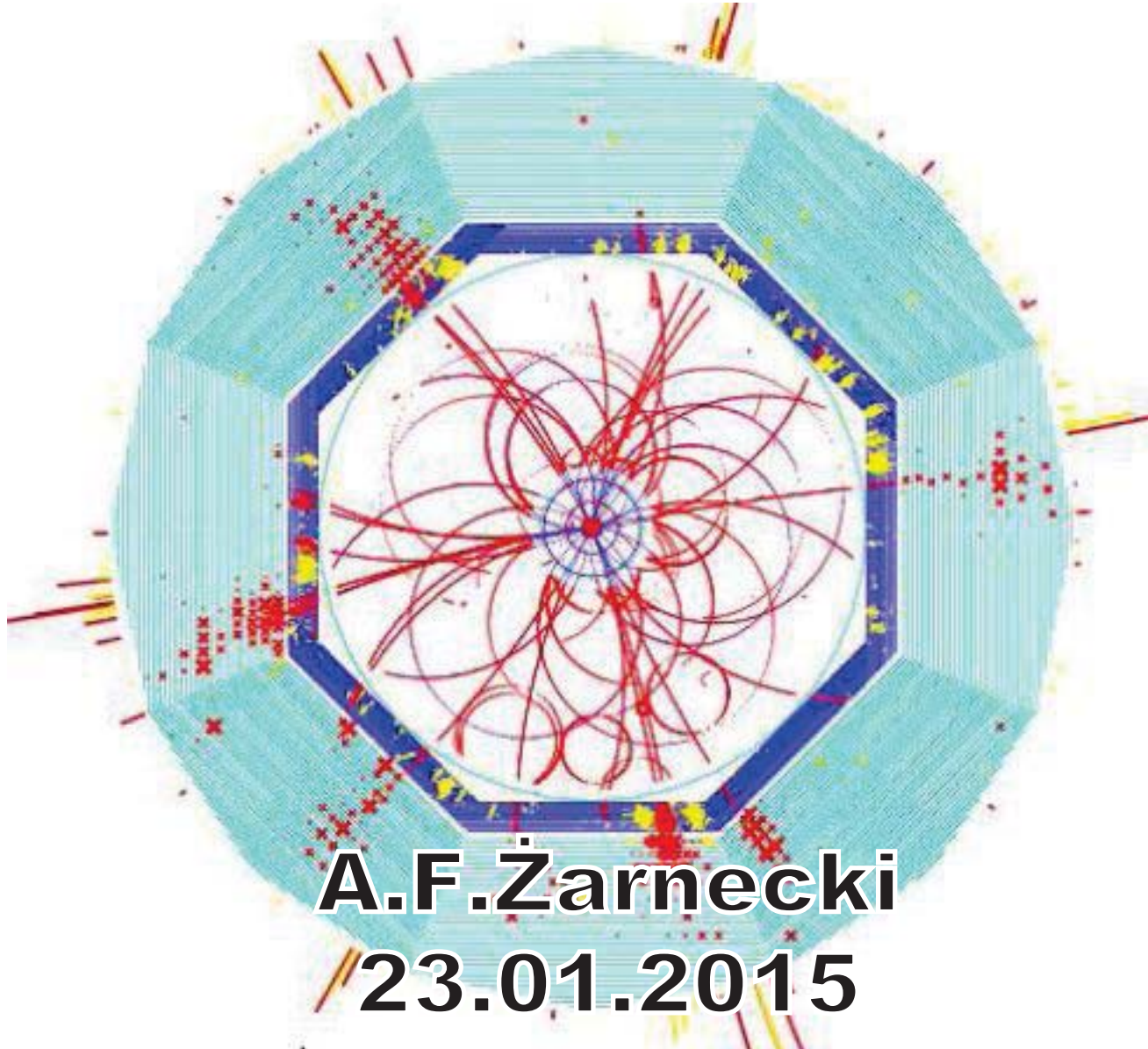


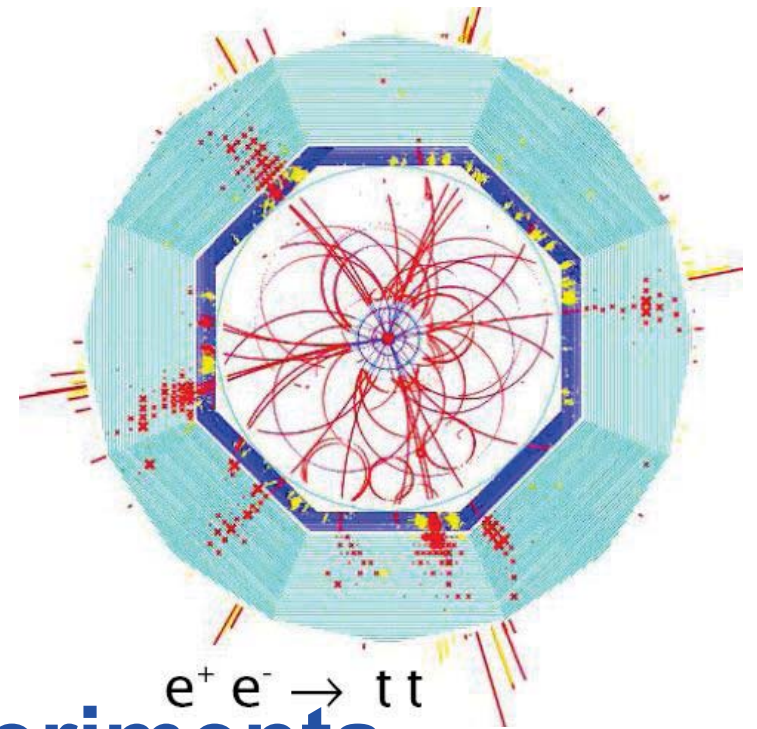
Top physics at future e^+e^- colliders



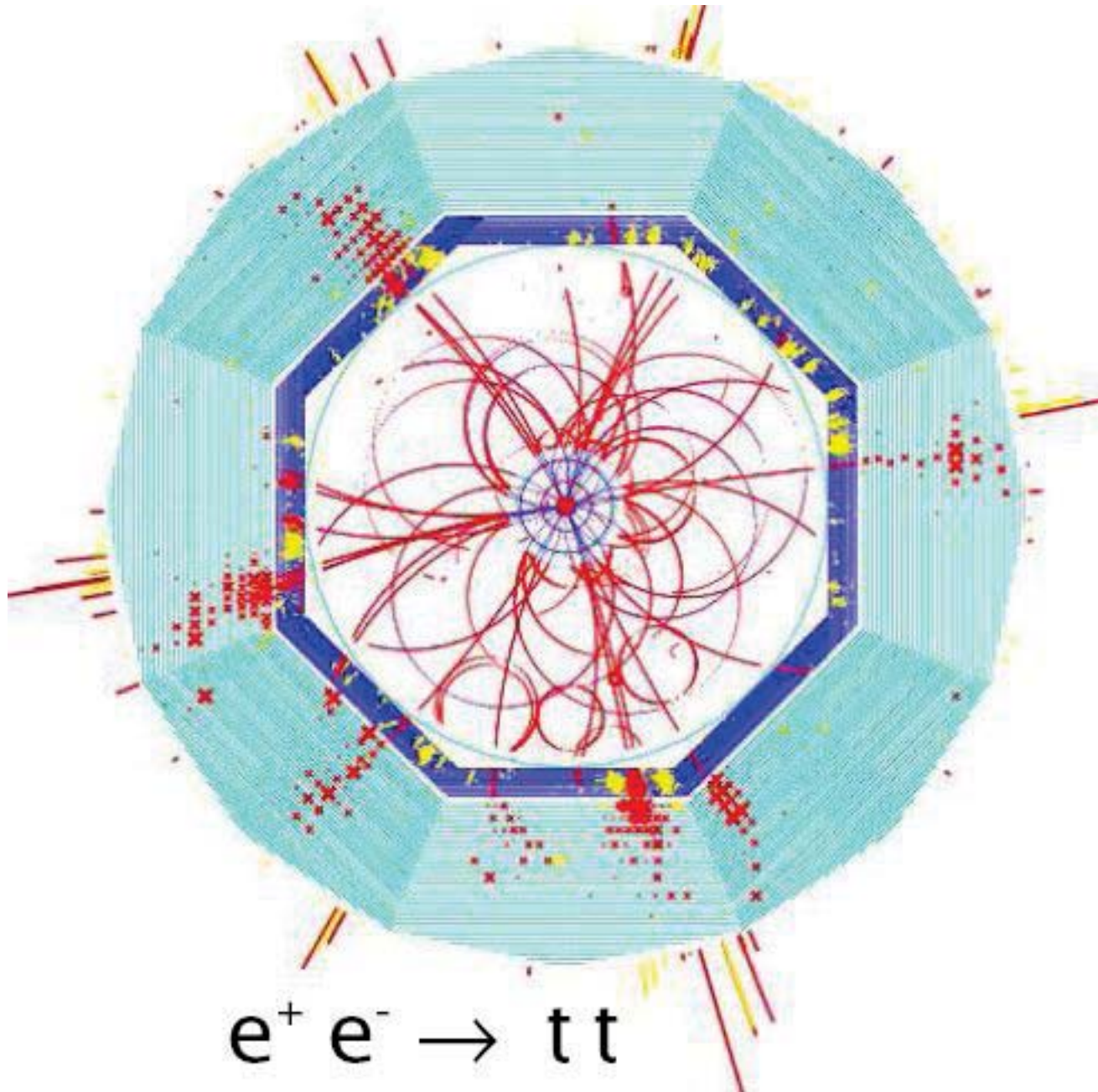
A.F. Żarnecki
23.01.2015

Outline

- ◆ Introduction
- ◆ Motivation
- ◆ Future colliders and experiments
- ◆ Prospects for precise measurements
 - threshold scan and top mass
 - Yukawa coupling
 - EW couplings



Introduction



$$e^+ e^- \rightarrow t t$$

Coronation of the Standard Model
and
First step on a road yet largely unexplored
Slightly modified citation of Barbieri arXiv:1309.3447



HIGGS

HIGGS

Particle Physics

Chip Brook, Snowmass Summary Talk

Where do we go from here?



Peeping through the Higgs window!

P5 likes Particle Astrophysics

Five science drivers:

- Higgs boson
- Neutrino mass
- Dark matter
- Cosmic acceleration
- Explore the unknown



Particle Astrophysics experiments address all but the first

Many experiments address more than one

Report supports expanded dark matter program, new cosmic surveys, and a new multi-agency program in CMB



Probing the Standard Model

- SM is self-consistent model accounting all particle physics phenomena at energy of current accelerators
 - with m_H all parameters of SM are known

$m_W = 80385 \pm 15 \text{ MeV}$

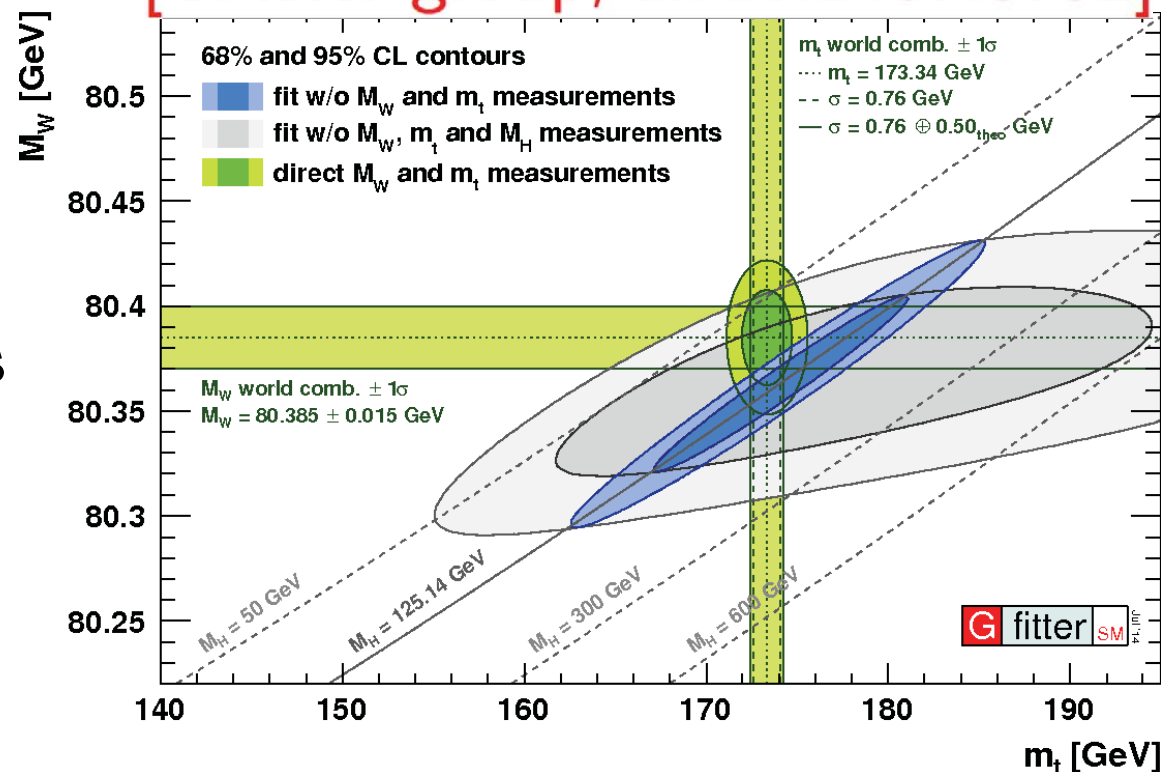
$m_t = 173.34 \pm 0.76 \text{ GeV}$

$m_H = 125.36 \pm 0.41 \text{ GeV}$

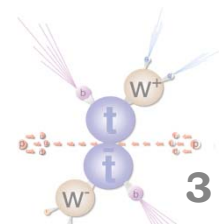
Current p-value for (data | SM)=0.2

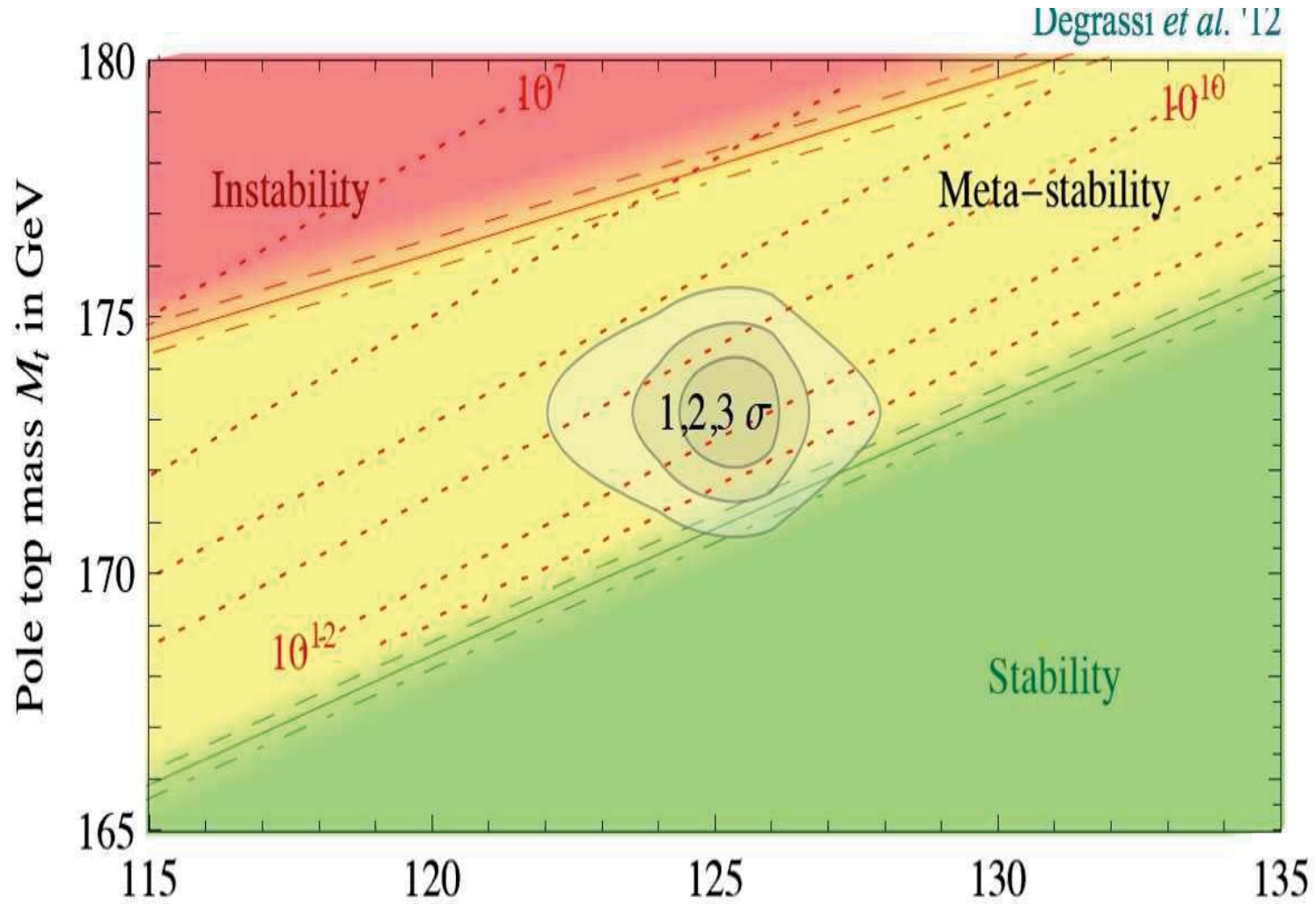
Need to improve m_W , m_t and m_H

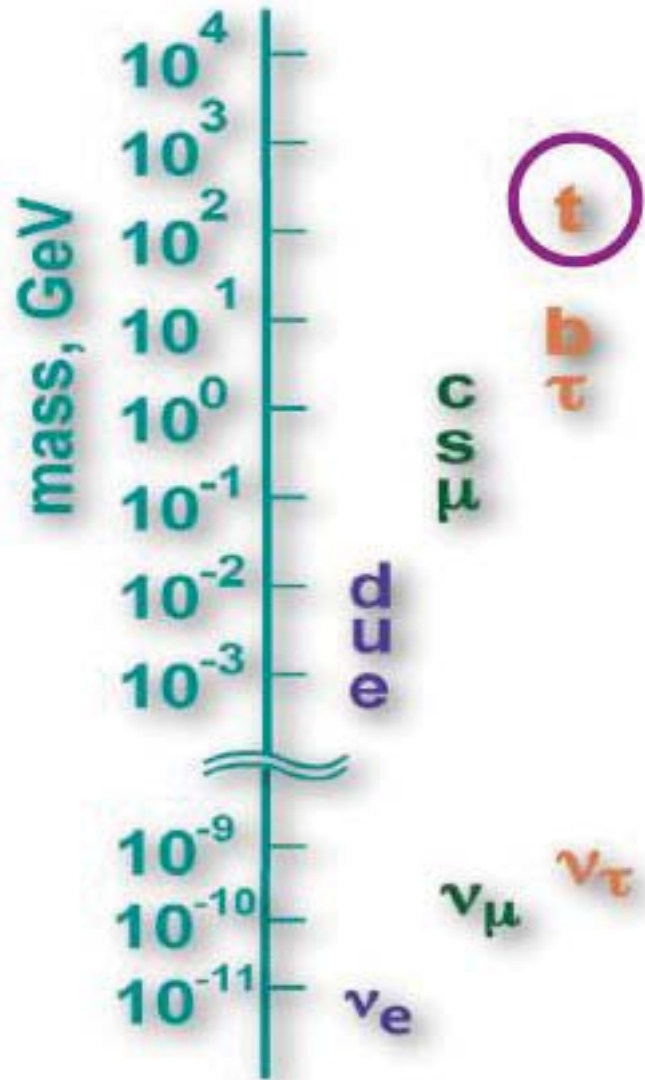
[GFitter group, arXiv:1407.3792]



Precision tests of further consistency of the SM are mandatory



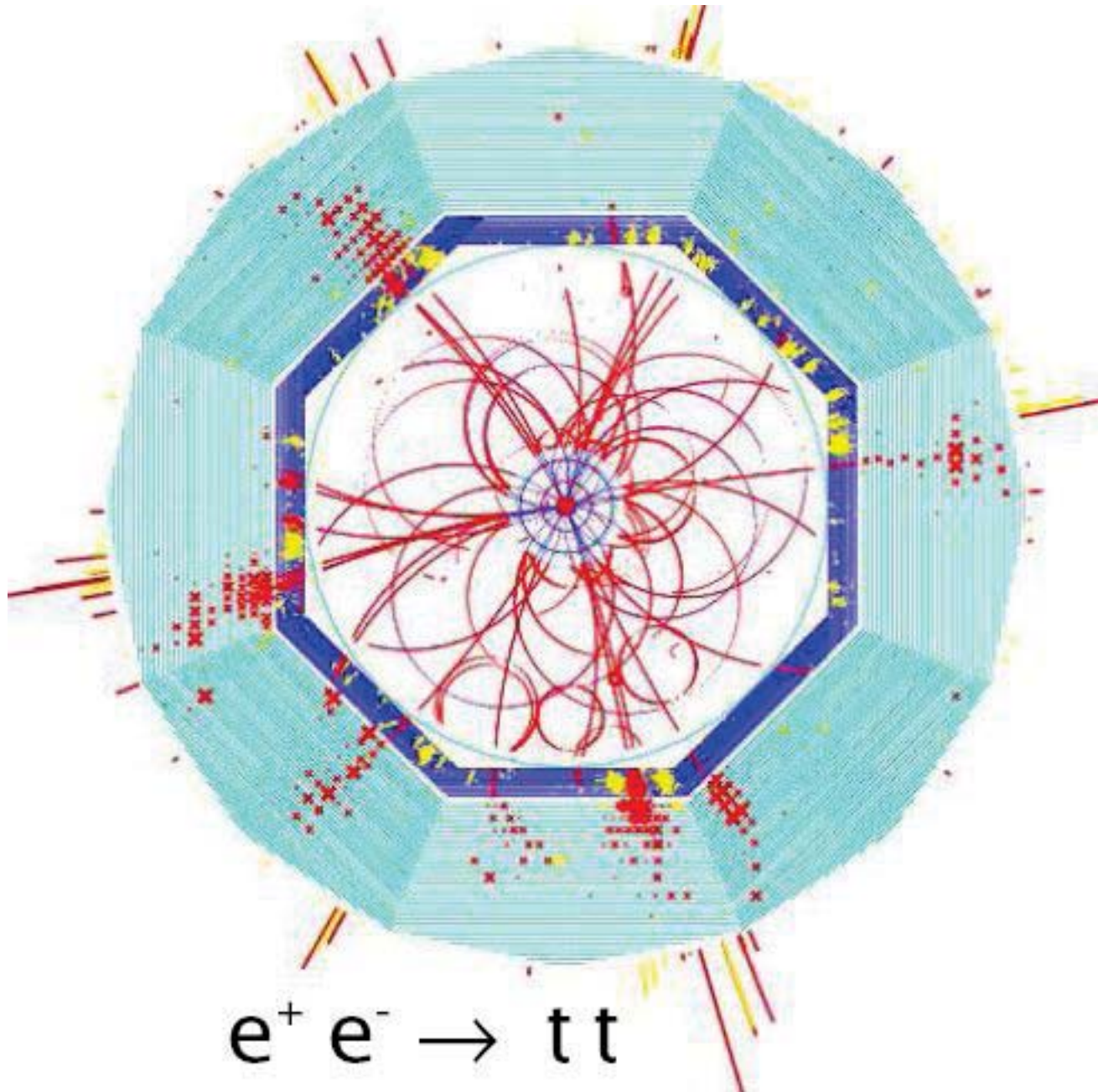




- SM does not provides no explanation for mass spectrum of fermions (and gauge bosons)
- Fermion mass generation closely related to the origin electroweak symmetry breaking
- Expect residual effects for particles with masses closest to symmetry breaking scale

