**, **ILC** and **PC**, and the phase measurement at **PC**.

Measurements compared assuming **CP-conserving 2HDM(II)** not consistent \Rightarrow "new physics":

- different **coupling structure** (eg. CP violation)
- existence of **new heavy particles** contributing to Γ_{gg} and $\Gamma_{\gamma\gamma}$

$$\chi_V = 0.7 \quad \chi_u = -1 \quad \Phi_{HA} = -0.2$$

NŻK



Summary

2HDM(II)

SM-like 2HDM(II)

Solution A

For light Higgs boson h :

$$\chi_u = \chi_d = \chi_V = 1$$

χ_i - couplings normalized to SM couplings

All couplings are the same as in SM.

$\Gamma_{\gamma\gamma}$ and $\phi_{\gamma\gamma}$ affected only by the H^+ loop

For heavy Higgs bosons H and A :

$$\chi_V \equiv 0$$

No decays to W^+W^- and ZZ ...

I. F. Ginzburg, M. Krawczyk and P. Osland,
Nucl. Instrum. Meth. A472:149, 2001
hep-ph/0101331; hep-ph/0101208.

Solution B_h

	h	H	A
χ_u	-1	$-\frac{1}{\tan\beta}$	$-i \gamma_5 \frac{1}{\tan\beta}$
χ_d	+1	$-\tan\beta$	$-i \gamma_5 \tan\beta$
χ_V	$\cos(2\beta)$	$-\sin(2\beta)$	0

$\tan\beta \rightarrow 0 \Rightarrow \text{sol. } B_u$

$\tan\beta \rightarrow \infty \Rightarrow \text{sol. } B_d$

Higgs production ($\Gamma_{\gamma\gamma}$ and $\phi_{\gamma\gamma}$)
and decays depend on $\tan\beta$.

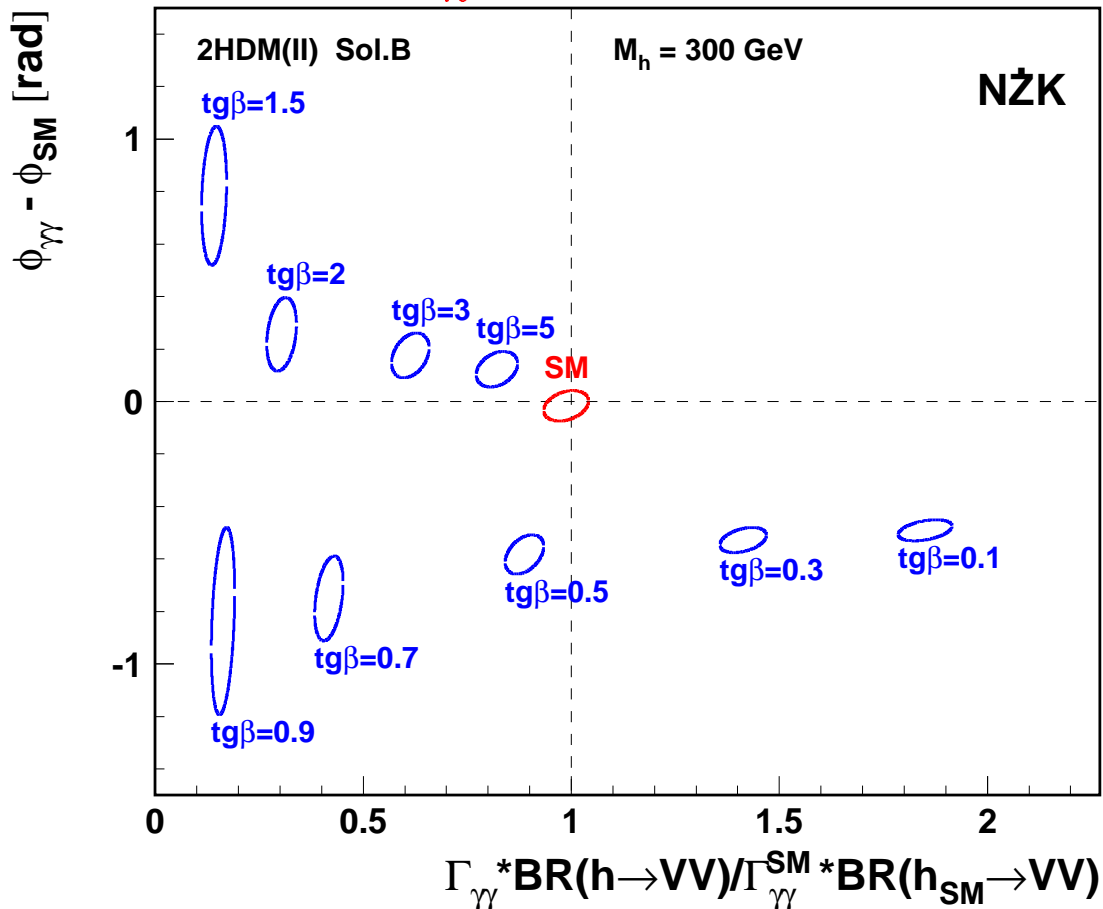
**Can we extract $\tan\beta$ value
from the measured W^+W^- and ZZ
invariant mass distributions ?**

SM-like 2HDM(II)

Light Higgs boson h

Two-photon width and phase measurement for different $\tan \beta$ $\chi_V = \cos 2\beta$

$M_h = 300 \text{ GeV}$



Measurement very sensitive to $\tan \beta$
 \Rightarrow precise determination possible.