Strong motivation to study chiral structure of top vertex in high energy e⁺e⁻ collisions

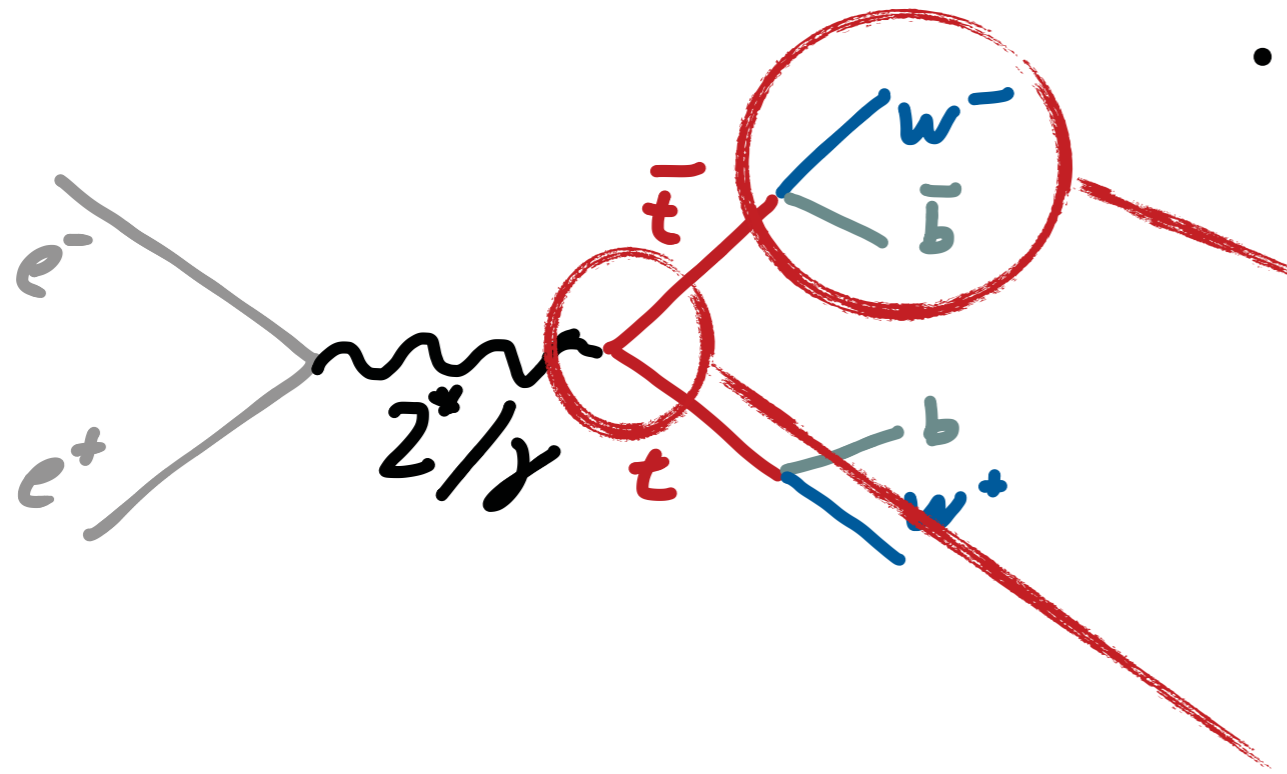
Motivation



$$e^+ e^- \rightarrow t t$$

Top Quark Physics at Linear Colliders

- The dominant production mechanism: Top pair production

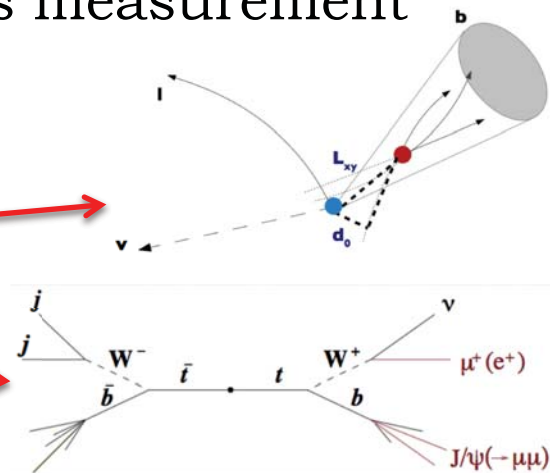


- Rich physics opportunities:
 - Top properties: **mass**, width, decay modes
 - BSM sensitivity: CP violation, flavor-changing decays,...
 - Top properties: **mass**, width,
 - Yukawa coupling, strong coupling constant
 - **Electroweak couplings** - sensitivity to BSM physics
- Measurements enabled by
 - known initial state & clean final state
 - Possibility for **polarized beams** - crucial for coupling measurements

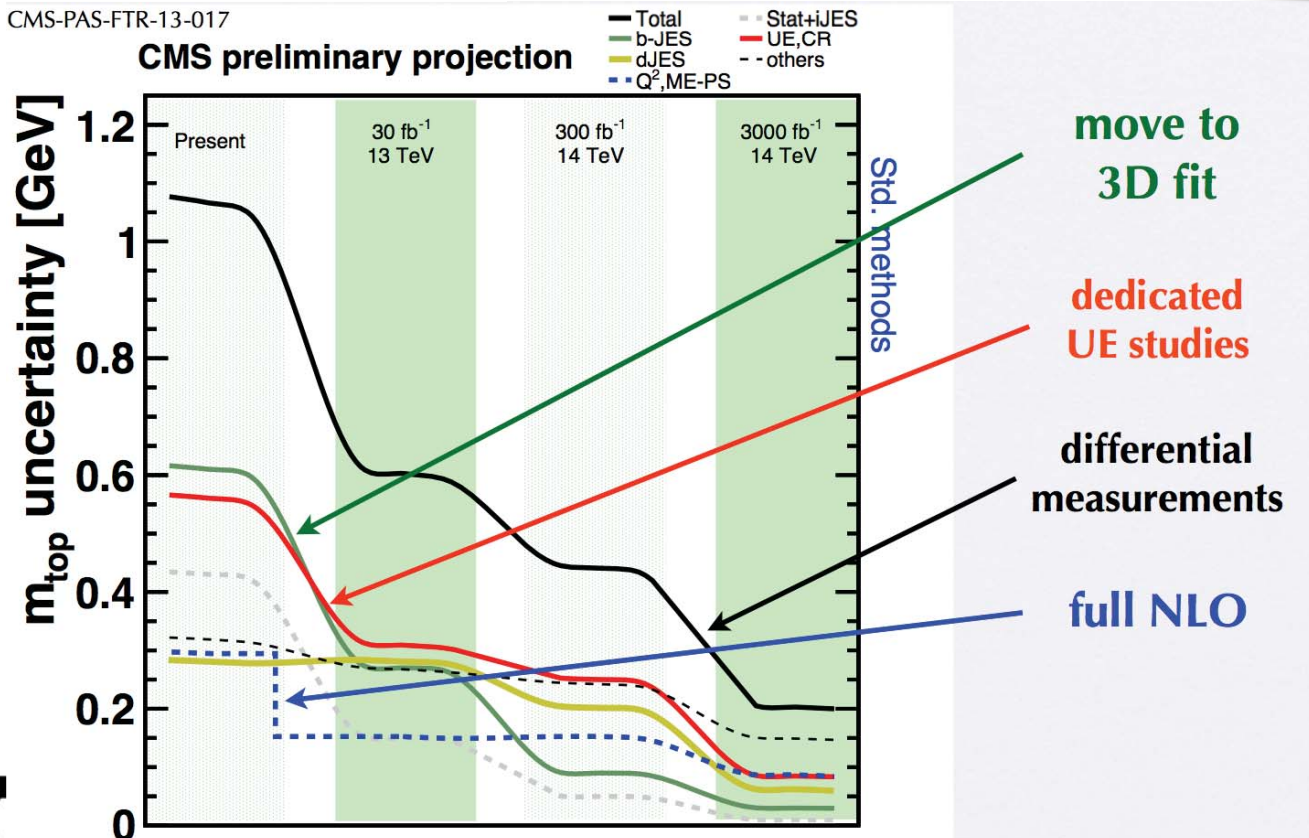


Top Mass Outlook

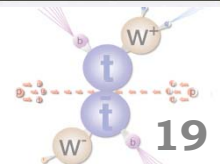
- For the projections used the baseline lepton+jets measurement at 7 TeV: JHEP **12** (2012) 105
 - [CMS PAS FTR-13-017](#)
 - And more methods with higher stats
 - Kinematic endpoints (**M_{lb} distribution**)
 - B-hadron lifetime (aka L_{xy})
 - J/Psi method



CMS-PAS-FTR-13-017



CM (TeV)	7	13	14
$L_{int} (fb^{-1})$	5	30	300
J/ψ		1.8	0.8
$L_{xy} (8 \text{ TeV})$	3.4	1.3	0.6
Endpoints	2.1	1.1	0.6
Standard	1.1	0.6	0.4



Concept of a Quark Mass

Short-distance mass schemes:

$$m^{\text{sd}}(R) = m^{\text{pole}} - R \left(a_1 \frac{\alpha_s}{4\pi} + a_2 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots \right)$$

Generic form of a short-distance mass scheme.

MS mass: $R = \bar{m}(\mu), \quad a_1 = \frac{16}{3} + 8 \ln \frac{\mu}{m}$

Processes where heavy quarks are off-shell and energetic.

Threshold masses (1S, PS, RS, kinetic masses)

$$R \sim m\alpha_s$$

Quarkonium bound states: heavy quarks are close to their mass-shell.

Jet masses (jet mass)

$$R \sim \Gamma_Q$$

Single quark resonance: heavy quarks are very close to their mass-shell.

The a_i 's are chosen such that the renormalon is removed.

The scale R is of order the momentum scale relevant for the problem.

Top Mass at e^+e^- Colliders

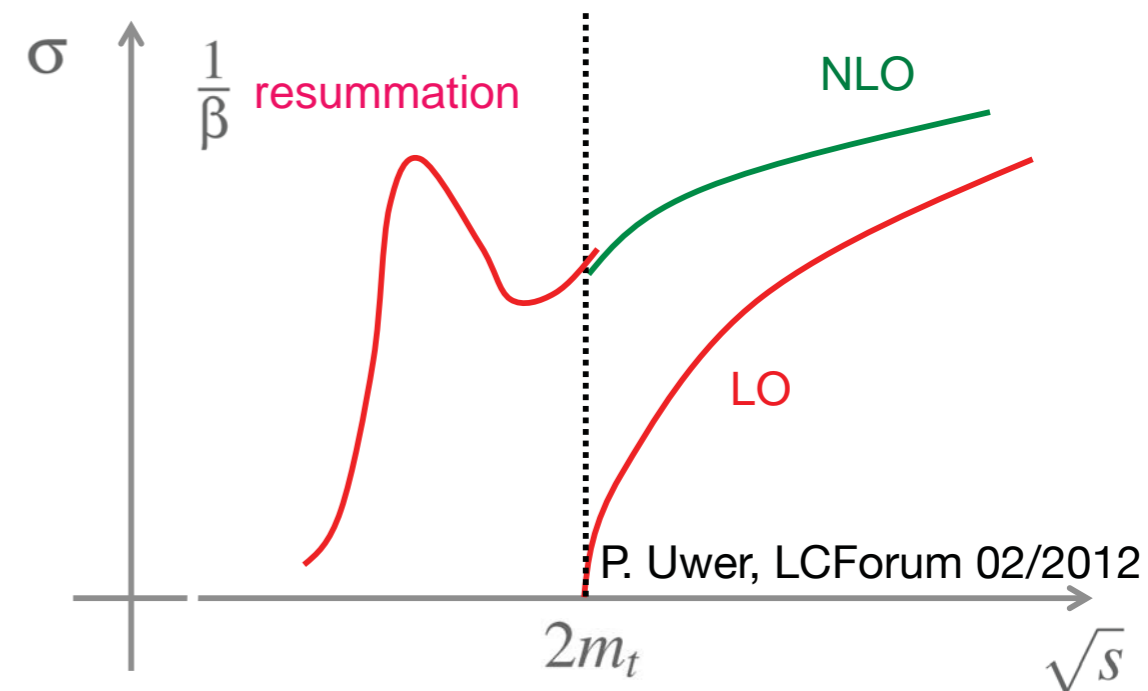
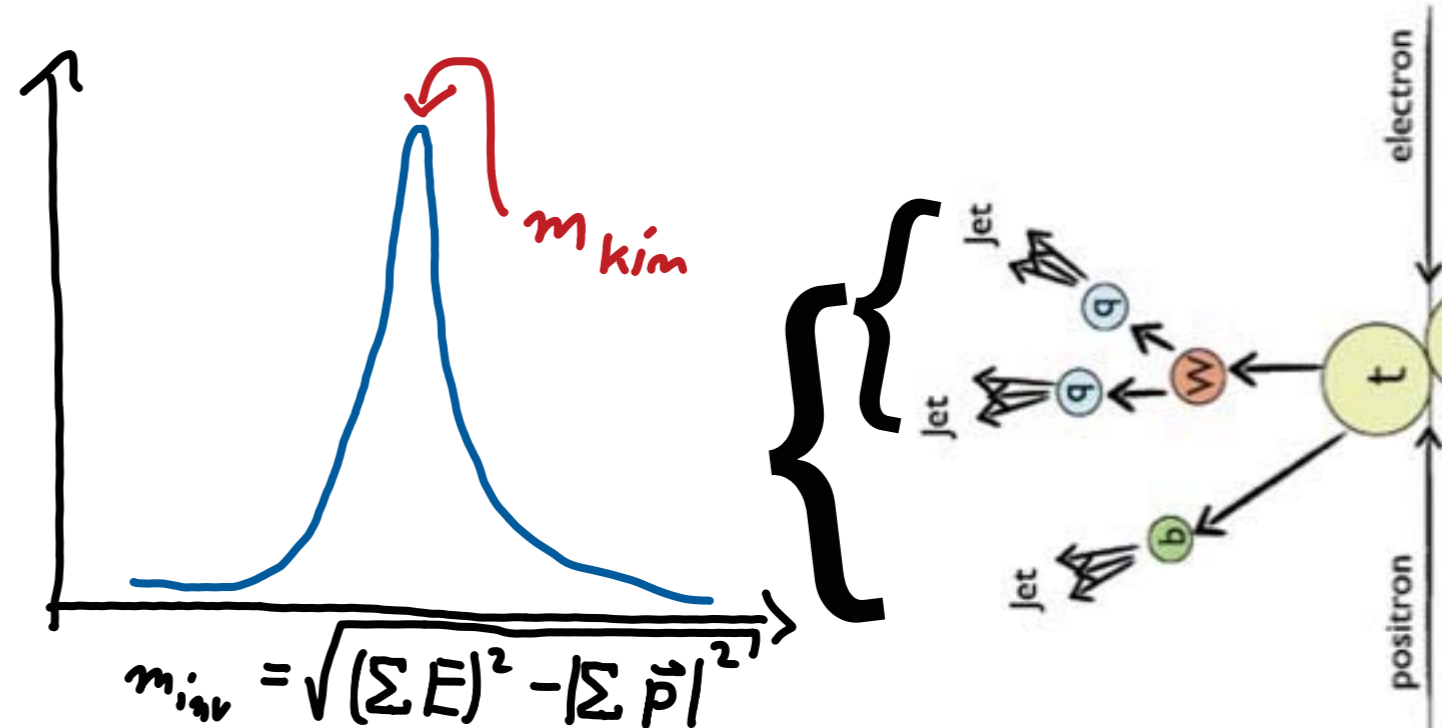
- Measurement in top pair production, two possibilities, each with advantages and disadvantages:

- Invariant mass

- experimentally well defined (but not theoretically: “PYTHIA mass”)
- can be performed at arbitrary energy above threshold: high integrated luminosity

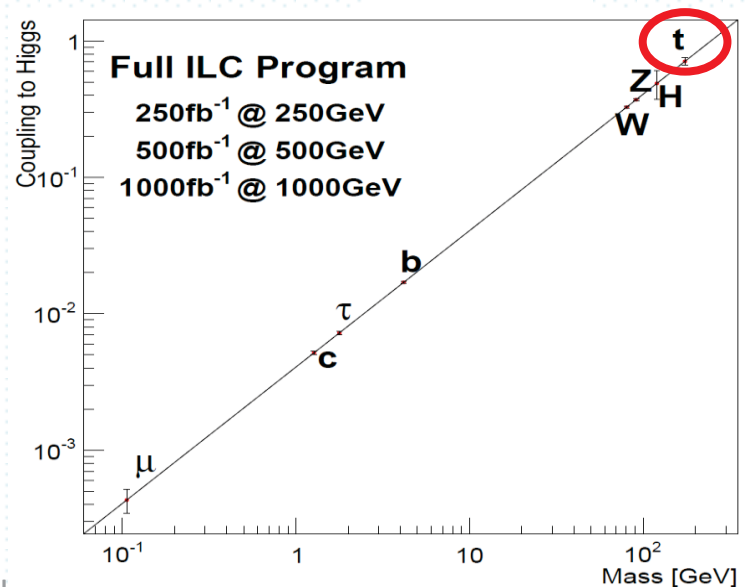
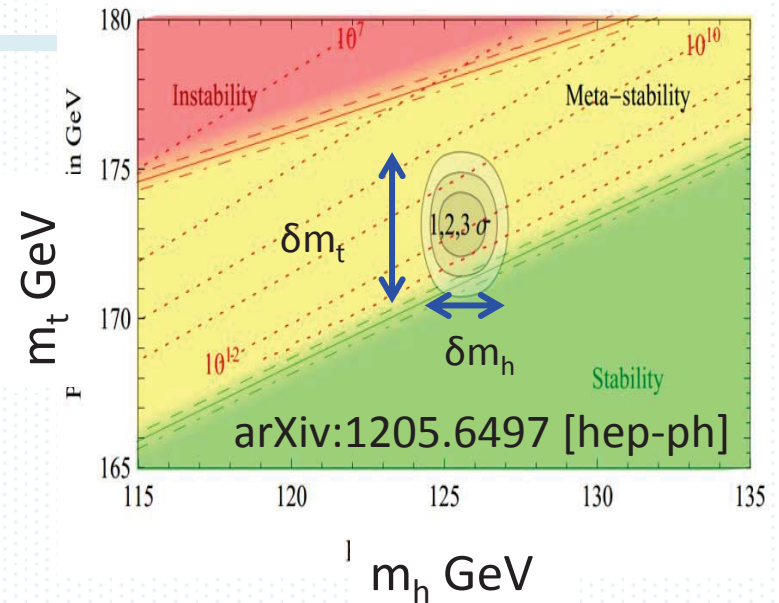
- Threshold scan

- theoretically well understood, can be calculated to higher orders
- needs dedicated running of the accelerator (but is also in a sweet spot for Higgs physics)
- ▶ The “ultimate” mass measurement at a LC!