Ambiguity resolved by the phase measurement (distinguishes between low $\tan \beta$ and large $\tan \beta$)

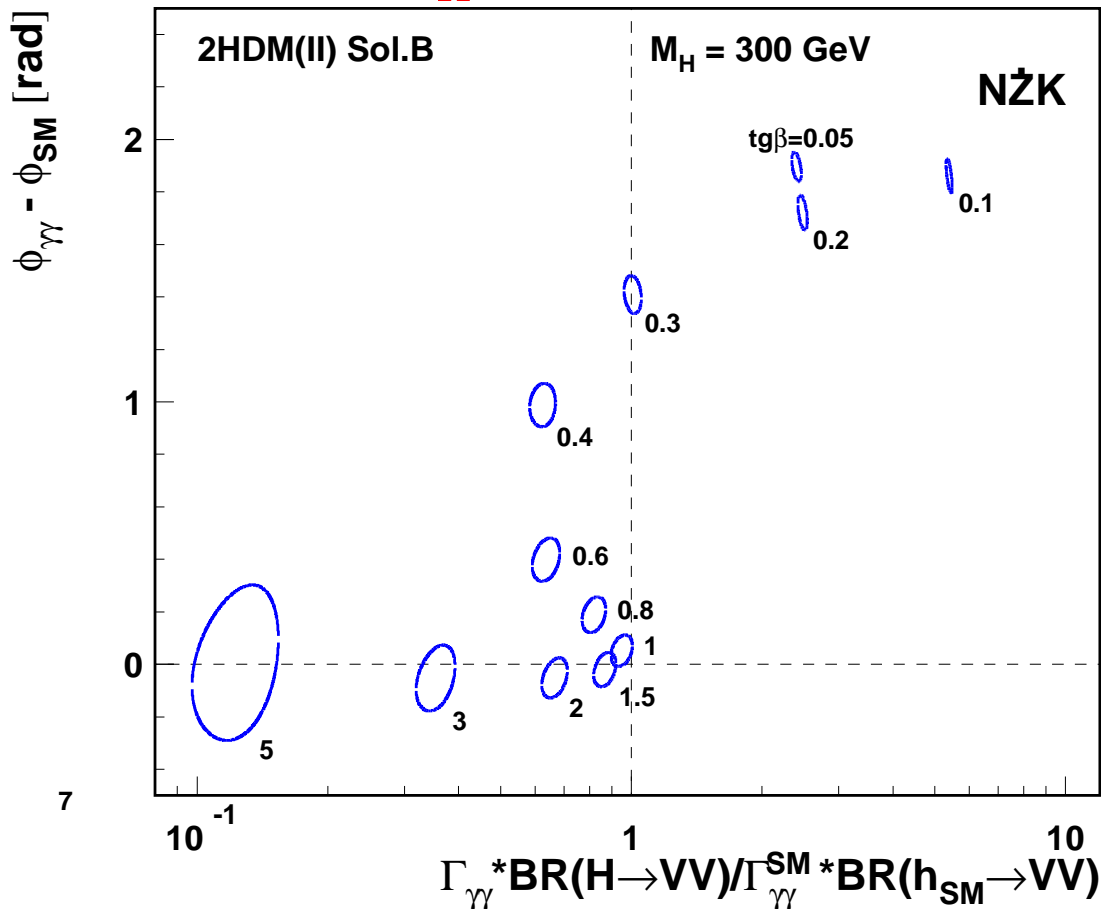
1σ contours for 1 year of PC running
 statistical errors only $M_{H^+} = 800 \text{ GeV}$

SM-like 2HDM(II)

Heavy Higgs boson H

Two-photon width and phase measurement for different $\tan \beta$ $\chi_V = -\sin 2\beta$

$M_H = 300 \text{ GeV}$



$\Gamma_{\gamma\gamma}$ enhancement for $\tan \beta < 1$ due to top contribution ($\chi_u = -\frac{1}{\tan \beta}$)