Target around 350GeV

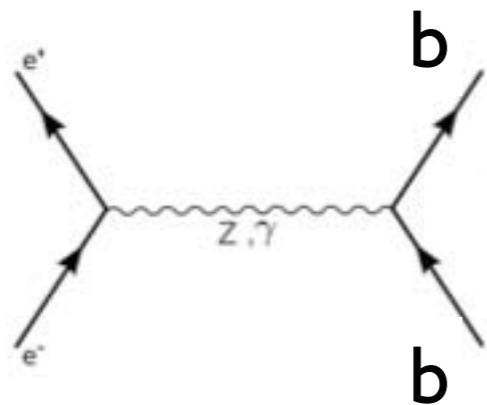
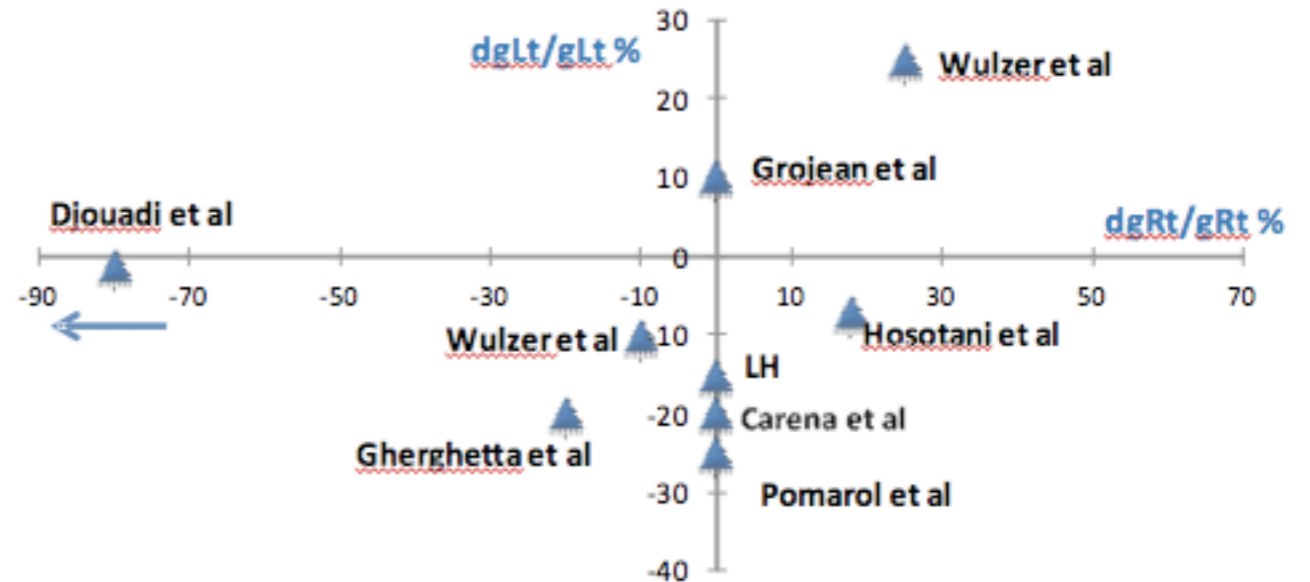
- **Top mass(m_t)**
 - Important input parameter
 - $\overline{\text{MS}}$ scheme mass ($m_t^{\overline{\text{MS}}}$)
 - ✓ $m_t^{\overline{\text{MS}}} = 160^{+5}_{-4}$ GeV (PDG)
 - Potential subtracted mass (m_t^{PS})
- **Decay width(Γ_t)**
 - anomalous coupling
 - exotic decay
- **Top yukawa coupling(y_t)**
 - Test of higgs mechanism
- α_s
- QCD wave function



Top-Z coupling

R. Rontsch
R. Poeschl

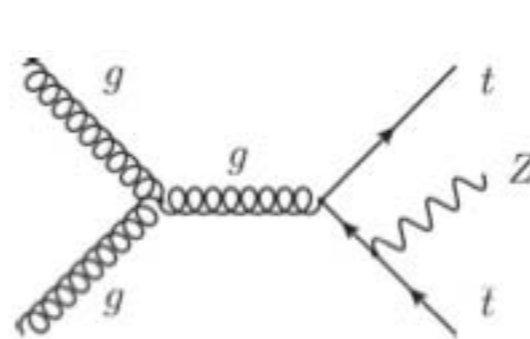
- Many models predict modifications in the coupling of the Z to top quarks.



LEP (+LC)

A_{FB} , constrain Zbb coupling, relate to Ztt through $SU(2)$

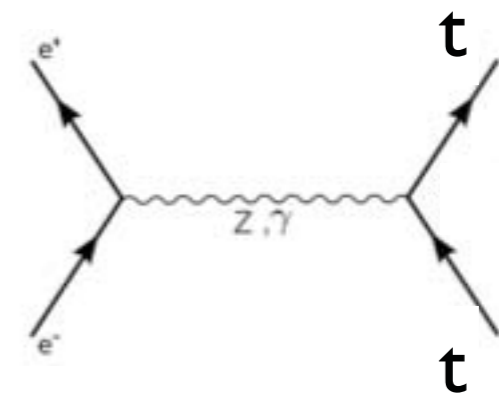
indirect



LHC

$\sigma(ttZ)$, rare process, previously limited by theory uncertainty

direct



future LC

A_{FB} , for precision must measure b-quark charge

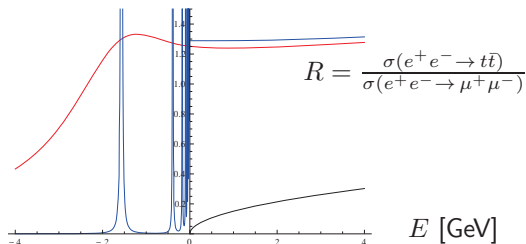
direct

Threshold Production

relative velocity of quark-antiquark pair is small:

$$\sqrt{s} = E + 2m_Q \approx 2m_Q \quad \Rightarrow \quad v = \sqrt{\frac{E}{m_Q}} \ll 1; \quad v \sim \alpha_s(m_Q v)$$

- multi-scale problem: mass m_Q , momentum $m_Q v$, energy $m_Q v^2$
- perturbation theory breaks down due to terms proportional to $\frac{\alpha_s}{v}$
 \rightsquigarrow Coulomb resummation
- formation of bound states below threshold
- $b\bar{b}$: bound-state resonances
- $t\bar{t}$: large width prevents existence of bound states



Framework: potential NRQCD

scales: mass, m : hard \gg momentum, mv : soft \gg energy, mv^2 : ultrasoft \gg Λ_{QCD}

$$\begin{array}{l} \text{potential quarks:} \\ \text{ultrasoft gluons:} \end{array} \left\{ \begin{array}{l} E_{\vec{p}} \sim mv^2 \\ |\vec{p}| \sim mv \\ E_{\vec{k}} \sim mv^2 \\ |\vec{k}| \sim mv^2 \end{array} \right. \quad \begin{array}{l} \frac{1}{E_{\vec{p}} - \frac{\vec{p}^2}{2m}} \\ \frac{1}{E_{\vec{k}}^2 - \vec{k}^2} \end{array}$$

QCD



NRQCD



pNRQCD (potential non-relativistic QCD)

↑
integrate out hard scale “ m ” from QCD [Caswell, Lepage’86;

Bodwin, Braaten, Lepage’95]

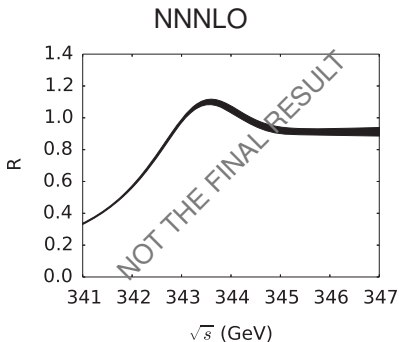
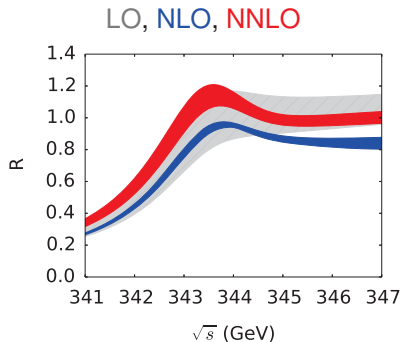
↑
integrate out all scales from NRQCD except potential quarks and ultrasoft gluons

[Beneke, Smirnov’97; Pineda, Soto’98; Brambilla, Pineda, Soto, Vairo’00]

alternative formulation: velocity NRQCD (vNRQCD)

[Luke, Manohar, Rothstein’00; Hoang, Stewart’03]

$$\sigma_{\text{tot}}(e^+e^- \rightarrow t\bar{t})$$

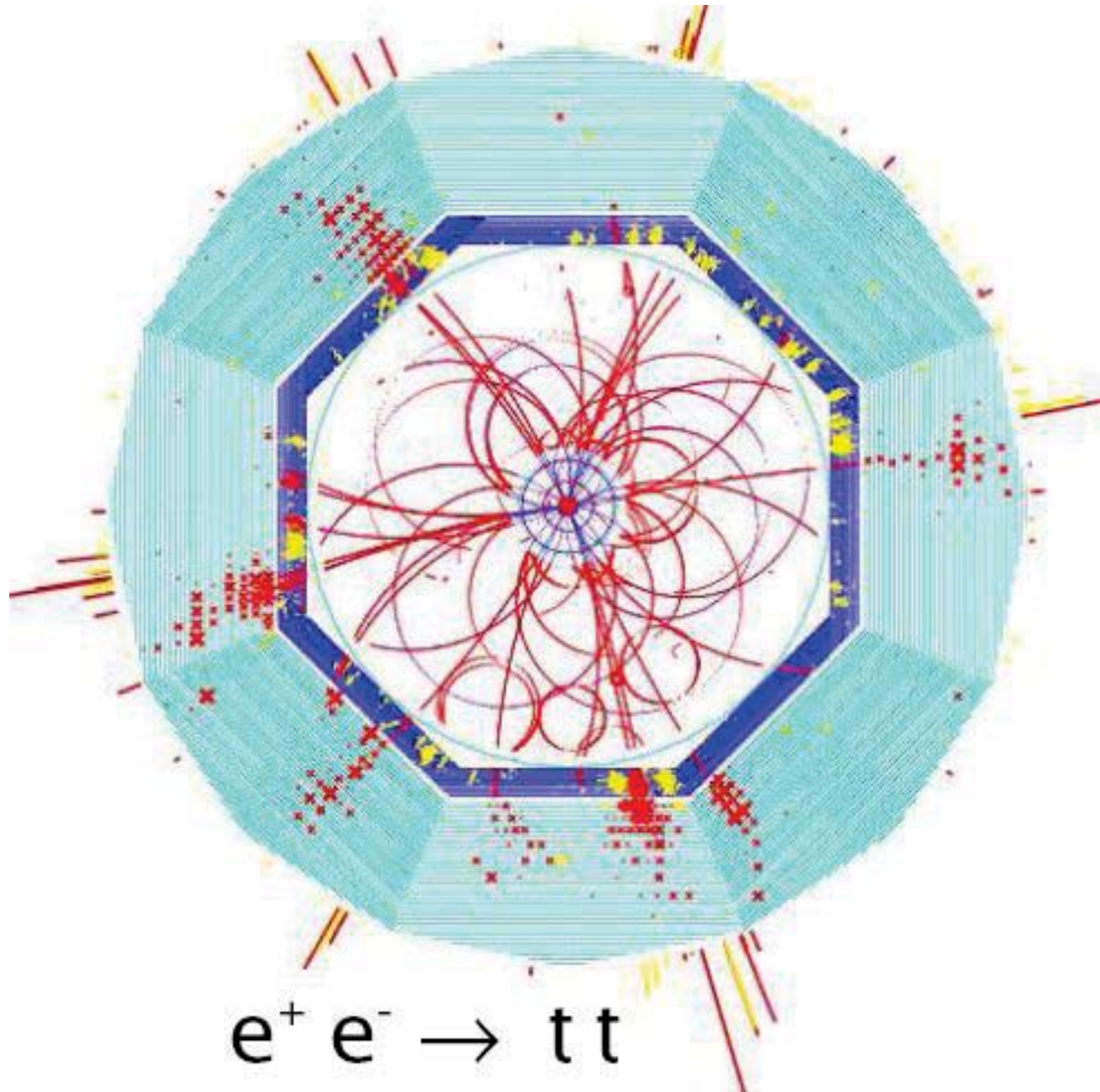


$$50 \text{ GeV} \leq \mu \leq 350 \text{ GeV}$$

$$m_t^{\text{PS}} = 171.3 \text{ GeV}$$

[Beneke et al.]

Future colliders and experiments

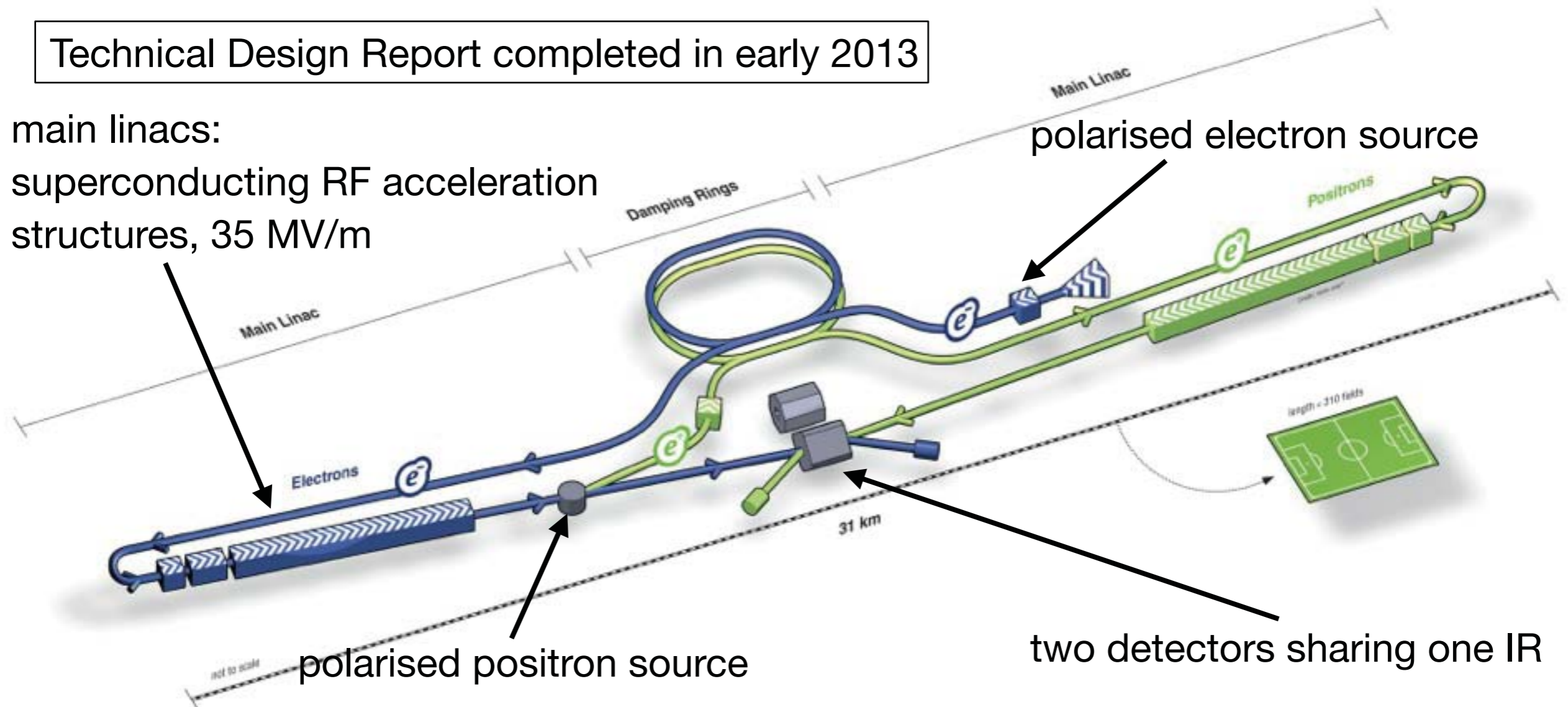


ILC - The International Linear Collider

- Currently the most advanced concept for a future energy frontier collider
 - e^+e^- collider, baseline energy 500 GeV, high luminosity: $2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
 - staged construction, starting from 250 GeV / 350 GeV
 - upgrade to 1 TeV possible (extension of linacs), luminosity upgrade by rate increase

Technical Design Report completed in early 2013

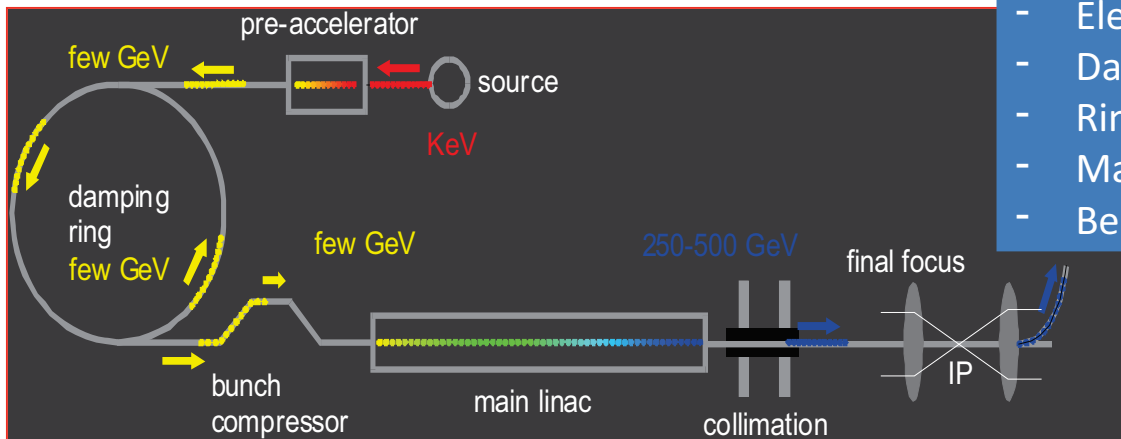
main linacs:
superconducting RF acceleration
structures, 35 MV/m





The ILC Accelerator Concept

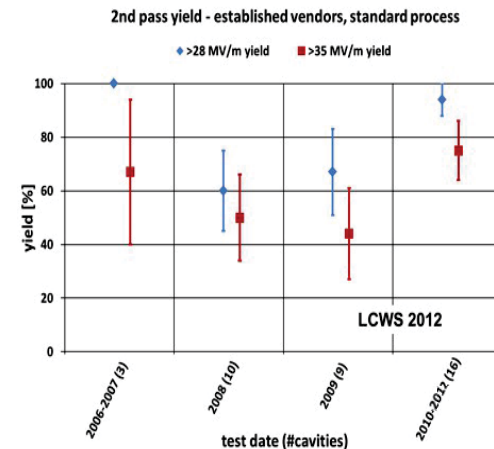
- Electron and Positron Sources (e^- , e^+)
- Damping Ring (DR)
- Ring to ML beam transport (RTML)
- Main Linac (ML) : SCRF Technology
- Beam Delivery System (BDS)



1.3 GHz Nb 9-cell Cavities	16,024
Cryomodules	1,855
SC quadrupole pkg	673
10 MW MB Klystrons & modulators	436

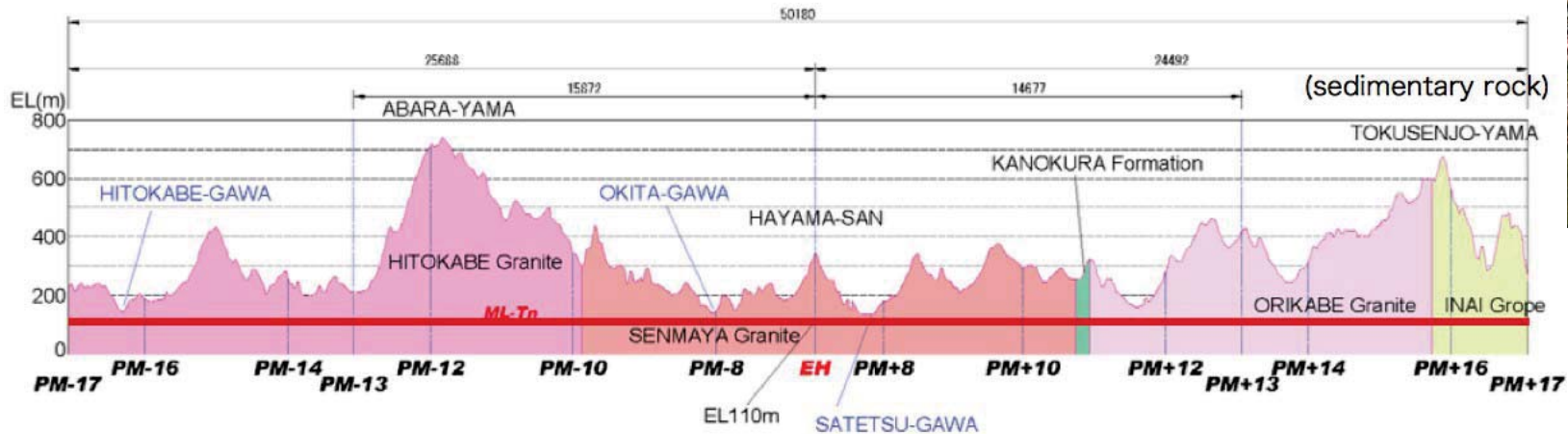


Production yield: 94 % at $> 35 \pm 20\%$
 Average gradient: 37.1 MV/m
 $>$ R&D goal of 35 MV/m reached (2012)





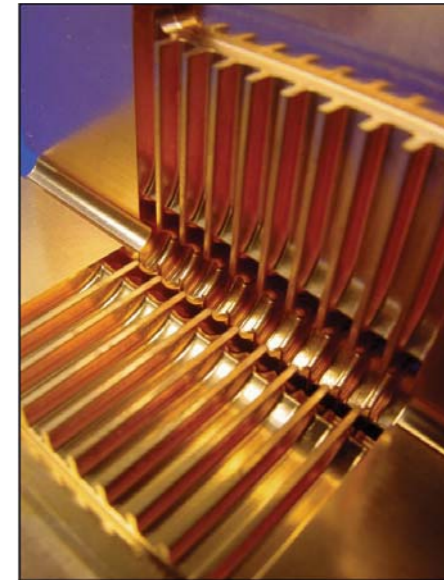
ILC preferred site - Kitakami



← 30km →

CLIC - The Compact Linear Collider

- A possible future energy frontier collider at CERN
 - e^+e^- collisions at up to 3 TeV with high luminosity ($\sim 6 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ at 3 TeV)
 - Staged construction 350 - 500 GeV, ~ 1.5 TeV, 3 TeV - detailed energies under study, based on physics and technical considerations
 - Based on two-beam acceleration: gradients of 100 MV/m
- Development phase until ~ 2018 - CDR completed in 2012



CLIC layout at 3 TeV

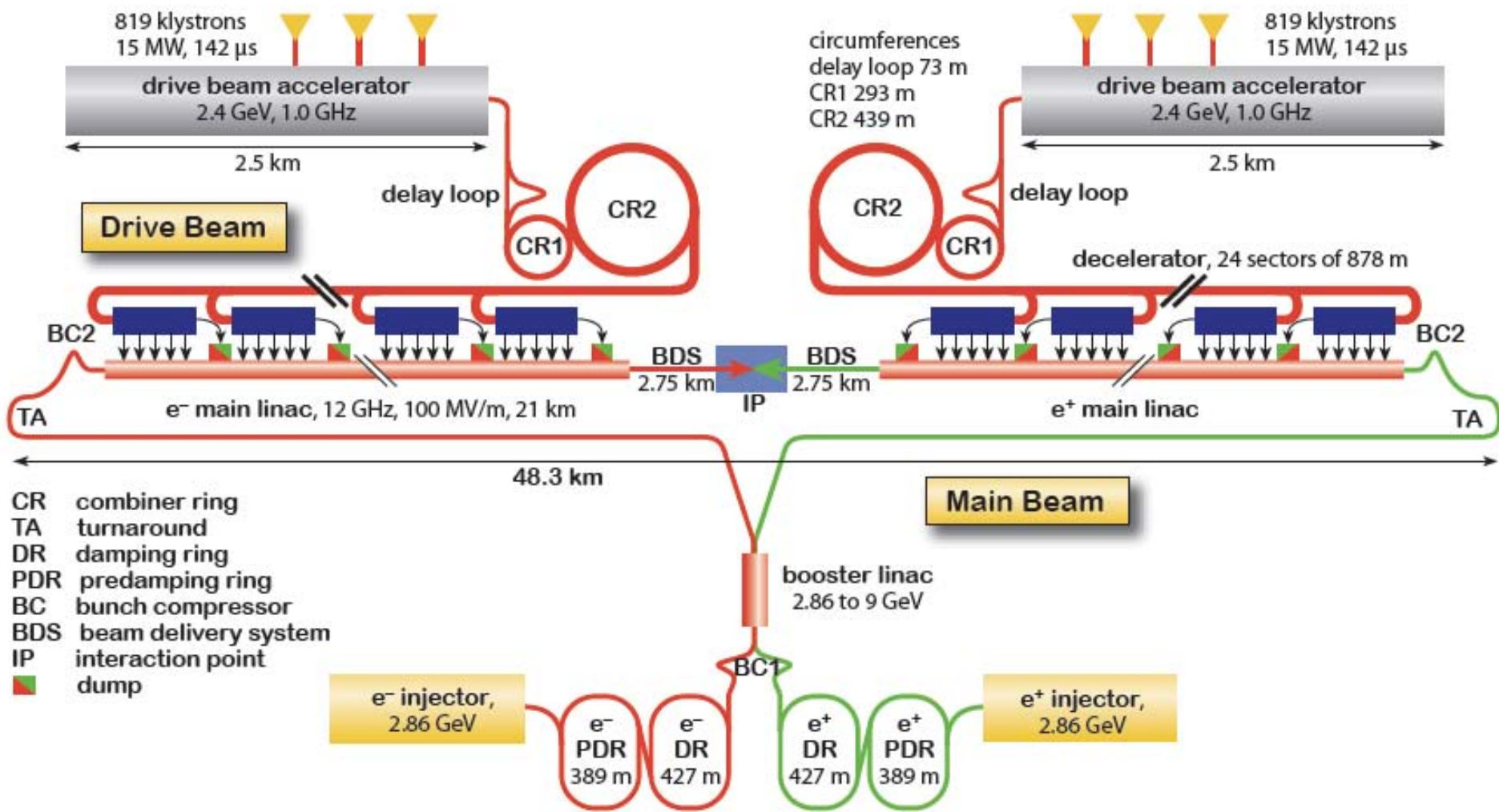


Fig. 3.1: Overview of the CLIC layout at $\sqrt{s} = 3$ TeV.

Future Circular Collider Study - SCOPE

CDR and cost review for the next ESU (2018)

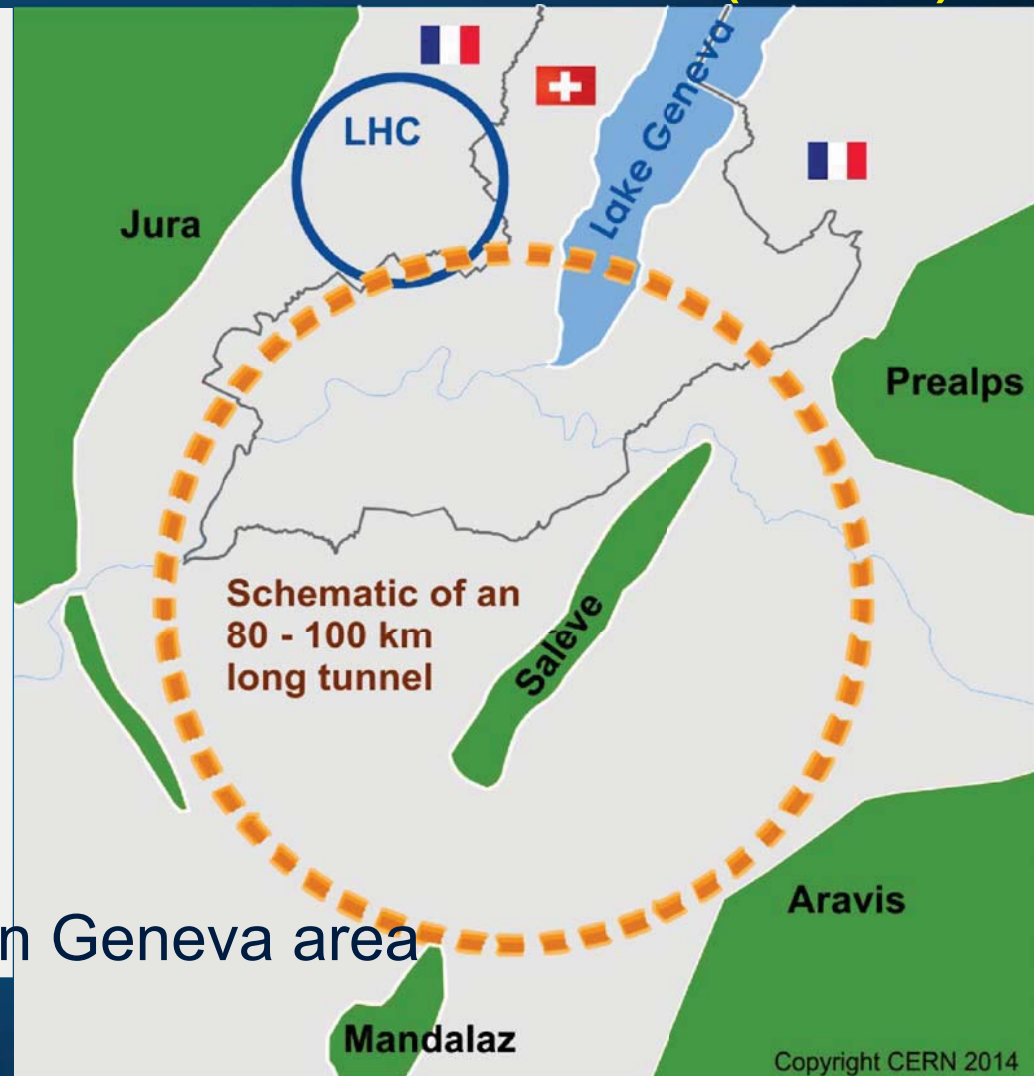
Forming an international collaboration to study:

- **pp -collider (*FCC-hh*)**
→ defining infrastructure requirements

~16 T \Rightarrow 100 TeV pp in 100 km

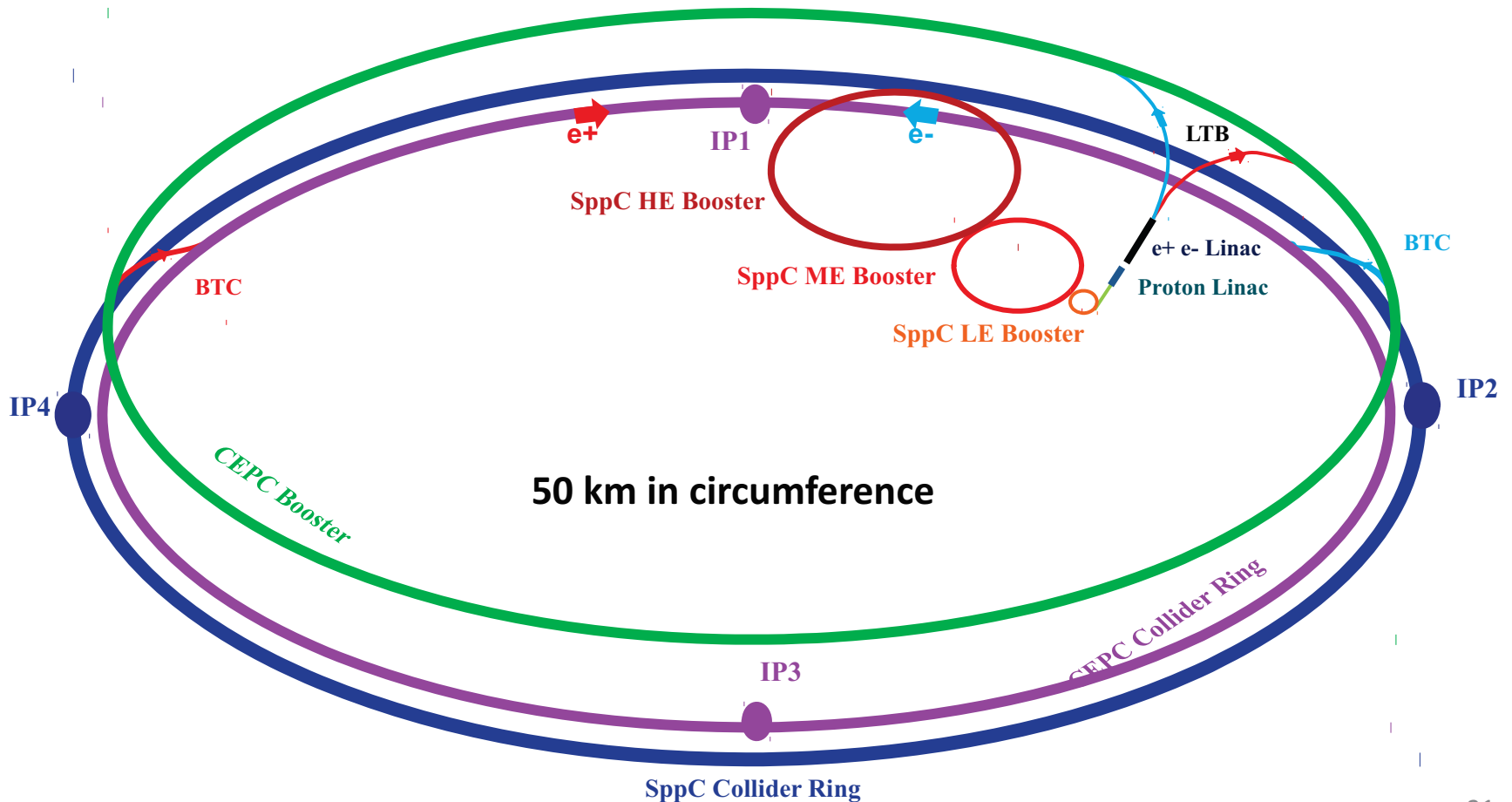
~20 T \Rightarrow 100 TeV pp in 80 km

- **e^+e^- collider (*FCC-ee*)** as potential intermediate step
- **p - e (*FCC-he*) option**
- **80-100 km infrastructure in Geneva area**

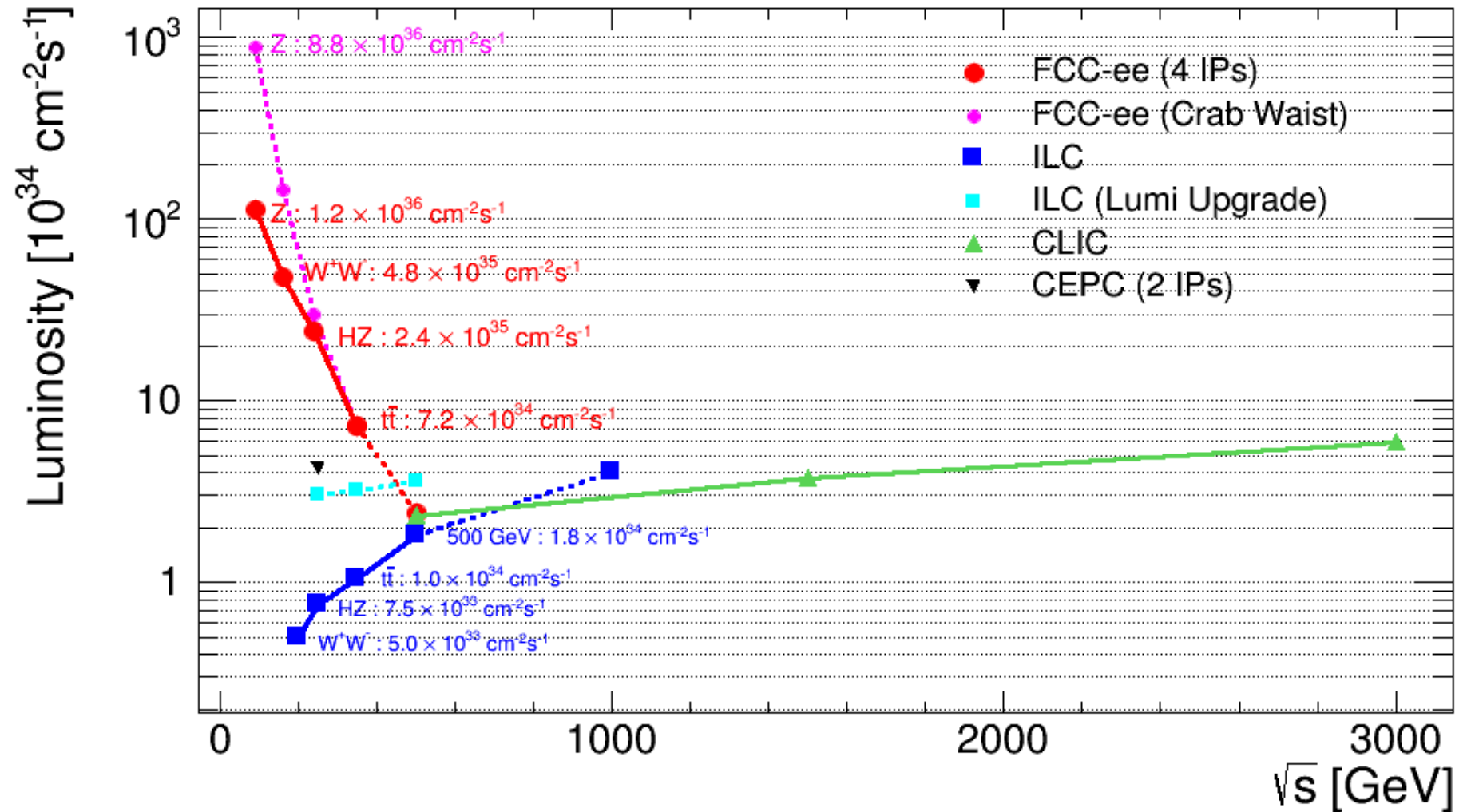


CEPC-SppC

CEPC is an 240 GeV Circular Electron Positron Collider, proposed to carry out high precision study on Higgs bosons, which can be upgraded to a 70 TeV or higher pp collider **SppC**, to study the new physics beyond the Standard Model.



Goal performance of e+ e- colliders



Identifying & Reconstructing Top Quarks

- Strategy depends on targeted $t\bar{t}$ final state

Semi-leptonic:

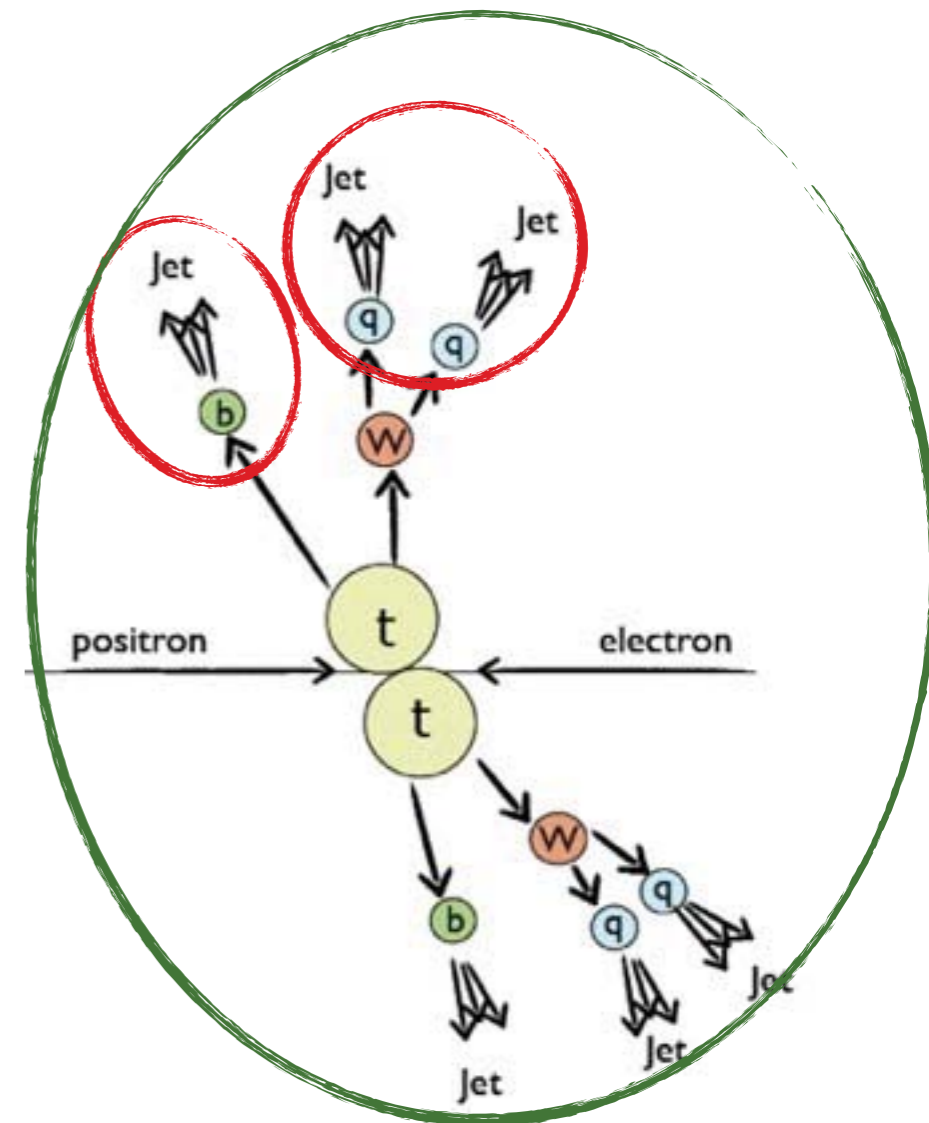
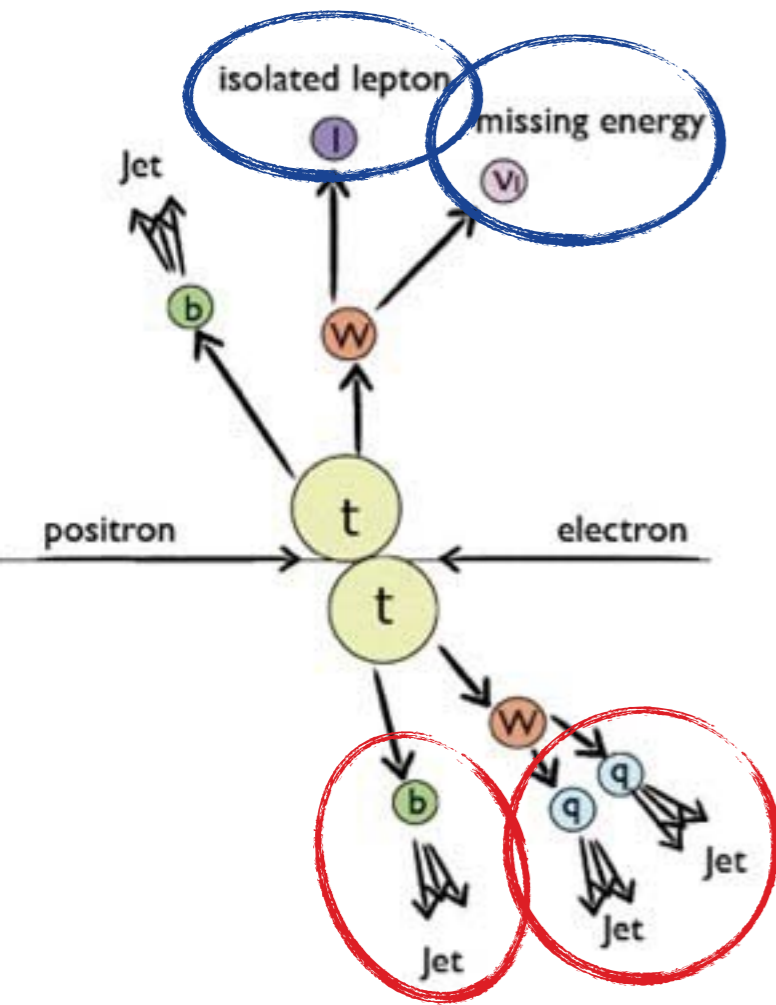
- isolated lepton ID, momentum measurement
- provides t / \bar{t} identification
- missing energy measurement

Universal

- Flavor tagging:
 - b - identification
 - b/c separation
- b -Jet energy measurement
- light Jet reconstruction & energy measurement

All-hadronic

- global hadronic energy reconstruction



LC Detector Requirements

Track Momentum: $\sigma_{1/p} < 5 \times 10^{-5} / \text{GeV}$ **(1/10 x LEP)**

(e.g. Z-Mass Measurement with charged Leptons)

Impact parameter: $\sigma_{d_0} < 5 \oplus 10 / (p[\text{GeV}] \sin^{3/2} \theta) \mu\text{m}$ **(1/3 x SLD)**

(c/b-tagging)

Jet energy : $dE/E = 3\text{-}4\%$

(Measurement of W/Z Mass with Jets)

Hermeticity : $\theta_{\min} = 5 \text{ mrad}$

(to detect of events with missing energy e.g. SUSY)

Events with large track multiplicity and a large number of Jets (6+) are expected.