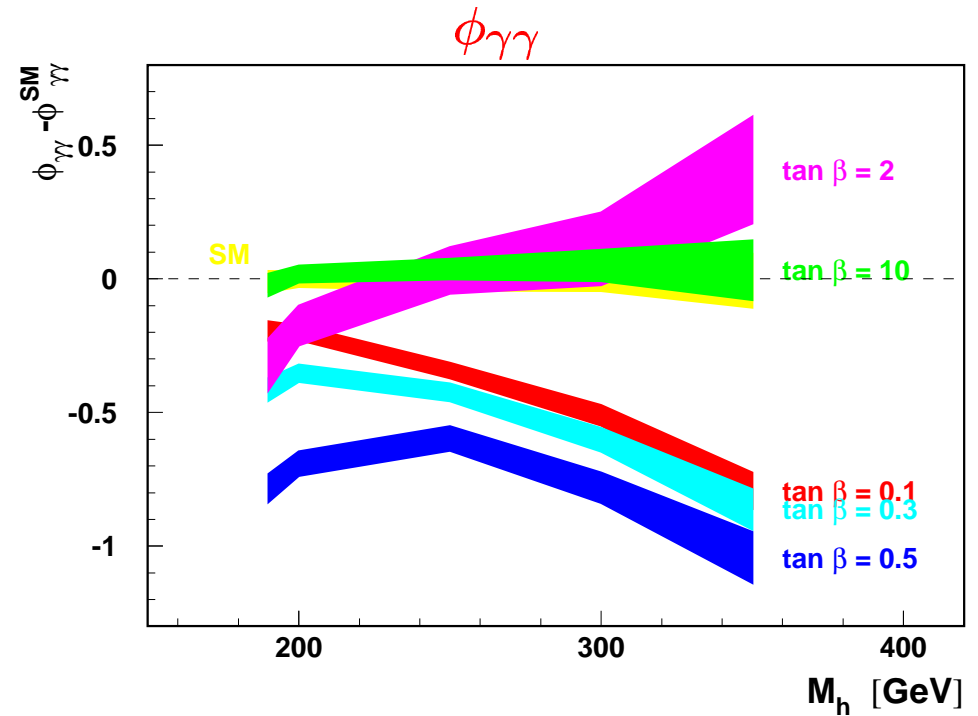
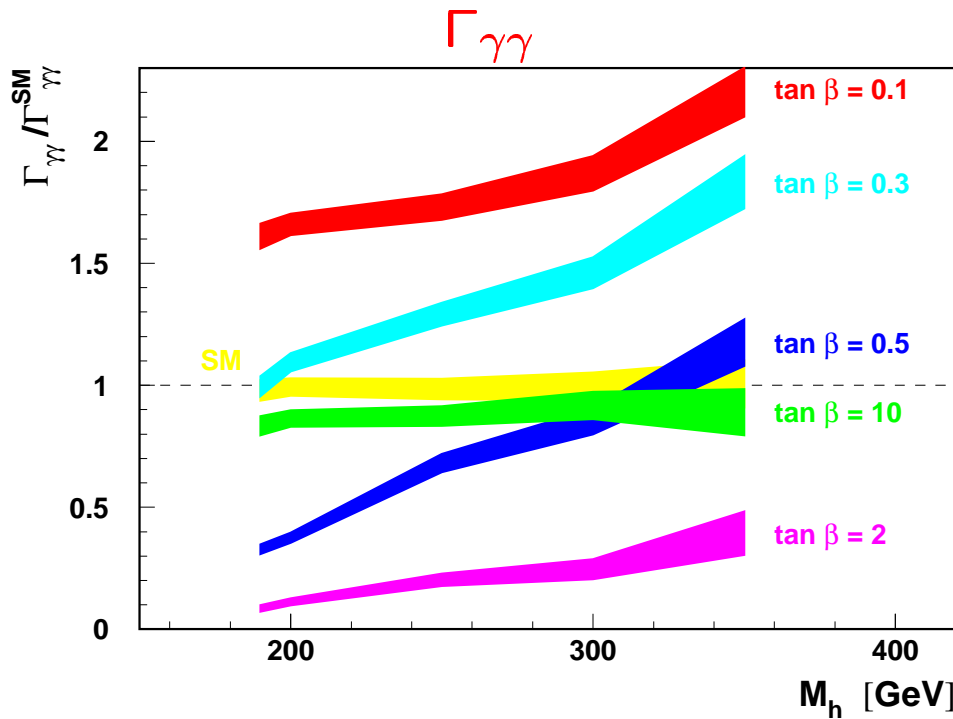
1σ contours for 1 year of PC running
statistical errors only

$M_h = 120 \text{ GeV}$, $M_{H^\pm} = 800 \text{ GeV}$

2HDM(II)

Solution B SM-like

Two-photon width and phase measurement for different M_h
 band width indicates statistical measurement error

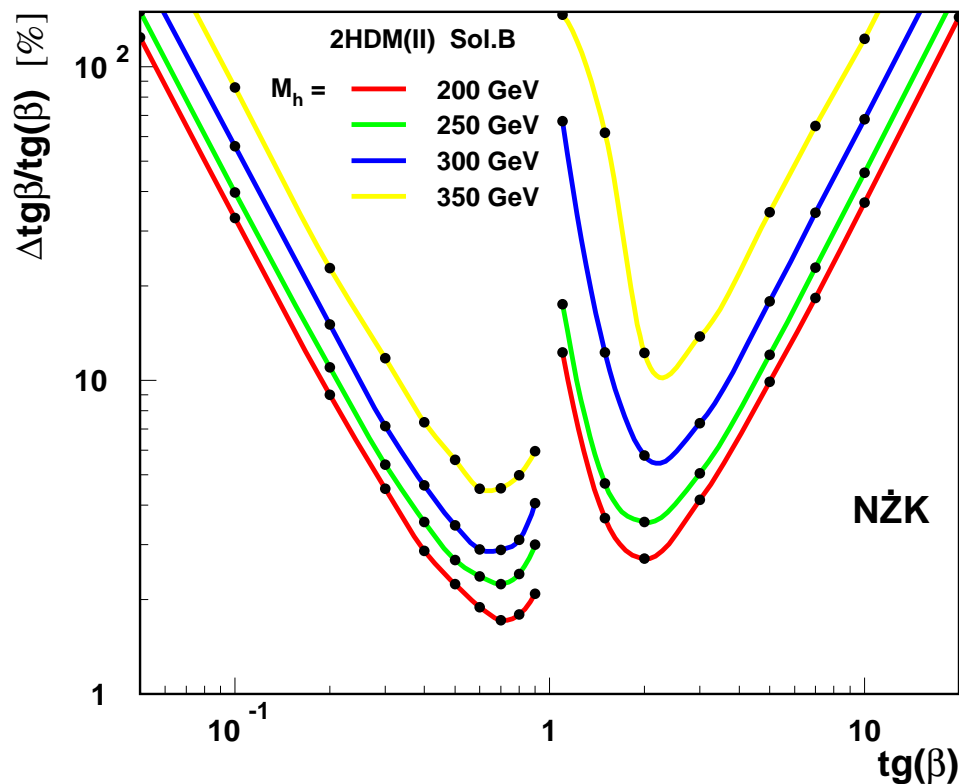


Measurement of both phase and width can help to test Higgs sector structure
 e.g. distinguish between large and small $\tan\beta$

SM-like 2HDM(II)