Therefore:

- **high Granularity**
- **good track Measurement**
- **good Track Separation**
- **"Particle Flow" detectors**

$e^+ e^- \rightarrow t \bar{t}$

Jet energy resolution and background rejection drive the overall detector design

=> => fine-grained calorimetry + Particle Flow Analysis (PFA)

What is PFA?

Typical jet composition:
60% charged particles
30% photons
10% neutral hadrons



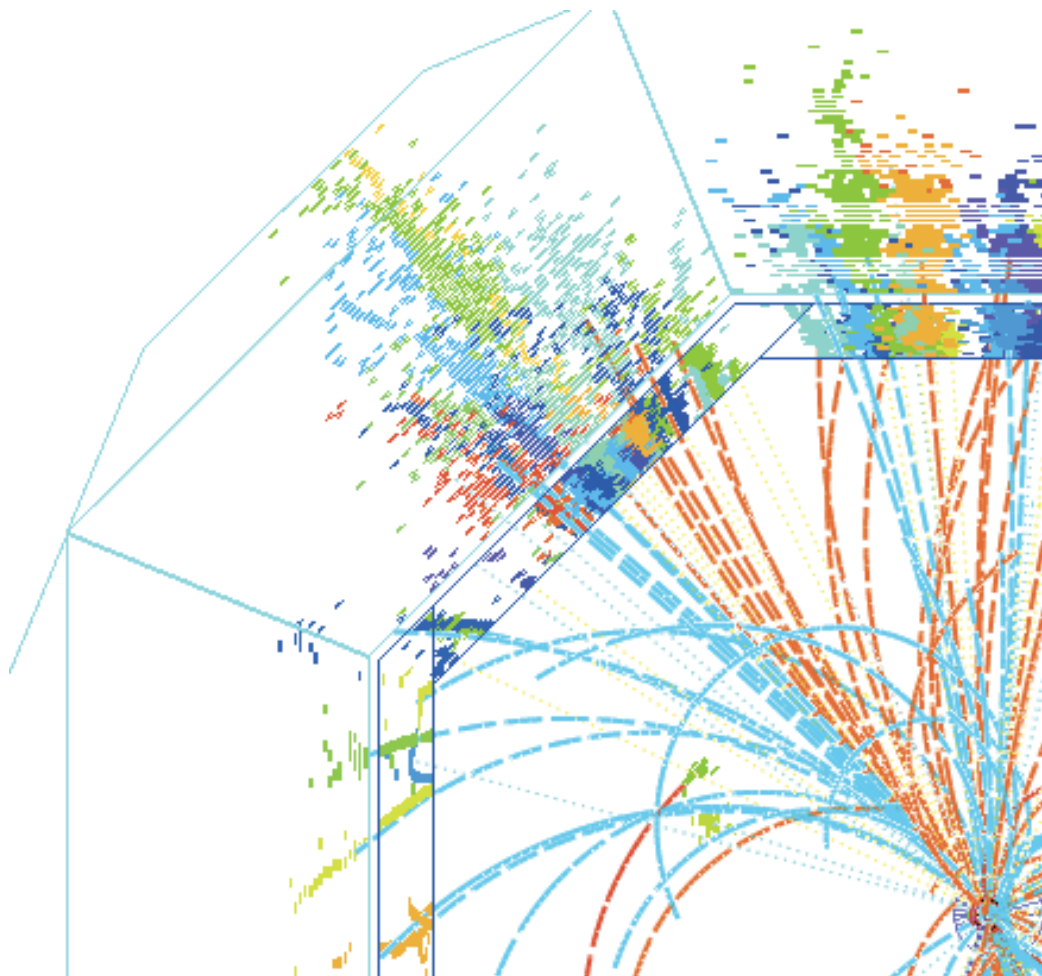
Always use the best info you have:

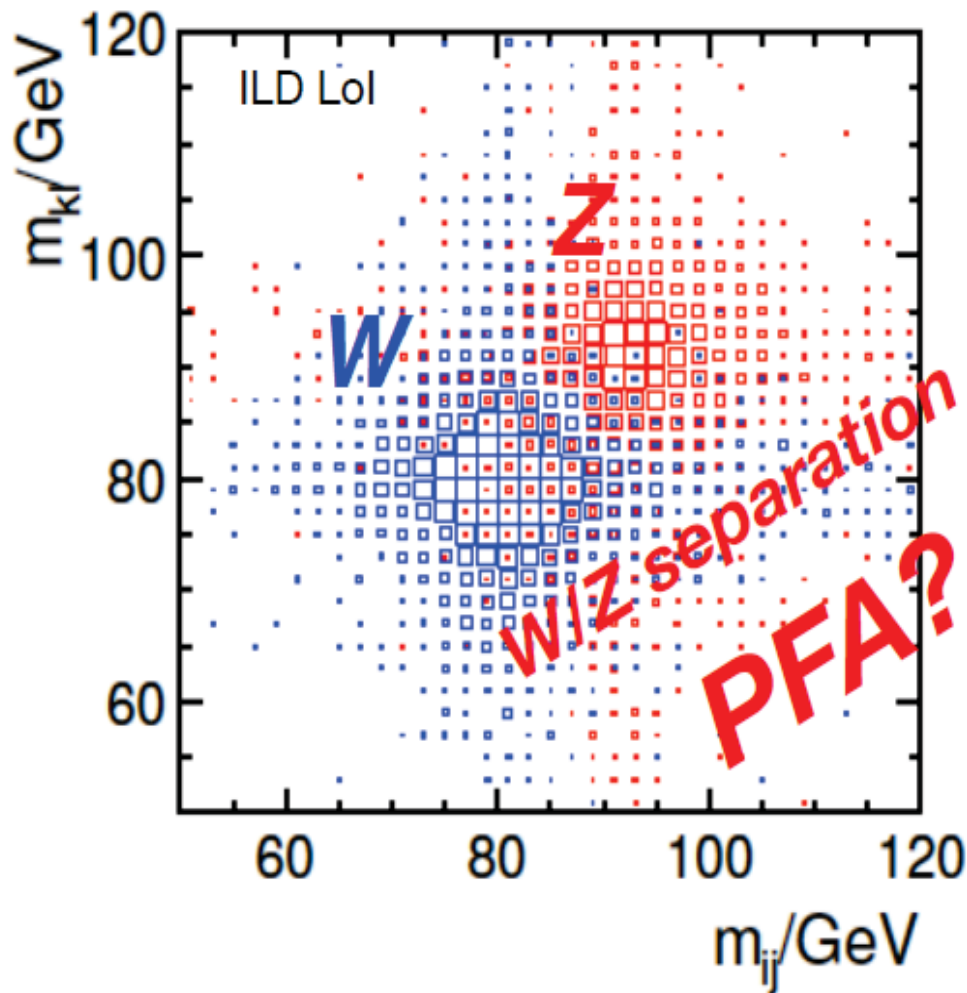
60% => tracker 😊 😊

30% => ECAL 😊

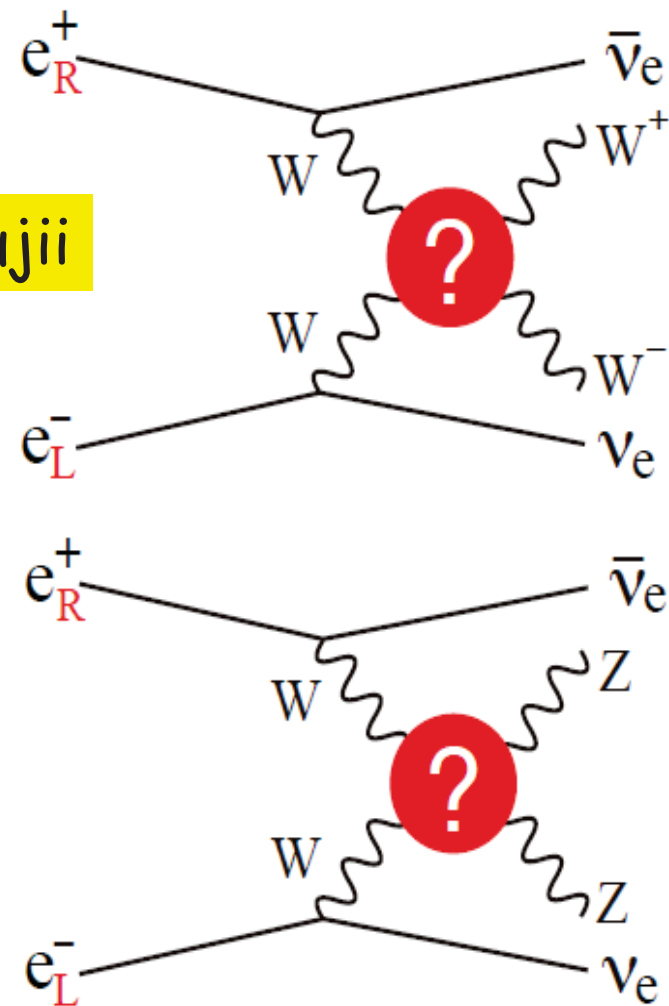
10% => HCAL 😞

Hardware + software !

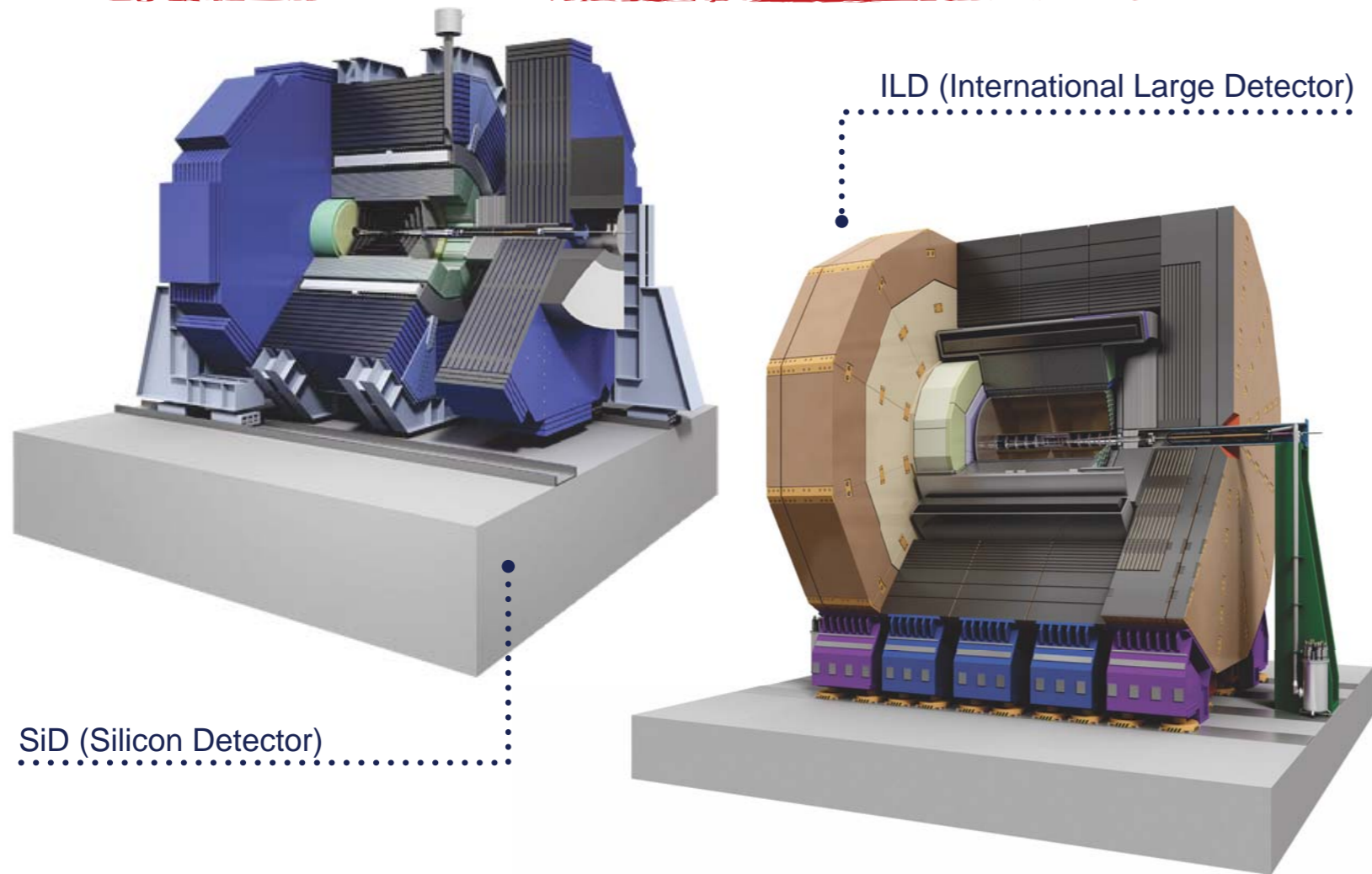




K. Fujii



Detector Systems at Linear Colliders

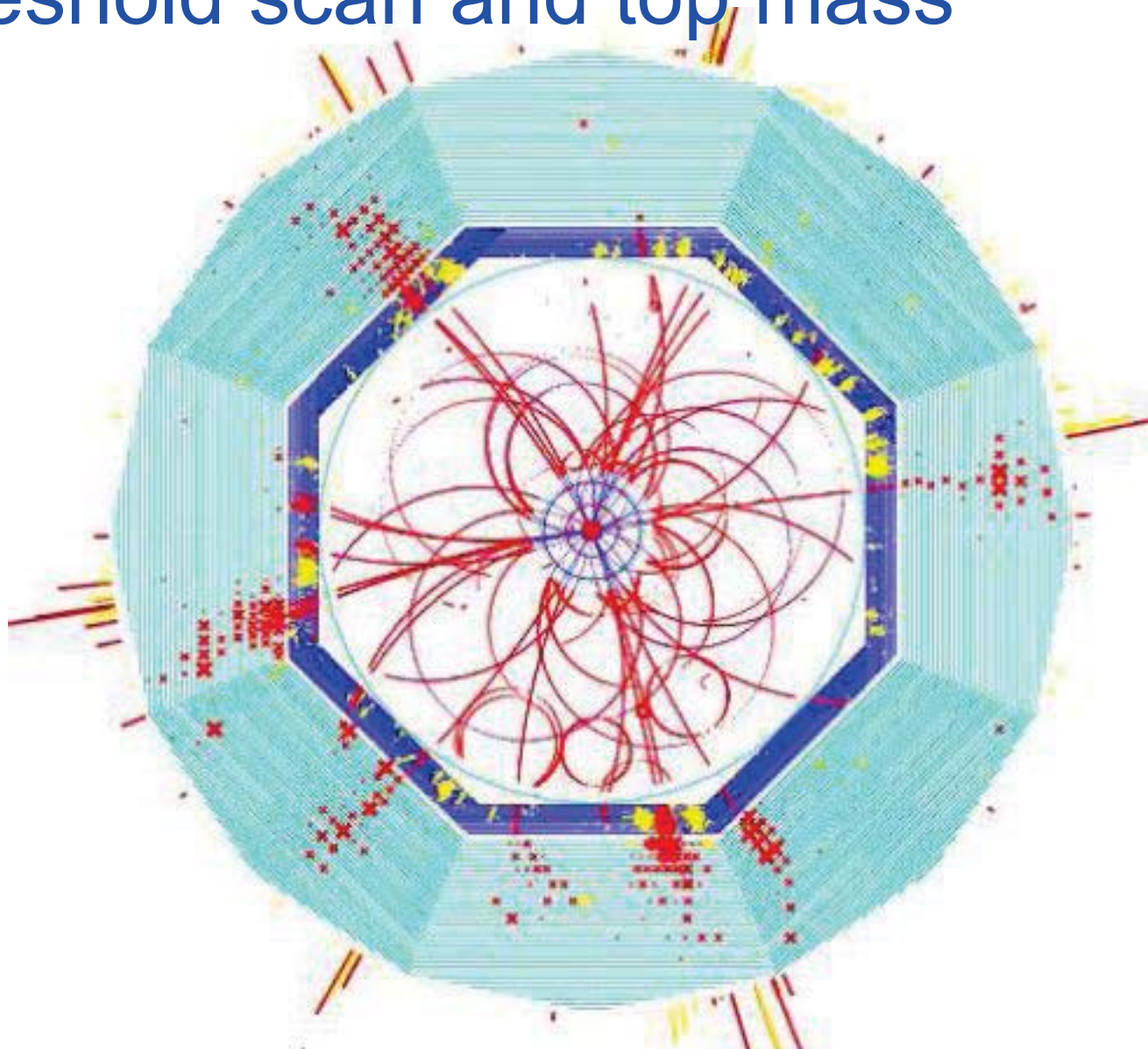


- Low-mass, high precision vertexing & tracking
- Highly granular calorimeters
- Particle flow event reconstruction

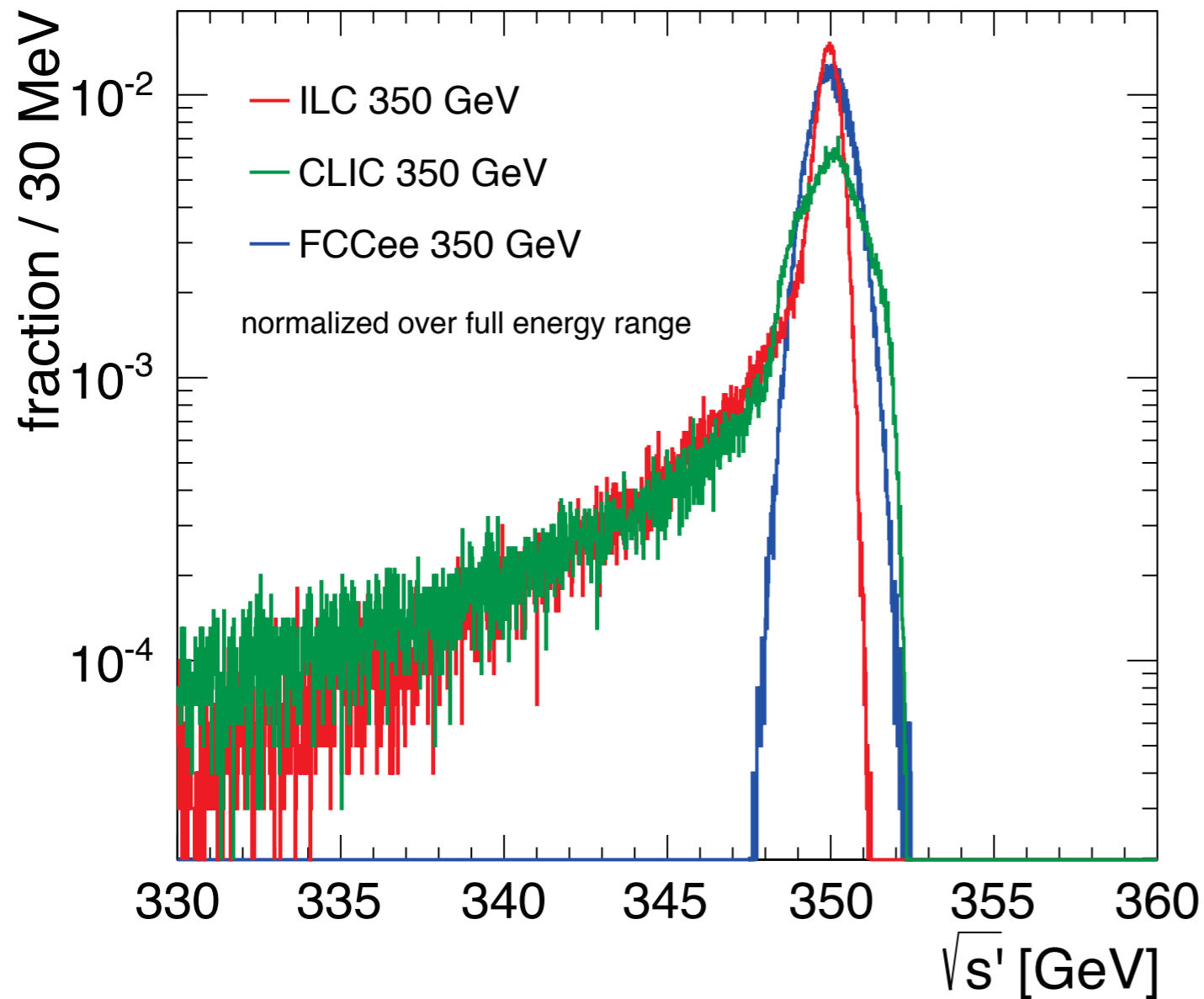
- CLIC detectors based on ILC concepts, with modifications in the calorimeters, vertex and forward regions to account for higher energy and higher backgrounds
- Detailed simulation models implemented in GEANT4
- Realistic event reconstruction including pattern recognition, tracking, PFA
- ▶ Full simulation studies used for all results presented here

Top measurements

- Threshold scan and top mass

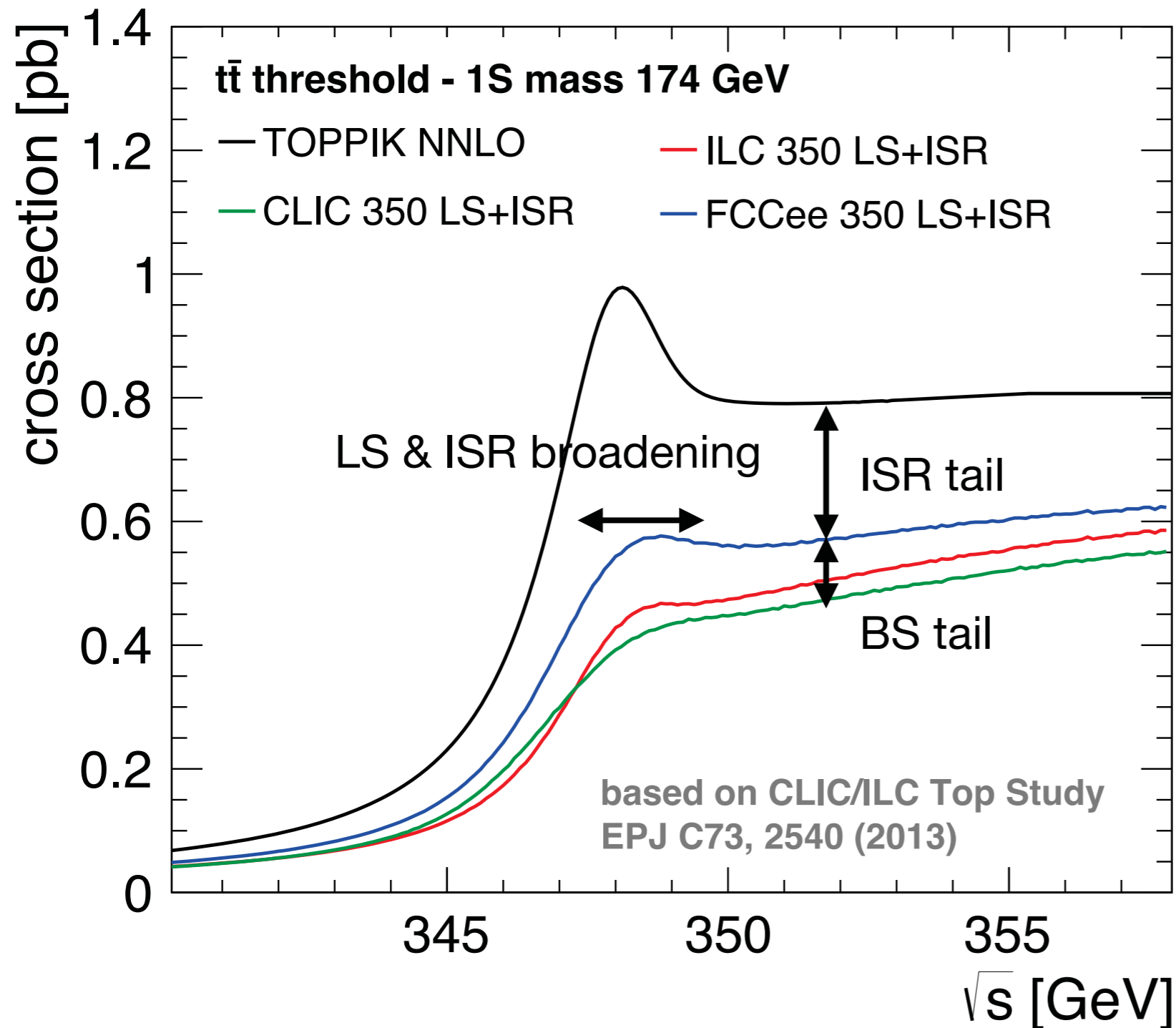


The Luminosity Spectrum of different Colliders...



- The luminosity spectrum of different e^+e^- colliders
 - ILC & CLIC - Full machine simulations (GuneaPig)
 - FCCee (TLEP): Gaussian, with 0.19% sigma (includes BS)

... and what it does to the Top Threshold



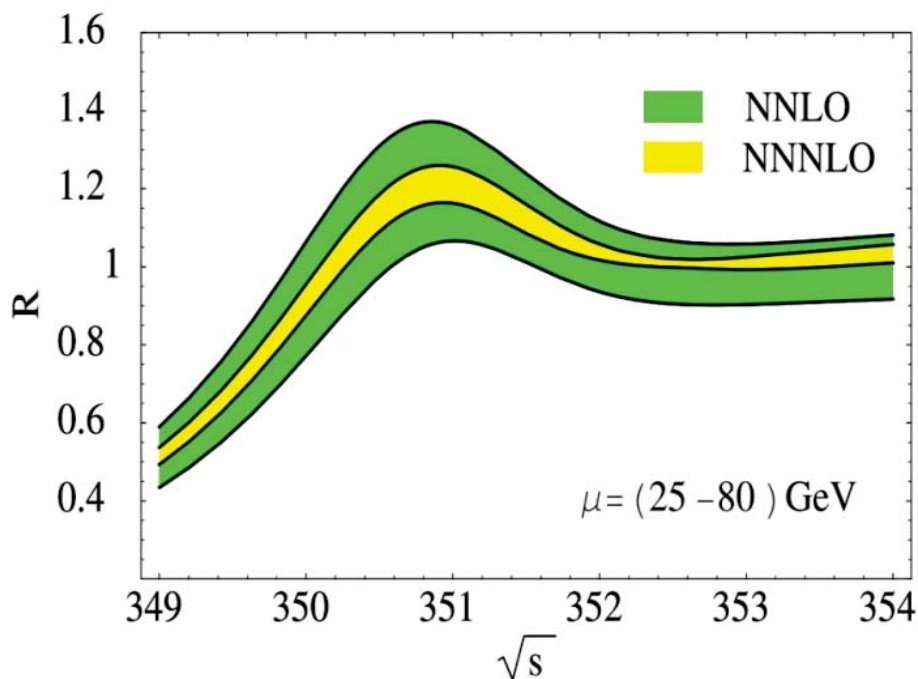
The effects:

- ISR tail: lowering of effective L at top energy
- BS tail: lowering of effective L at top energy - not at FCCee - Gaussian spectrum, 100% of L at $> 99\%$
- LS & ISR broadening: smearing of Xsection due to beam energy spread, BS tail and ISR - most pronounced at CLIC - comparable at ILC and FCCee

σ_{tt} Measurement

Near the threshold region of top pair production ($\sqrt{s}=2m_t$), the energy dependence of σ_{tt} is large. And σ_{tt} depend on fundamental parameters. Then, using threshold scan technic, measuring σ_{tt} precisely and fitting it, these parameters are determined !!

$$\sigma_{tt} \propto f(\sqrt{s}, m_t, \Gamma_t, \alpha_s, y_t, m_h)$$



$\sigma/\delta\sigma(\text{theoretical}) \sim 4\text{-}5 \%$

arXiv:0801.3464 [hep-ph]

M. Beneke, Y. Kiyo and K. Schuller,

Signal and background

Signal

6-Jet

4-Jet

Branching Ratio

6-Jet

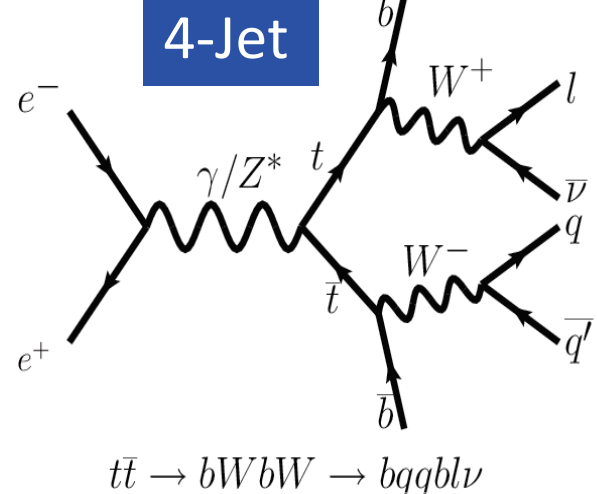
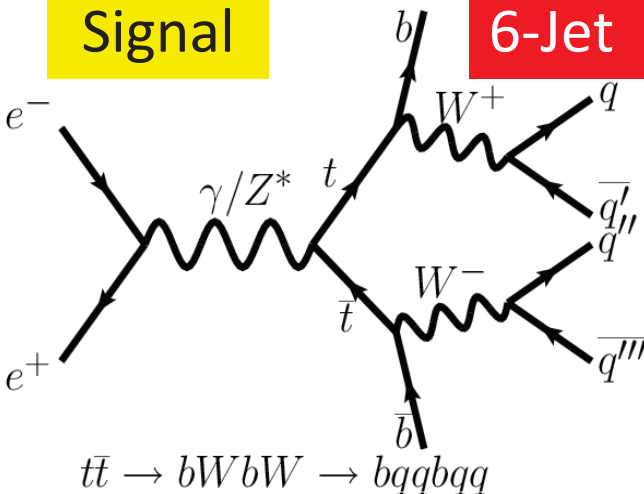
45%

4-Jet

44%

2-Jet

11%

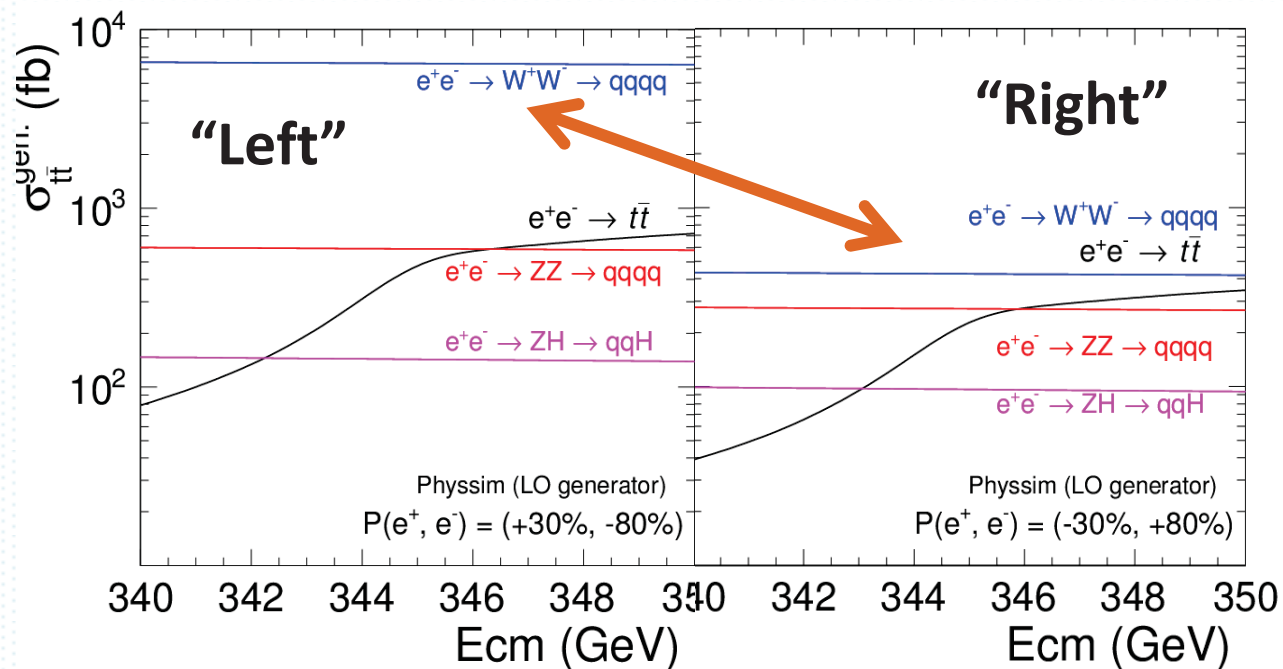
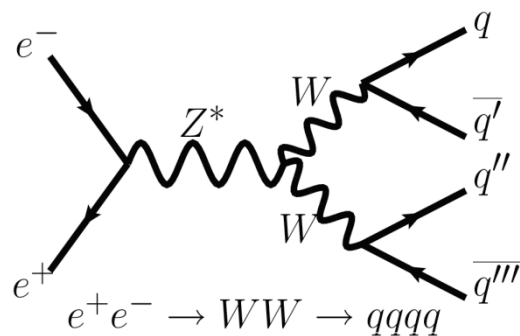


6-Jet	45%
4-Jet	44%
2-Jet	11%

background

SM bkg. which have 4 or 6 fermions in final state

Main bkg. : WW, ZZ, ZH



Top Quark Reconstruction (6-Jet & 4-Jet)

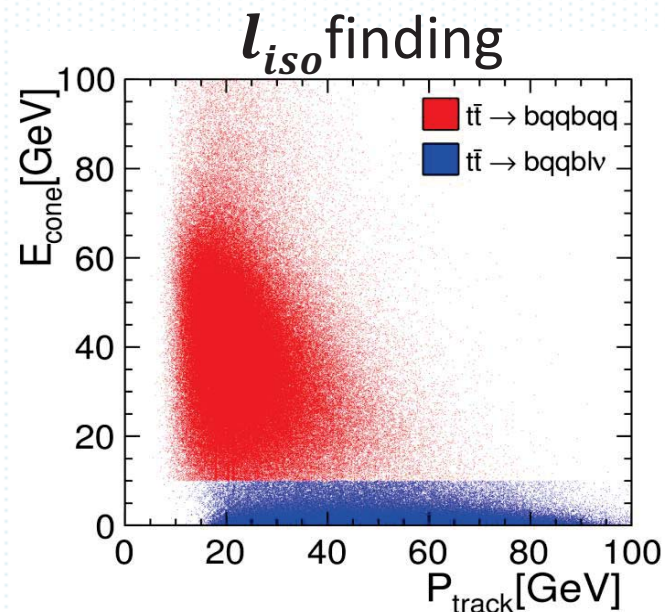
Reconstruction method	6-Jet	4-Jet
Suppressing the background overlay using anti- k_T algorithm ($R=0.7$)		
Isolated Lepton(l_{iso}) finding using cone energy cut	# of $l_{iso} = 0$	# of $l_{iso} = 1$
Jet clustering using Durham algorithm	Cluster to 6jets	Cluster to 4jets
2 b-likeness Jets were found using LCFIPlus	-	-
Reconstruction of two W bosons	$q_1 + q_2$ & $q_3 + q_4$	q_1+q_2 & $l_{iso}+v$
Reconstruction of two top quarks	-	-
Minimizing the χ^2	①	②

①

$$\chi_{6\text{-Jet}}^2 = \frac{(m_{3j^a\text{reco.}} - m_t)^2}{\sigma_t^2} + \frac{(m_{3j^b\text{reco.}} - m_t)^2}{\sigma_t^2} + \frac{(m_{2j^a\text{reco.}} - m_w)^2}{\sigma_w^2} + \frac{(m_{2j^b\text{reco.}} - m_w)^2}{\sigma_w^2}$$

②

$$\chi_{4\text{-Jet}}^2 = \frac{(m_{3j\text{reco.}} - m_t)^2}{\sigma_t^2} + \frac{(m_{jl\nu\text{reco.}} - m_t)^2}{\sigma_t^2} + \frac{(m_{2j\text{reco.}} - m_w)^2}{\sigma_w^2}$$



Selection Table 6-Jet @350GeV

$$\int \mathcal{L}(t) dt = 5(\text{fb}^{-1}) \quad S = \frac{N_{Sig}}{\sqrt{N_{Sig} + N_{BG}}}$$

Left	tt6j	tt4j	tt2j	WW	ZZ	ZH	6f+4f	S_{6j}
Generated	1643	1583	381	32664	3004	694	71691	4.9
# of lepton = 0	1592	357	19	32079	2957	638	39983	5.7
btag > 0.1 × 2	1515	340	18	3601	1398	471	7399	12.5
Thrust < 0.84	1485	313	13	398	433	383	1084	23.2
Evis > 290 GeV	1481	159	1	218	310	309	90	29.2
missPt < 38 GeV	1473	72	0	217	307	303	80	29.7
$m_{\tau} > 100 \text{ GeV} \times 2$	1467	69	0	180	253	255	63	30.7
$y_{45} > 0.0015$								
$y_{56} > 0.0007$	1419	45	0	68	71	80	36	34.2
# of pfos > 86	1406	38	0	45	59	73	33	34.6

S/N

5.67

$\delta\sigma/\sigma$

2.9%

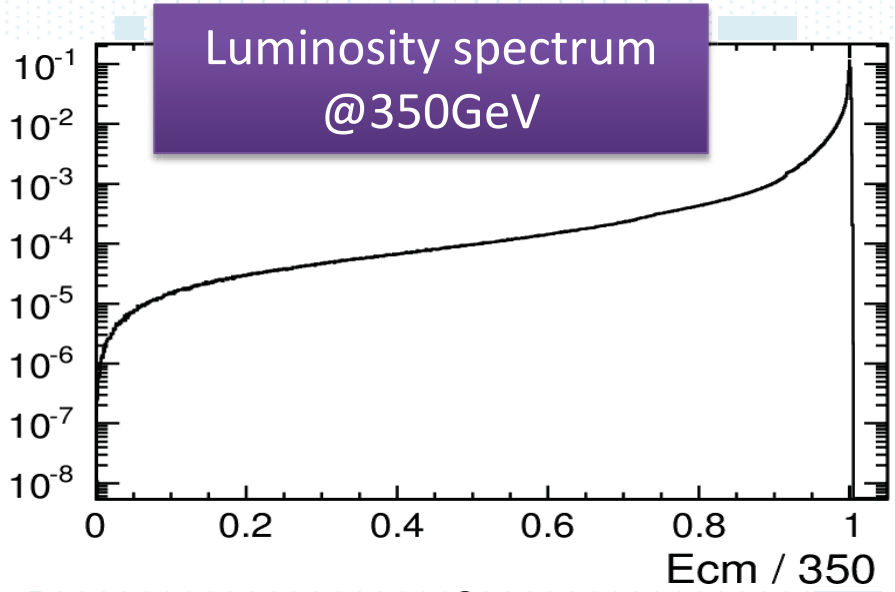
6f: 6 fermion final state except ttbar

Fit - convolution -

© We must consider **“Beam effects”** around threshold.

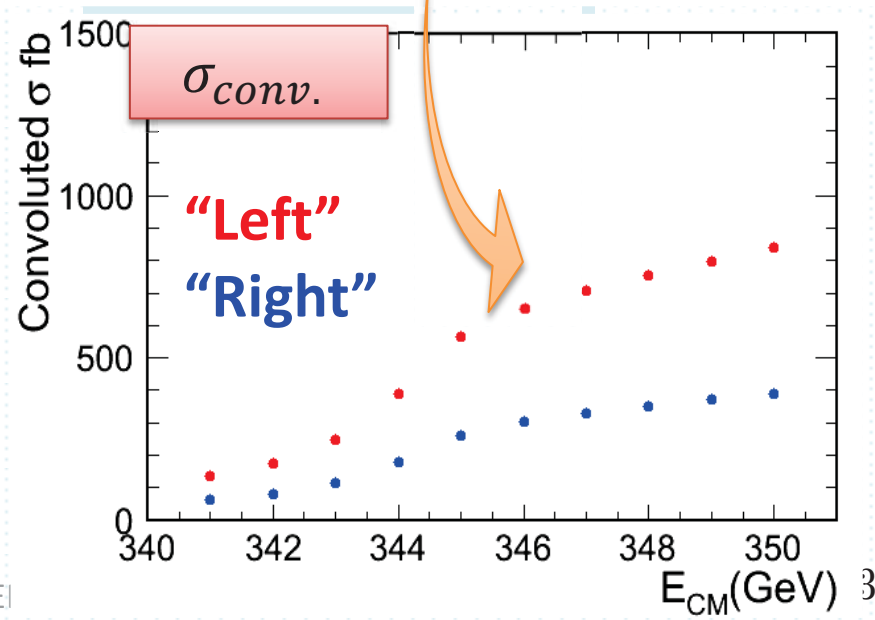
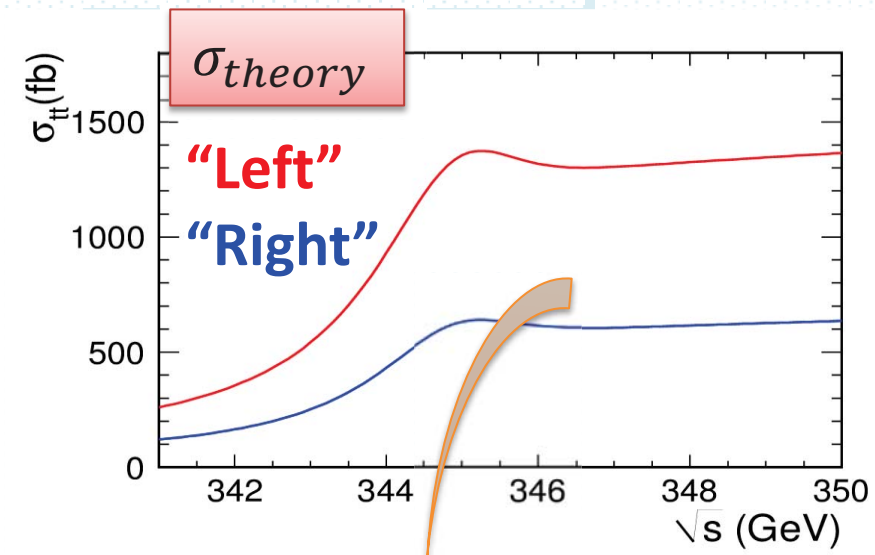


Using luminosity spectrum,
theoretical cross section is convoluted.



$$\sigma_{conv.}(\sqrt{s}) = \int \mathcal{L}(t) \sigma_{th}(t) dt$$

\mathcal{L} : luminosity spectrum, \sqrt{s} : nominal, σ_{th} : theoretical σ ,
 $\sigma_{conv.}$: convoluted σ , $t(=\sqrt{s'/s})$ where $\sqrt{s'}$ is collision energy



Fit -Result-

Stat. Error (MeV)	6-Jet		4-Jet	
	m_t^{PS}	Γ_t	m_t^{PS}	Γ_t
Left(50fb ⁻¹)	28	40	33	48
Right(50fb ⁻¹)	42	63	48	67
Left (50fb ⁻¹) + Right(50fb ⁻¹)	23	34	27	39