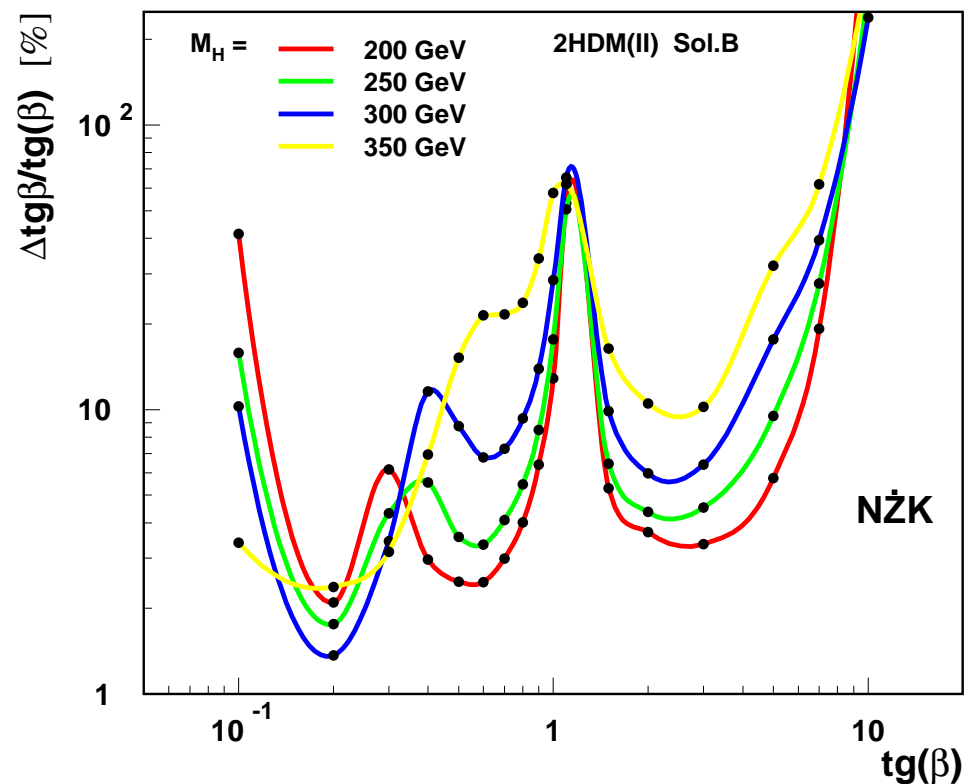
Light Higgs boson h

Expected precision in $\tan\beta$ determination
stat. + sys. errors



Heavy Higgs boson H

Expected precision in $\tan\beta$ determination
stat. + sys. errors



$\tan\beta$ can be determined with precision better than 10% in wide parameter range

2HDM(II)

SM-like 2HDM(II)

We consider SM-like **solution** B_h

Basic couplings, relative to SM:

$$\chi_x = g_{\mathcal{H}xx} / g_{\mathcal{H}xx}^{SM} \quad \mathcal{H} = h, H, A$$

	h	H	A
χ_u	-1	$-\frac{1}{\tan\beta}$	$-i \gamma_5 \frac{1}{\tan\beta}$
χ_d	+1	$-\tan\beta$	$-i \gamma_5 \tan\beta$
χ_V	$\cos(2\beta)$	$-\sin(2\beta)$	0

CP conserving model:

Higgs production ($\Gamma_{\gamma\gamma}$ and $\phi_{\gamma\gamma}$) and decays depend on $\tan\beta$ only.

For charged Higgs boson couplings (loop contribution to $\Gamma_{\gamma\gamma}$) we set

$$M_{H^\pm} = 800 \text{ GeV} \quad \mu = 0$$

CP violation

Mass eigenstates of the neutral Higgs-bosons h_1, h_2 and h_3 do not need to match CP eigenstates h, H and A .

We consider **weak CP violation** through a small mixing between H and A states:

$$\begin{aligned} \chi_X^{h_1} &\approx \chi_X^h \\ \chi_X^{h_2} &\approx \chi_X^H \cdot \cos\Phi_{HA} + \chi_X^A \cdot \sin\Phi_{HA} \\ \chi_X^{h_3} &\approx \chi_X^A \cdot \cos\Phi_{HA} - \chi_X^H \cdot \sin\Phi_{HA} \end{aligned}$$

⇒ additional model parameter:

CP-violating mixing phase Φ_{HA}

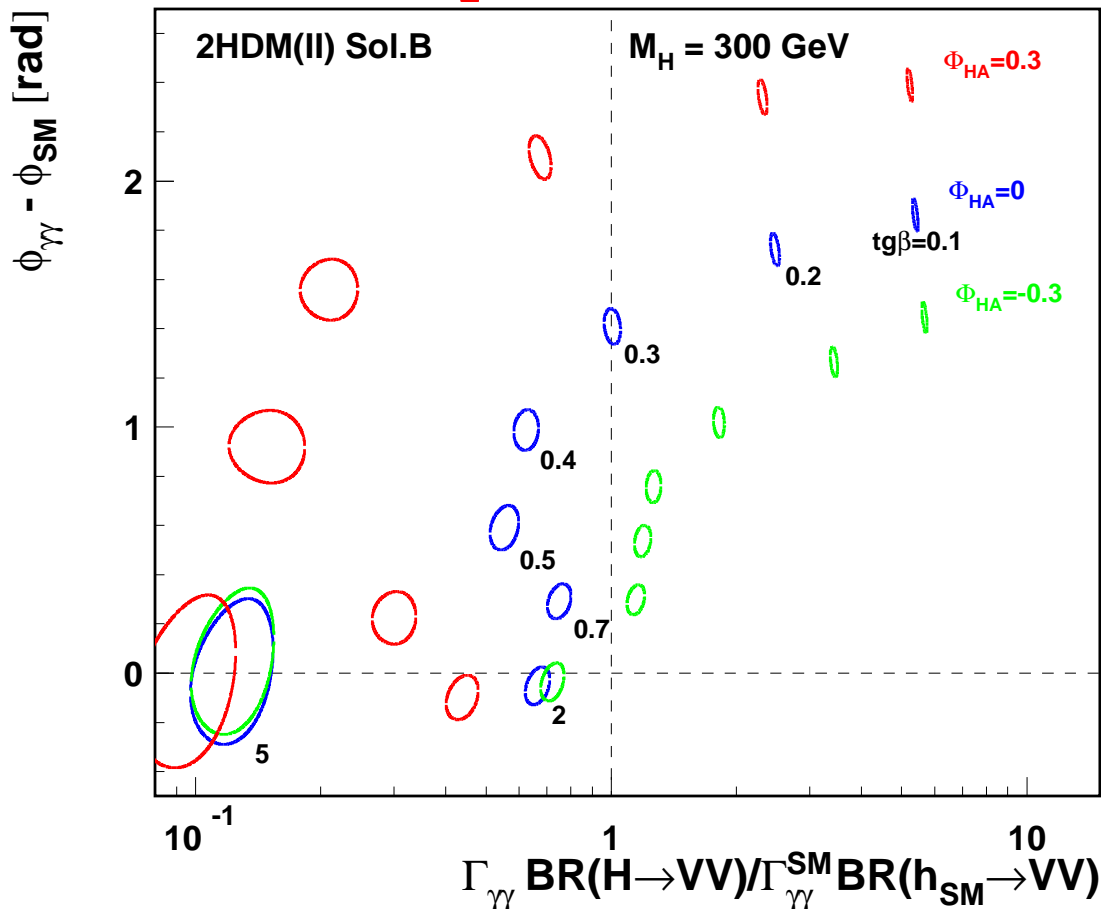
We consider h_2 production and decays

2HDM(II)

Higgs boson h_2

Two-photon width and phase measurement for different $\tan \beta$ and Φ_{HA}

$$M_{h_2} = 300 \text{ GeV}$$



1σ contours for 1 year of PC running
statistical errors only

$$M_h = 120 \text{ GeV}, M_{H^\pm} = 800 \text{ GeV}$$

Expected precision at PLC:
(for small mixing i.e. $\Phi_{HA} \sim 0$)

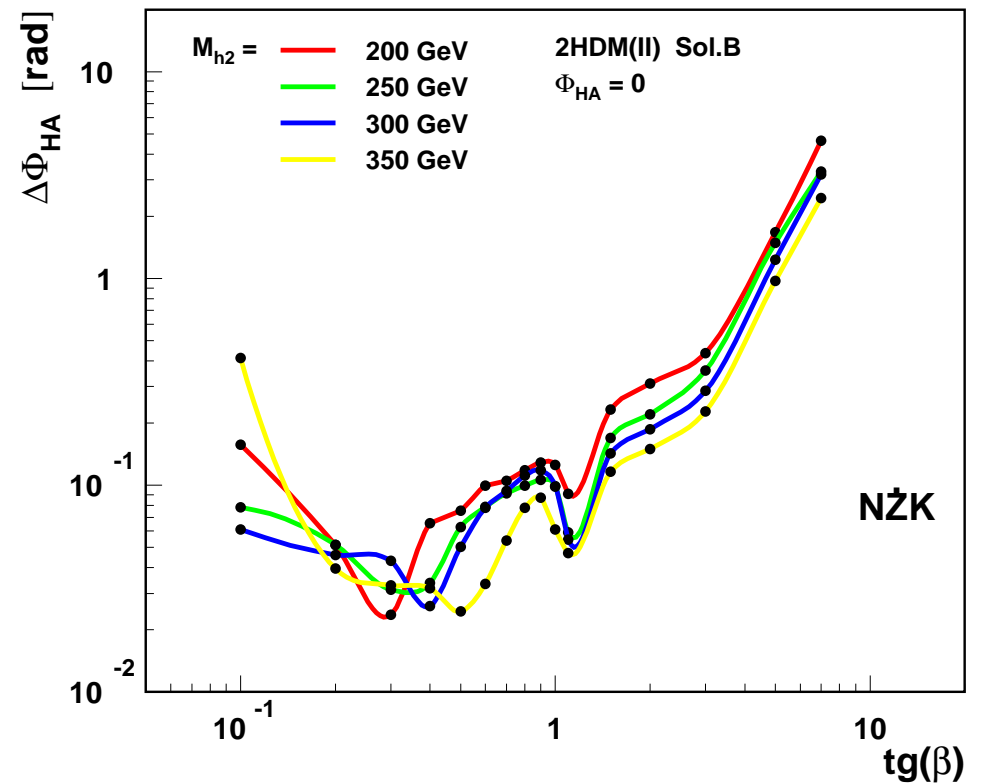
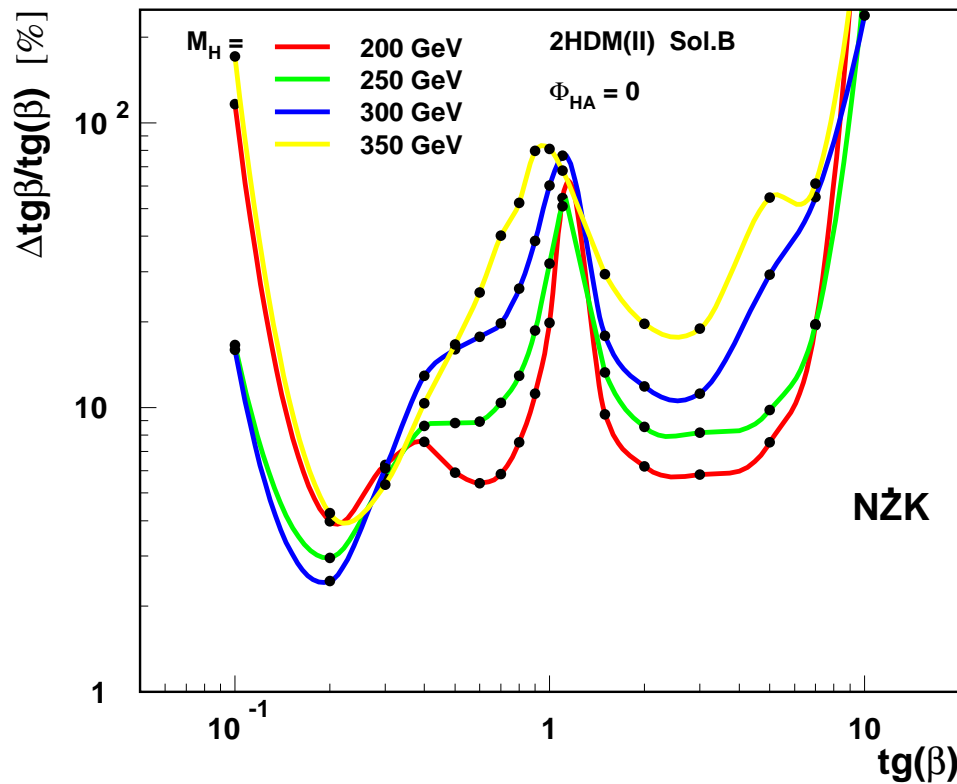
- $\sim 10\%$ for $\tan \beta$
- $\sim 100 \text{ mrad}$ for Φ_{HA}
(for low $\tan \beta$)