Center value: $m_t^{PS} = 172.000$ GeV, $\Gamma_t = 1.400$ GeV

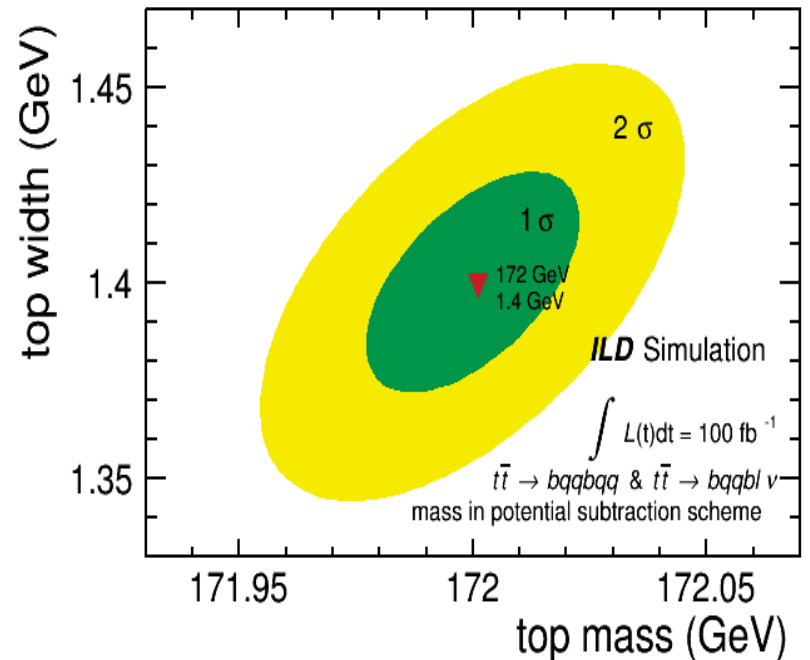
Combined ALL

m_t^{PS} (GeV)	Γ_t (GeV)
172.001 ± 0.018	1.399 ± 0.026

PS → MS

$$m_t^{\overline{MS}} \sim m_t^{PS} - \frac{4}{3\pi} (m_t^{PS} - 20) \alpha_s + \dots$$

$$m_t^{\overline{MS}} = 163.800 \pm 0.017 \text{ (stat.) (GeV)}$$

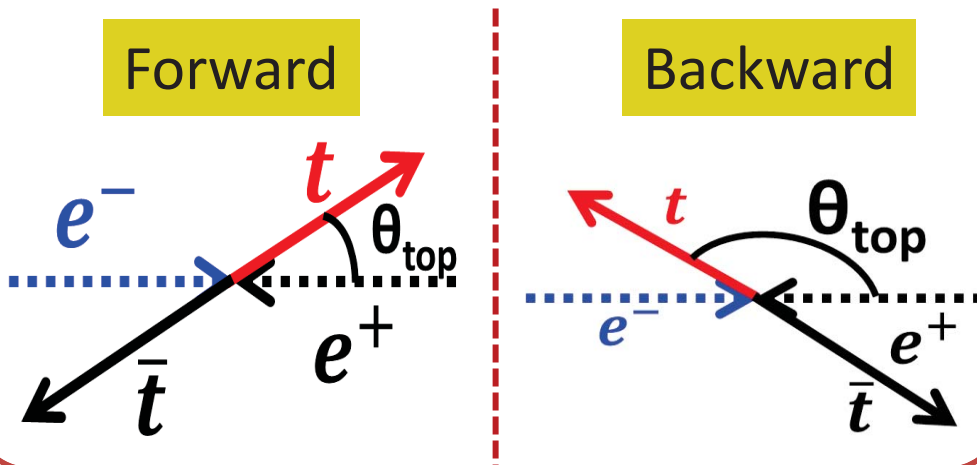


A_{FB} near $t\bar{t}$ threshold

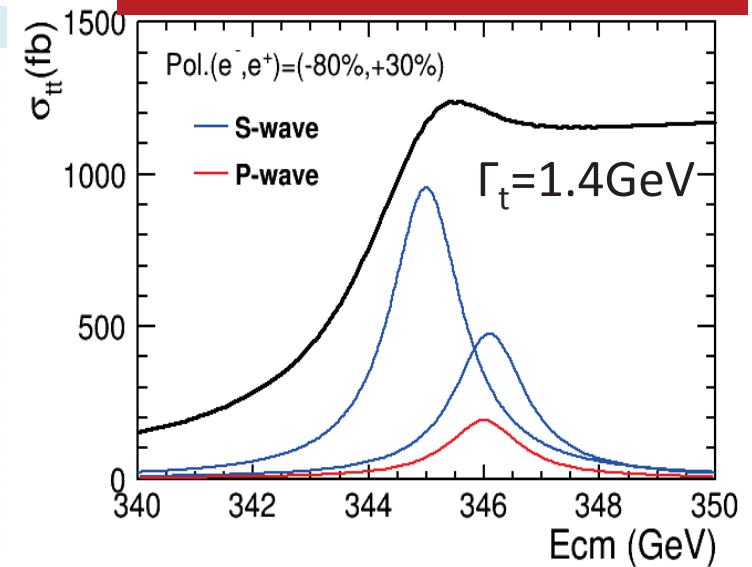
★ Forward backward asymmetry of top quark (A_{FB})

- Since top has large Γ_t , we can measure A_{FB} by interfering the resonance of S- and P- wave.
- The level split which is separation of two resonances depends on α_s .

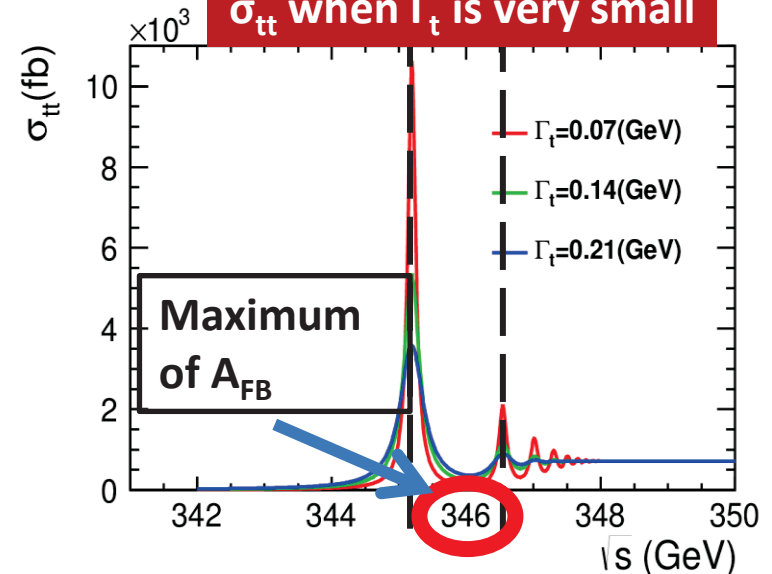
$$A_{FB} \equiv \frac{N(\cos\theta_{top} > 0) - N(\cos\theta_{top} < 0)}{N(\cos\theta_{top} > 0) + N(\cos\theta_{top} < 0)}$$



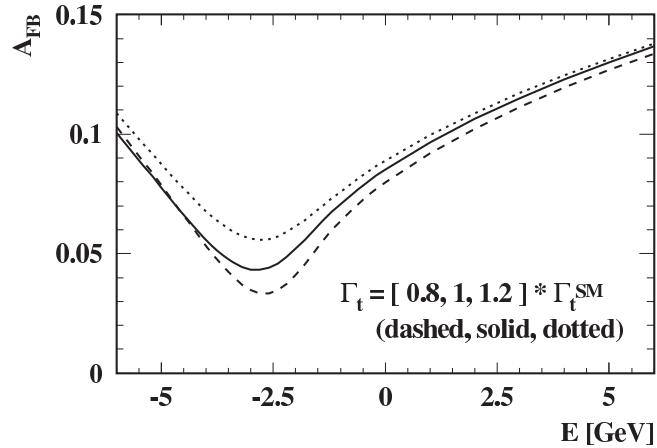
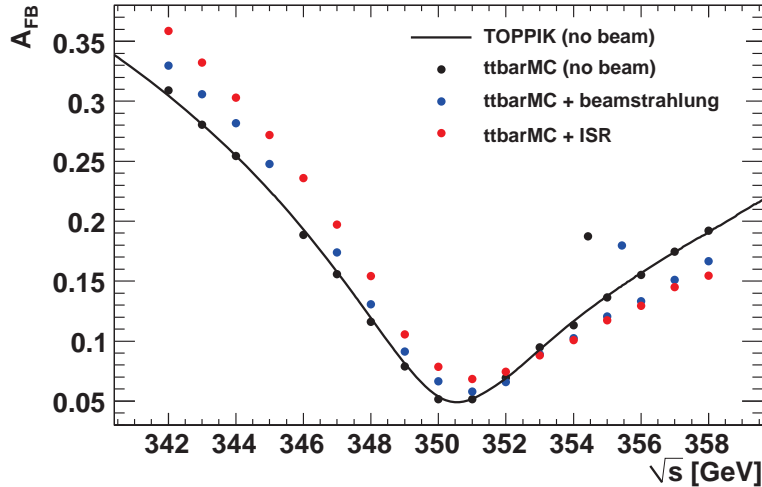
Interference of S- and P-wave



σ_{tt} when Γ_t is very small

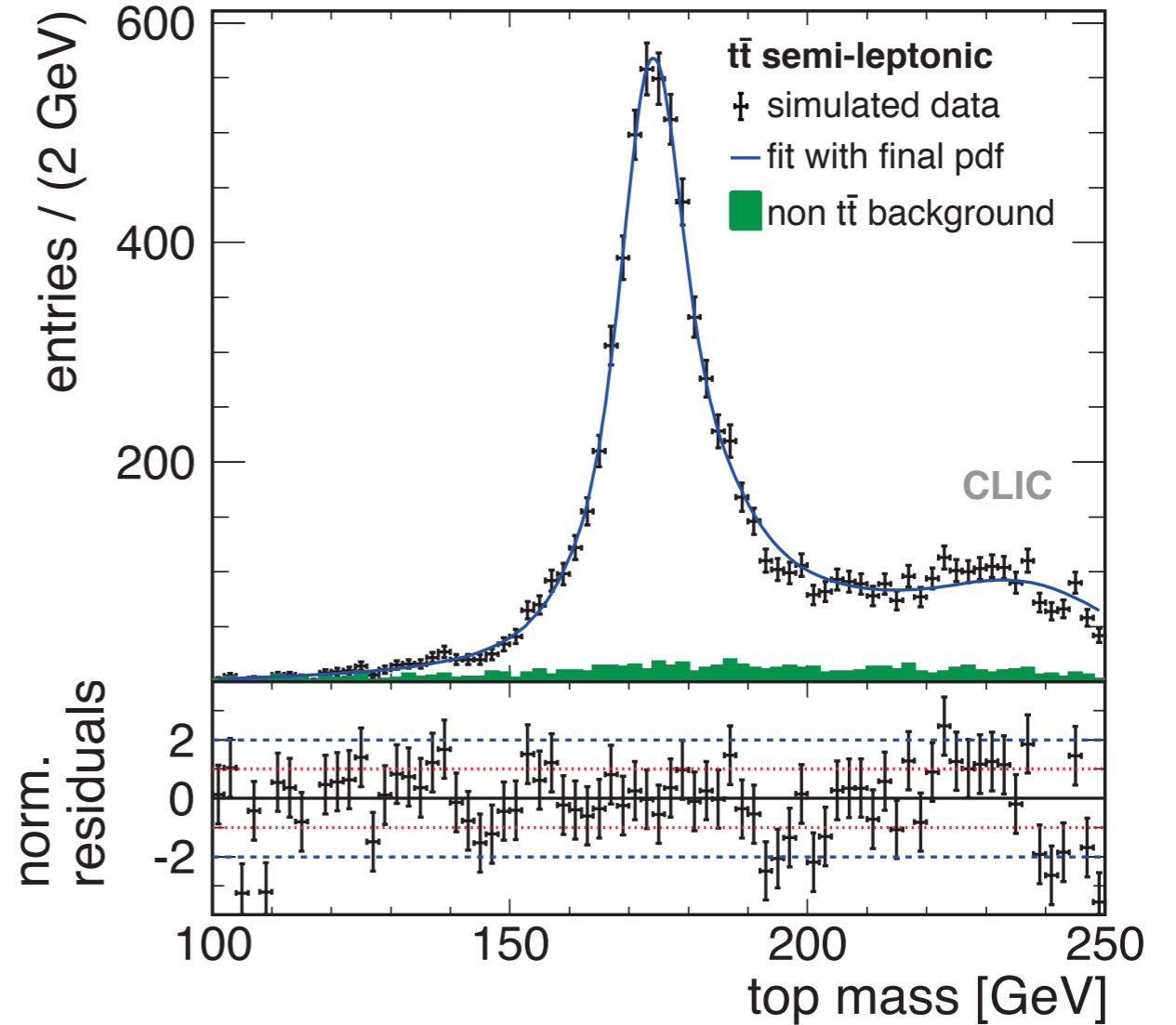
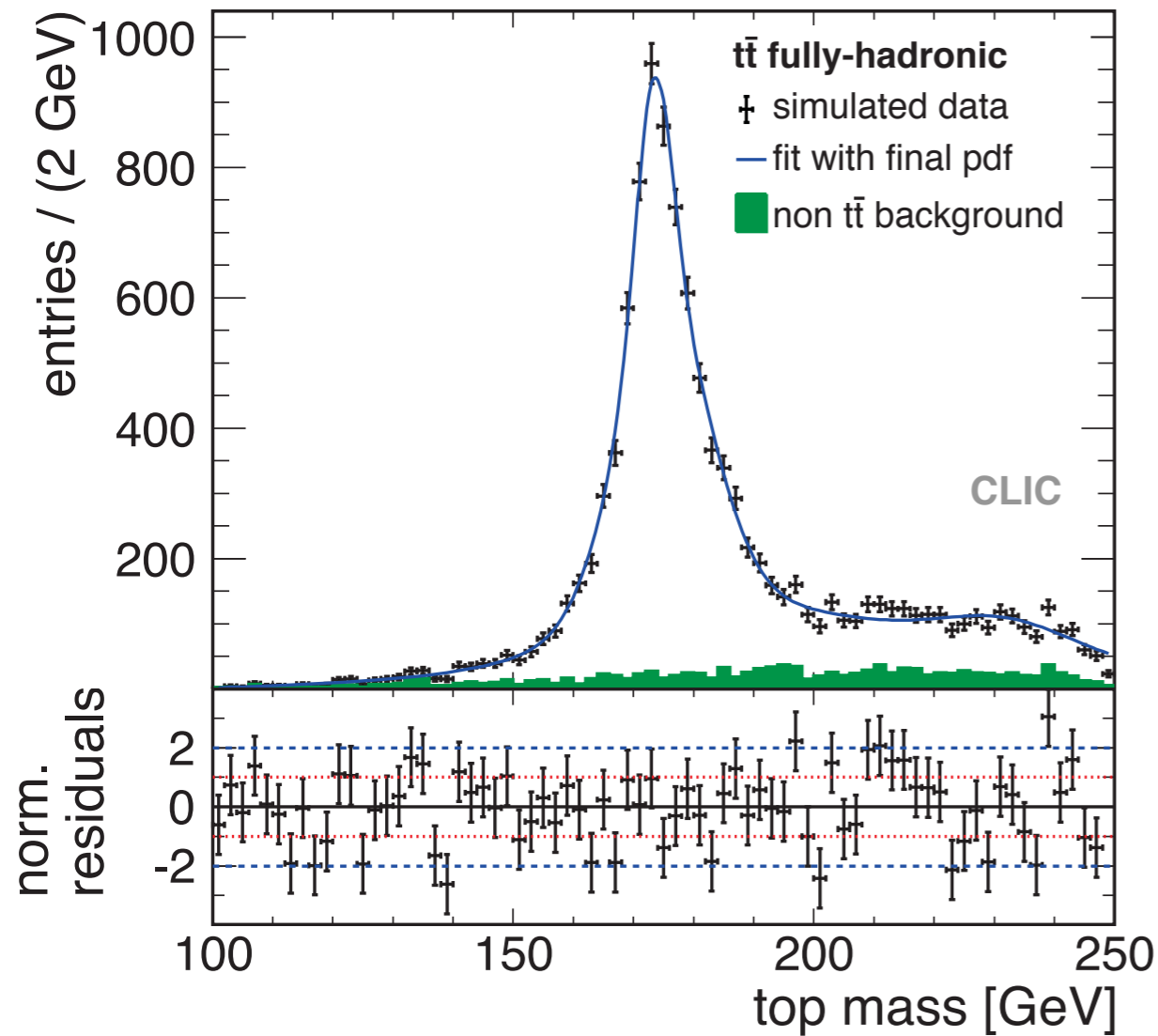


A_{FB} at the threshold



Can contribute to m_t and Γ_t measurement
☹️ Detector level analysis still missing (?)

Mass Reconstruction Above Threshold



- Width less constrained than mass: substantial detector effects (peak width ~ 5 GeV compared to 1.4 GeV top width)

channel	m_{top}	Δm_{top}	Γ_{top}	$\Delta\Gamma_{\text{top}}$
fully-hadronic	174.049	0.099	1.47	0.27
semi-leptonic	174.293	0.137	1.70	0.40
combined	174.133	0.080	1.55	0.22

The Valencia jet algorithm

A new clustering jet reconstruction algorithm that combines the good features of lepton collider algorithms, in particular the **Durham-like distance criterion**;

$$d_{ij} = \min(E_j^{2\beta}, E_i^{2\beta}) (1 - \cos \theta_{ij}) / R^2$$

with the **robustness against background** of the longitudinally invariant k_t algorithm

$$d_{iB} = p_T^{2\beta}$$

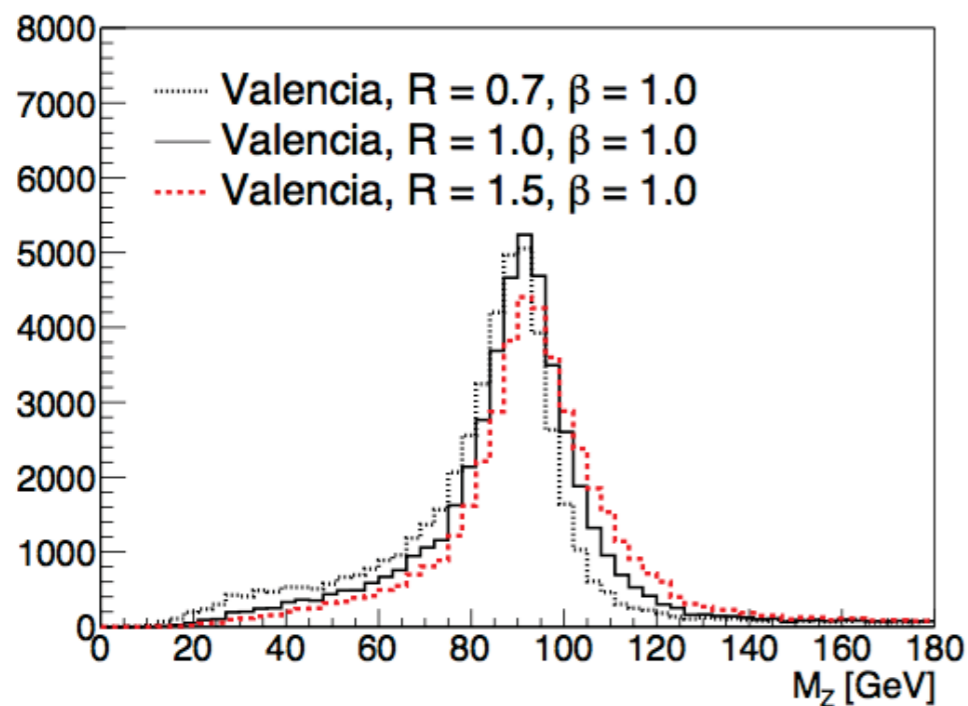
The exponent β allows to **tune the background rejection** level

The algorithm has been implemented as a **plugin for the *FastJet*** package and is available in `fjcontrib`

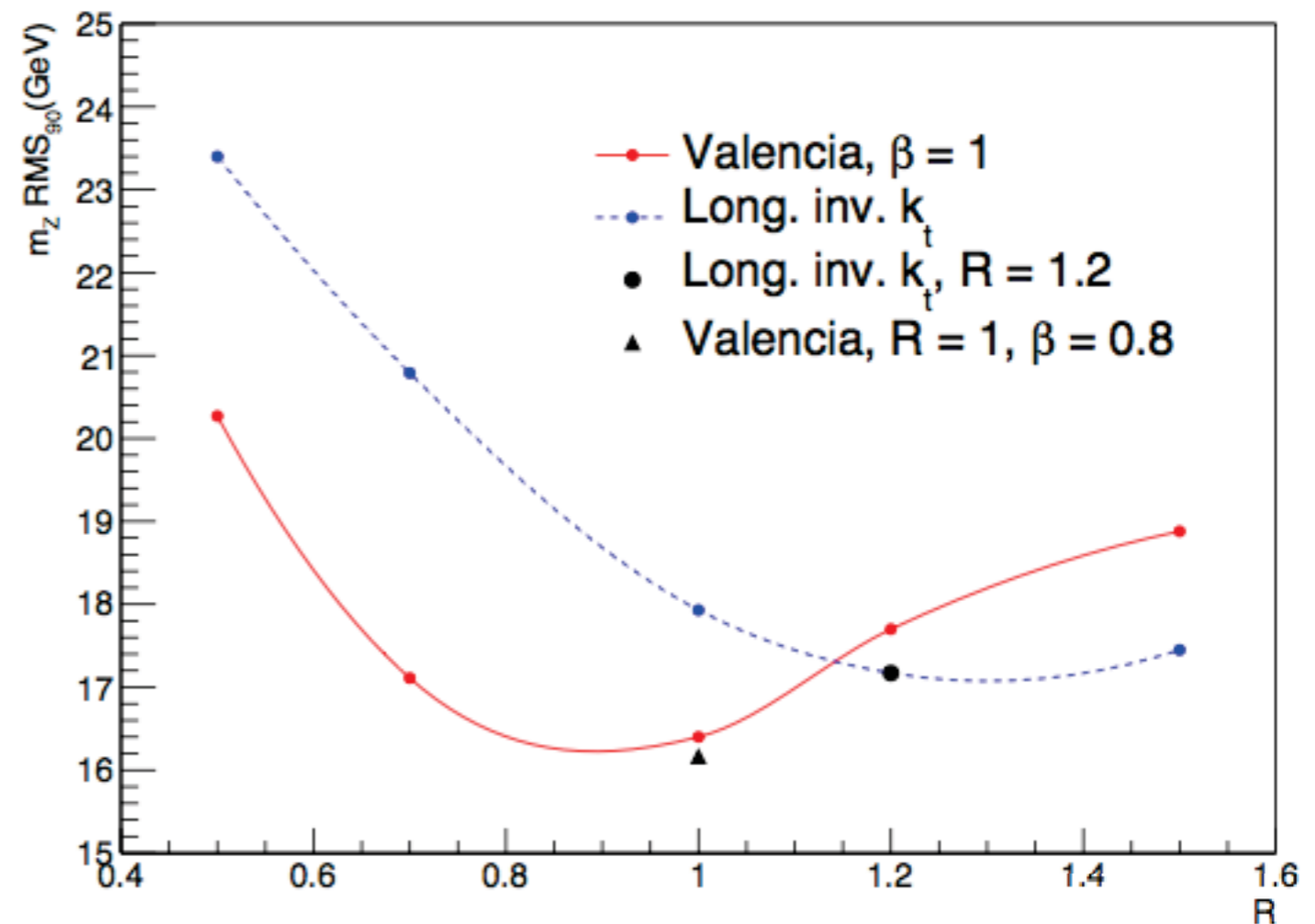
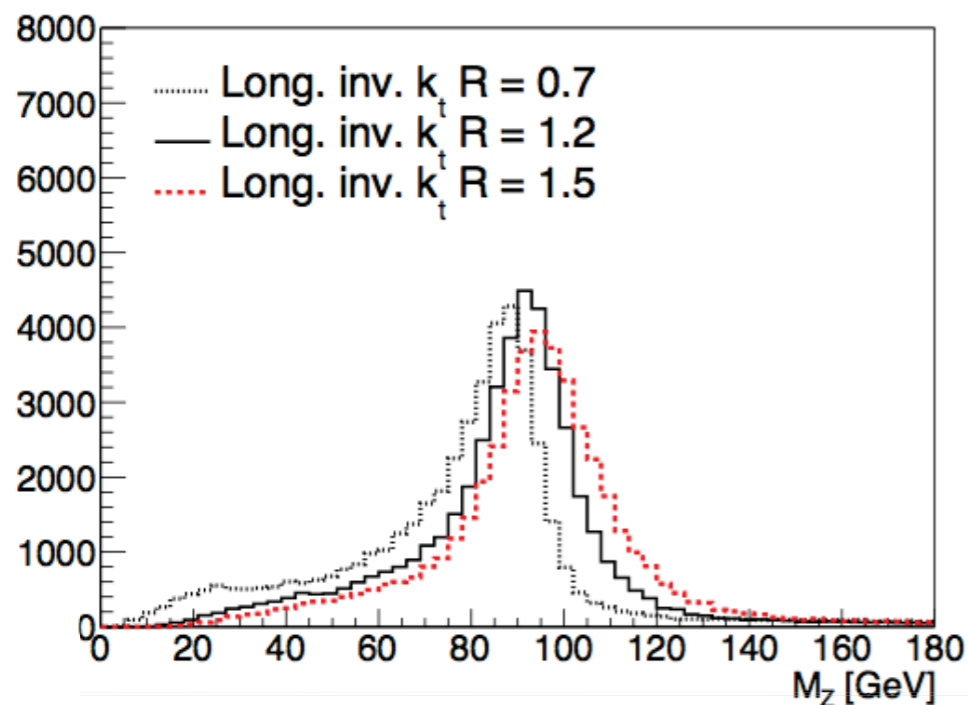
<https://fastjet.hepforge.org/trac/browser/contrib/contribs/ValenciaJetAlgorithm>



Parameter optimisation: R scan



The choice of parameters corresponds to the optimal setting determined in a scan over a broad range of parameters.