2HDM(II)

Higgs boson h_2

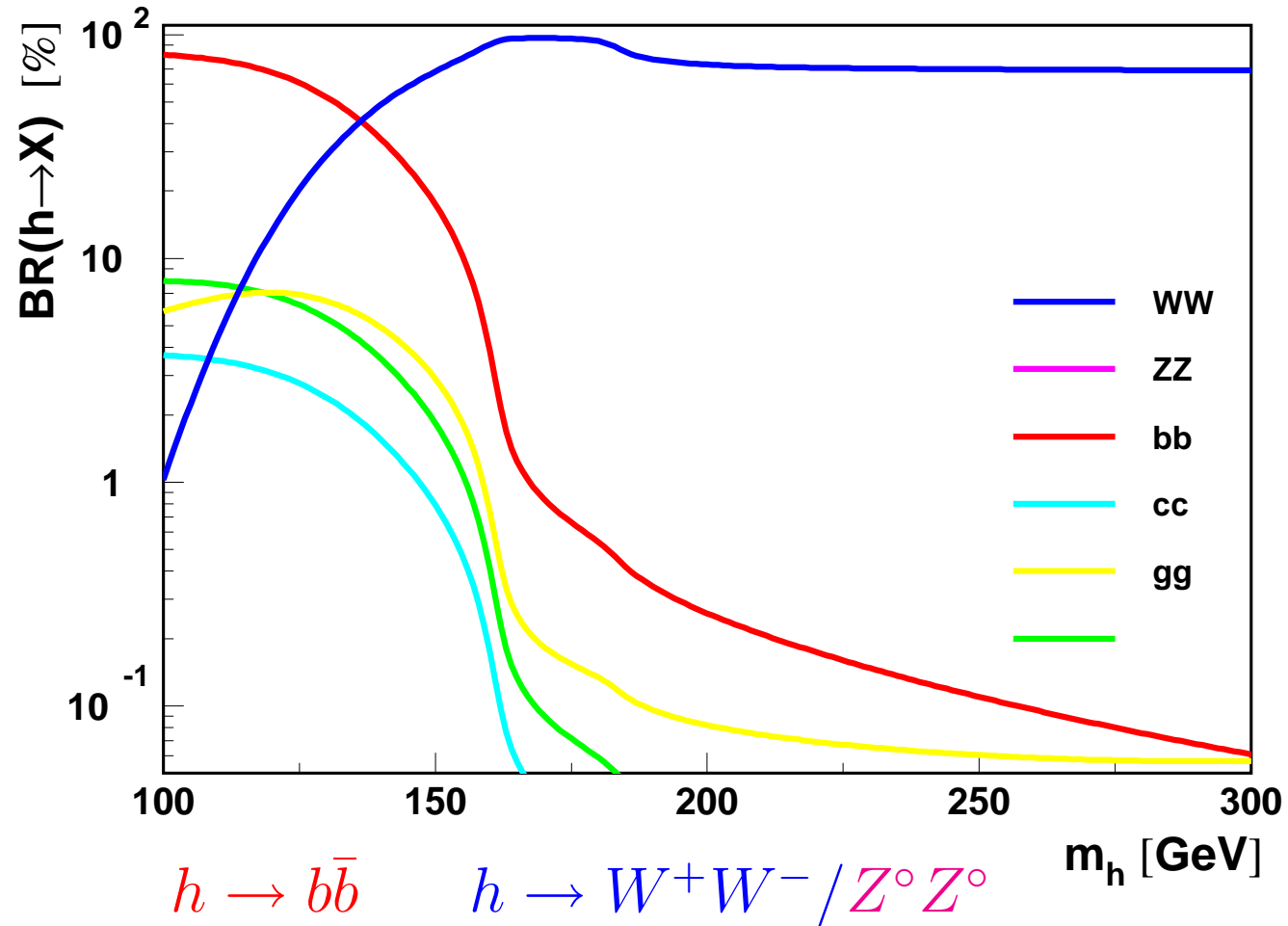
Solution B_h (with CP violation) \Rightarrow two free parameters ($\tan \beta$ and Φ_{HA})

Expected precision in $\tan \beta$ and Φ_{HA} determination at PLC (stat.+sys. errors)



CP violating H–A mixing angle can be precisely measured, if $\tan \beta$ is not too large

Higgs decays



Analysis

Event selection

W^+W^-

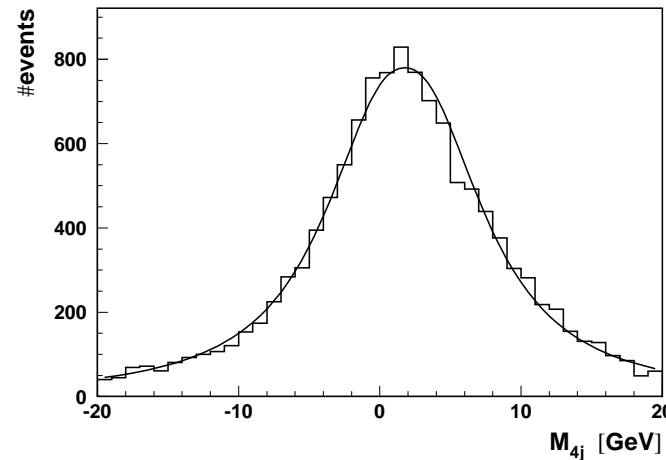
- balanced transverse momentum:
 $P_T/E_T < 0.1$
- 4 hadronic jets reconstructed
(Durham algorithm)
- cut on jet angle $\cos \theta_{jet} < 0.95$
to preserve good mass resolution
- two W^\pm can be reconstructed
with probability $P_W > 0.001$

ZZ

- balanced transverse momentum:
 $P_T/E_T < 0.1$
- 2 leptons (e^\pm or μ^\pm) + 2 hadronic jets
too large background in 4-jet channel
- cut on lepton and jet angle $\cos \theta_{l,jet} < 0.95$
- leptons and jets reconstruct into two Z°
with probability $P_Z > 0.001$

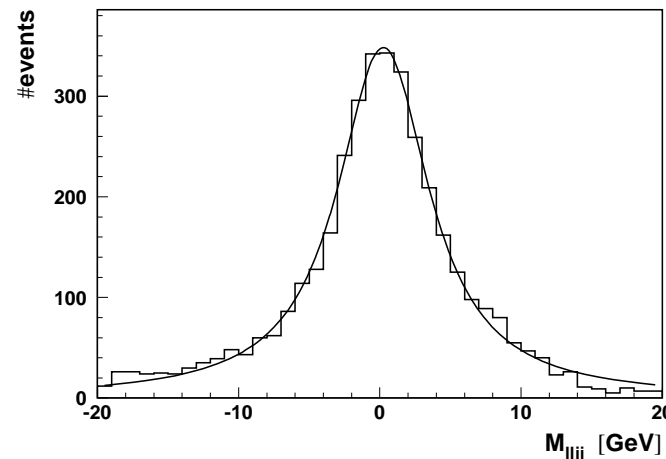
Mass resolution

$\gamma\gamma \rightarrow W^+W^- \rightarrow 4 \text{ jets}$



$\Gamma \sim 6.5 - 13 \text{ GeV}$

$\gamma\gamma \rightarrow ZZ \rightarrow l^+l^- + 2 \text{ jets}$



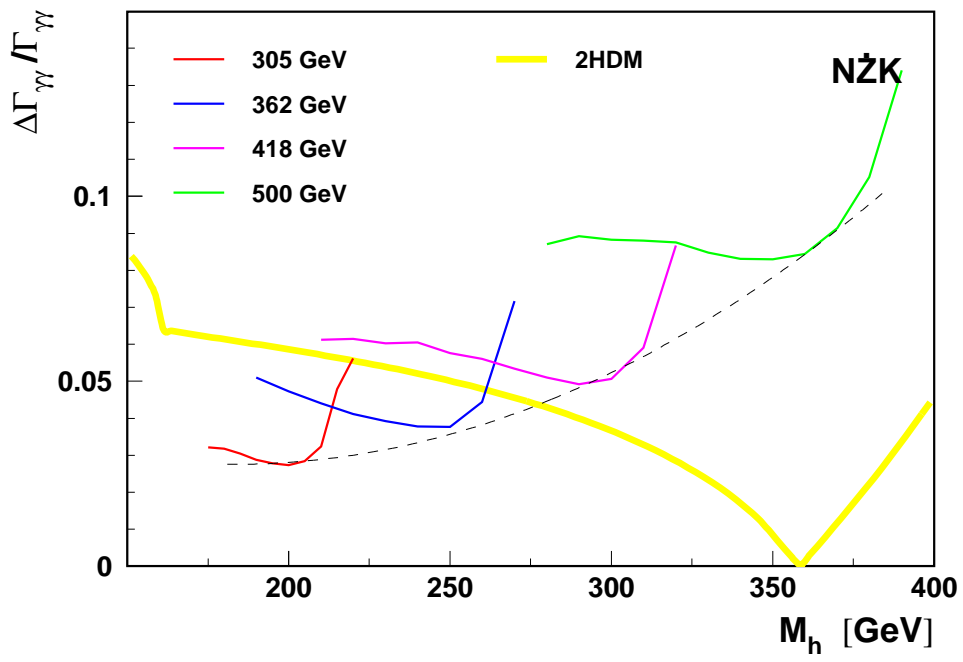
$\Gamma \sim 5.5 - 7.5 \text{ GeV}$
for
 $W_{\gamma\gamma} = 200 - 400 \text{ GeV}$

SM results

$\Gamma_{\gamma\gamma}$ measurement

One parameter fit to invariant mass distribution for W^+W^- and ZZ events

⇒ Average statistical precision (1 PC year):



assuming SM branching ratios

Sensitive to possible “new physics” only up to $M_h \sim 280$ GeV

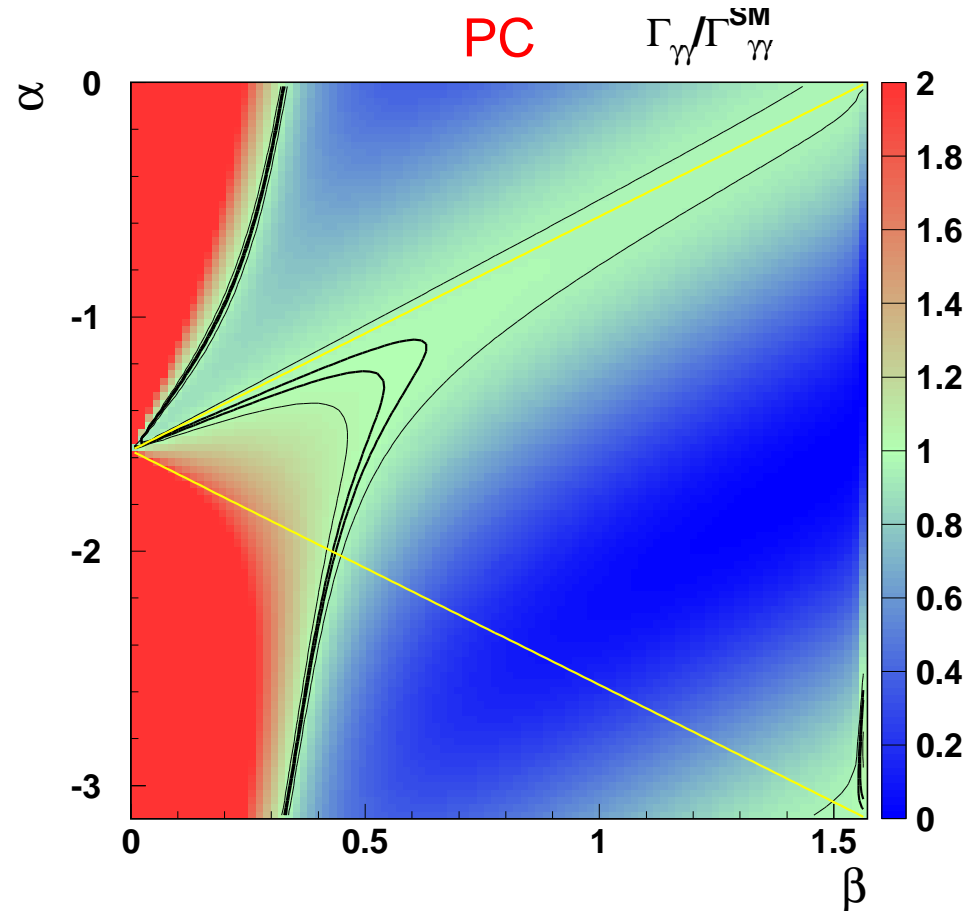
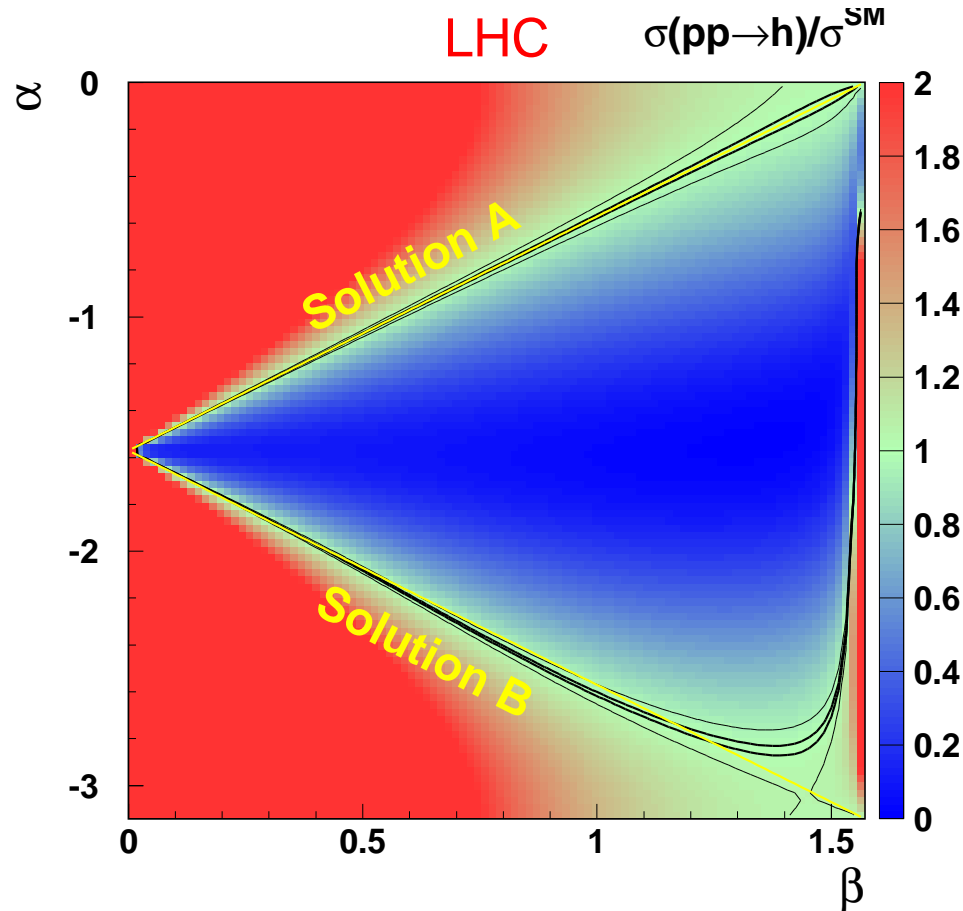
For higher Higgs masses $\Gamma_{\gamma\gamma}$ is little sensitive to contribution of new heavy charged particles !

“new physics” modeled by SM-like 2HDM (II) with $M_{H^+} = 800$ GeV

2HDM(II)

Higgs production cross section compared to SM

$M_h=300$ GeV



black contours: ± 1 and $\pm 5\%$ deviations

\Rightarrow significant effects in solution B