Jet reconstruction performance

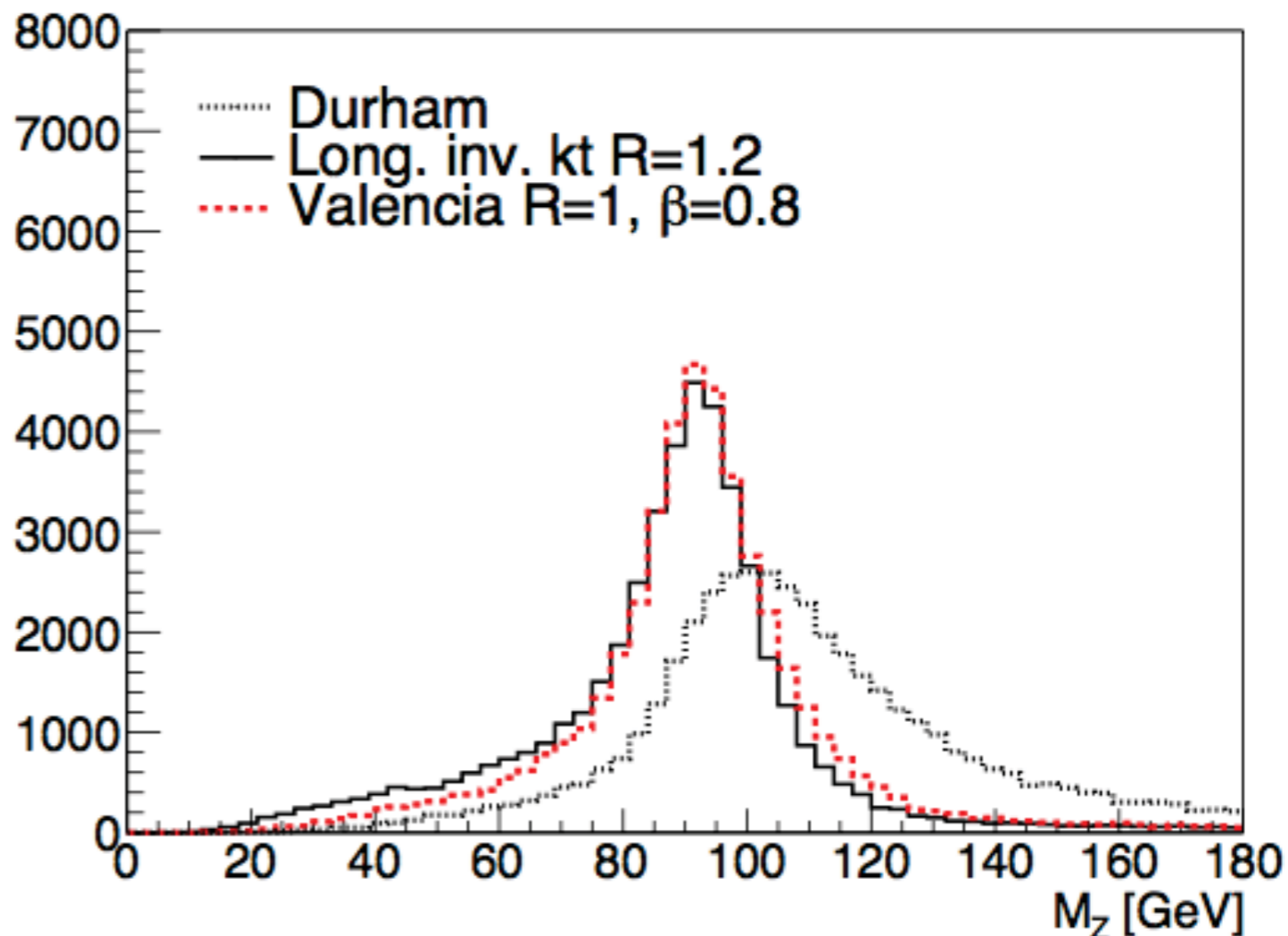
CLIC di-boson (ZZ)
production @ 500 GeV
+ 300 BX of $\gamma\gamma \rightarrow \text{hadrons}$

Reconstruct Particle Flow
objects using PANDORA
+ quality and timing cuts

Reconstruct jets
(exclusive, $n=4$)

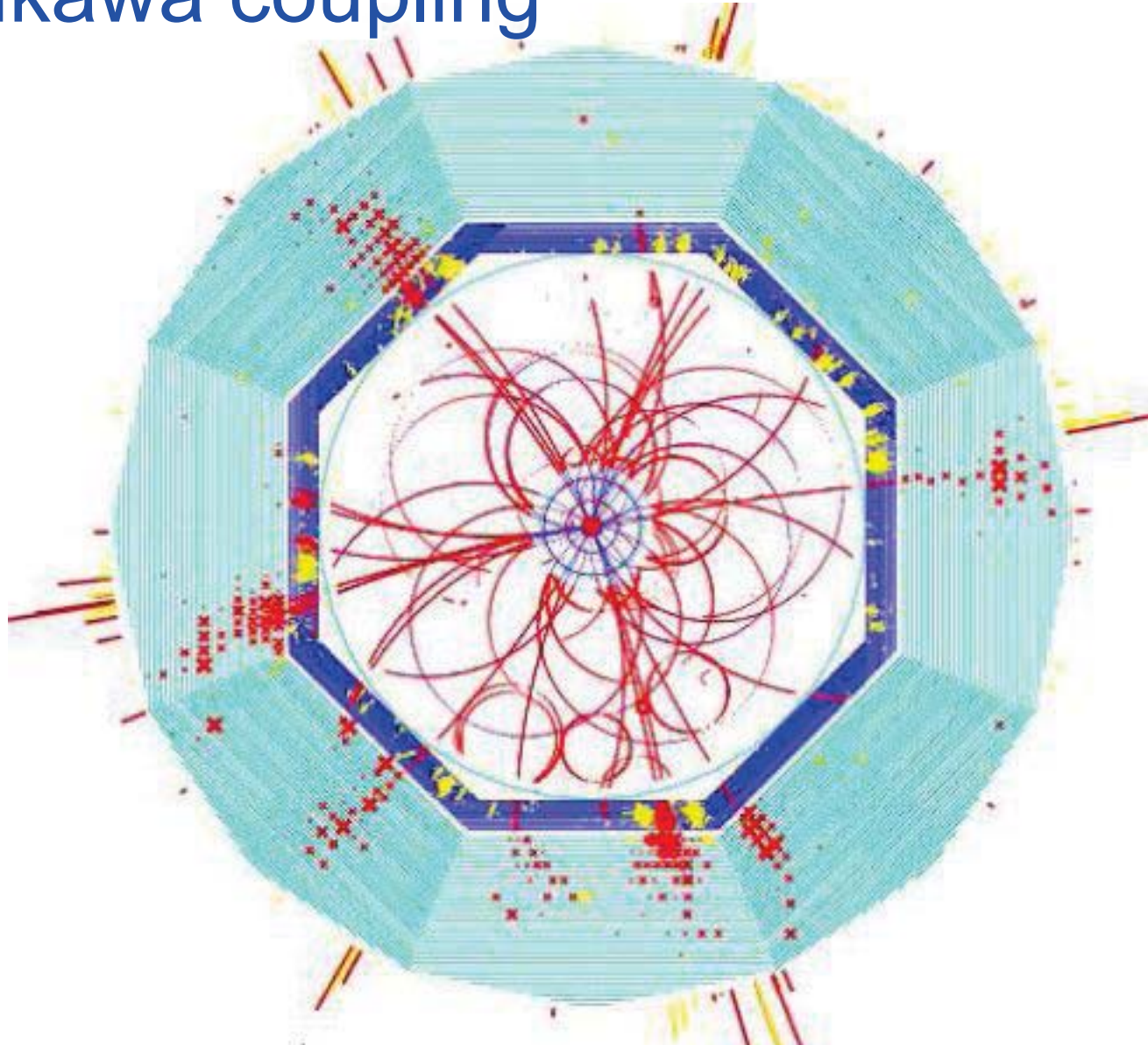
Form Z boson candidates,
selecting best jet pairs

**Nominal background: Durham is severely affected,
longitudinally invariant k_t and Valencia OK**



Top measurements

- Yukawa coupling

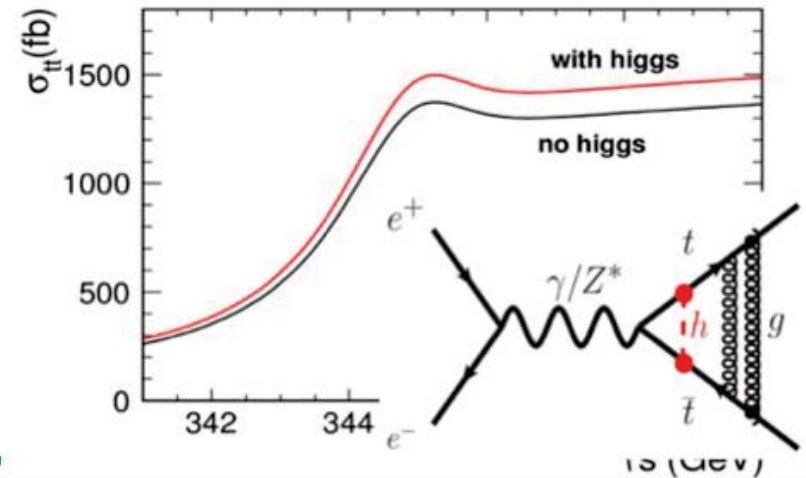


Top Yukawa coupling at threshold

The cross section is enhanced about **9%** by exchanging the Higgs boson !!

$$\sigma_{tt} \propto |\mathcal{M}_{w/o \text{ higgs}} + y_t^2 \mathcal{M}_{w/ \text{ higgs}}|^2$$

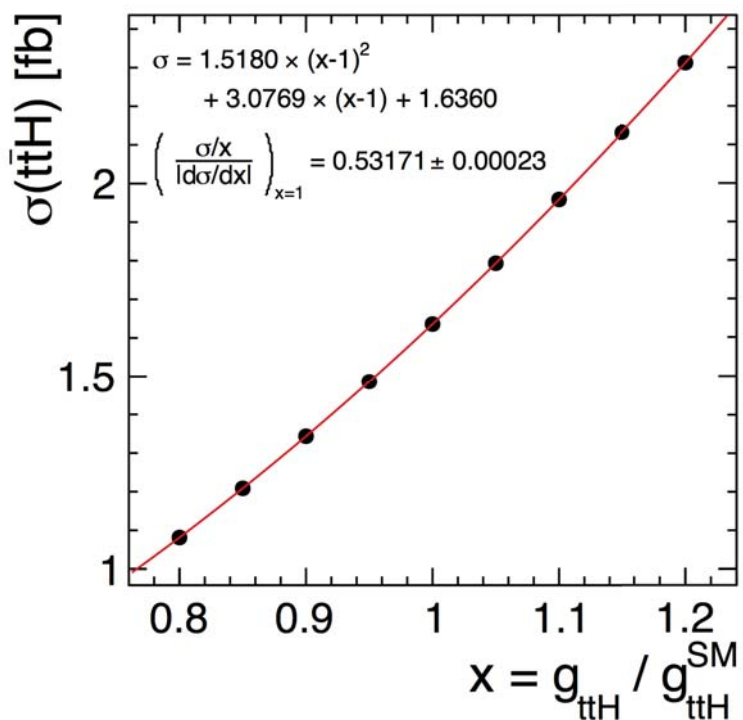
$$\frac{\delta y_t}{y_t} \sim \frac{109 \times \frac{1}{2} \times \frac{\delta \sigma}{\sigma}}{9}$$



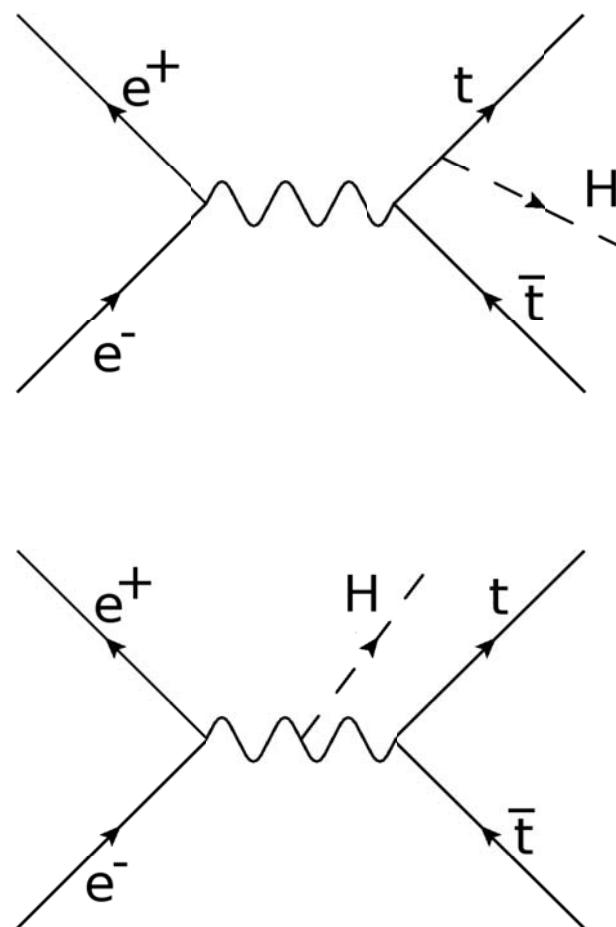
$$\int \mathcal{L} dt = \mathbf{100 \text{ fb}^{-1}}$$

	(2 + 1) param fit	3 param fit
mt	19 MeV	29 MeV
Γ_t	38 MeV	39 MeV
yt	4.6%	5.9%

Stat. Uncertainties
'add'
Theoretical
uncertainties $\sim 70 \text{ MeV}$



$$\frac{\Delta y_t}{y_t} = 0.53 \frac{\Delta \sigma}{\sigma}$$



(0.5 without contribution from Higgsstrahlung)

Expected # of events @ 500fb⁻¹

- $\sqrt{s} = 500 \text{ GeV}$, $M_h = 125 \text{ GeV}$, $(P_{e^-}, P_{e^+}) = (-0.8, +0.3)$
- production cross section
- Branching ratio

Process	σ (fb)
$e^-e^+ \rightarrow tth$	0.485
$e^-e^+ \rightarrow ttZ$	1.974
$e^-e^+ \rightarrow ttg(bb)$	1.058
$e^-e^+ \rightarrow tbW$	979.8

Decay mode	Branching ratio
$h \rightarrow bb$	0.577
$tt \rightarrow bq qbqq$	0.457
$tt \rightarrow blvbqq$	0.438
$tt \rightarrow blvblv$	0.105

- expected # of signals and Backgrounds(@500fb⁻¹)

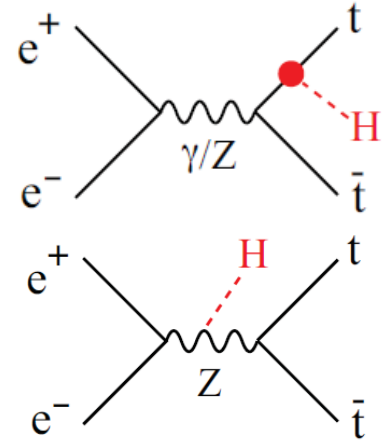
tth(tt6j, hbb)	63.9	tth(ttl n4j, hbb)	61.3
tth(ttall, hnobb)	102.6	ttZ	987
tth(ttlvlv2j, hbb)	14.6	ttg(bb)	529
		tbW	489902

tth \rightarrow 8jets(lv+6jets) analysis

- interference term is negligible
- counting analysis with cut based event selection

In this analysis, higgs decays into two b jets

- **4 b jets** out of 8(6) jets (b tagging: LCFIPlus)
- **No (one) isolated lepton**
- Use Kt clustering only for removing low Pt background

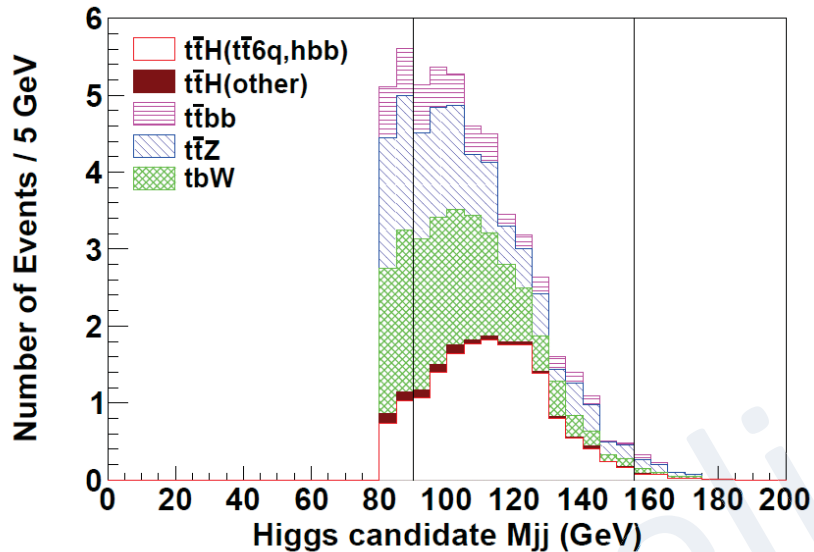


Event Selection

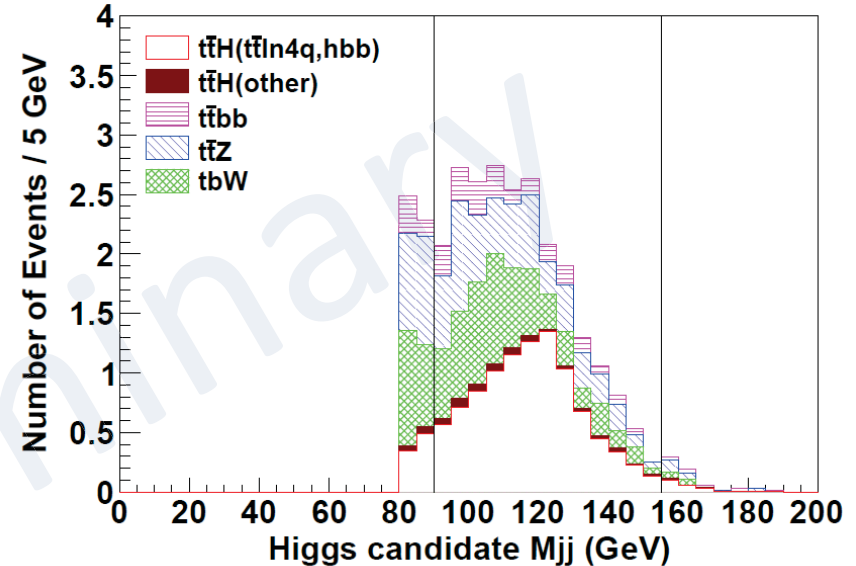
- **signal topology**
 - ✓ Y cut (6, 8 jet event)
 - ✓ No(one Isolated Lepton)
 - ✓ B jet candidate ≥ 4
- **detector acceptance**
 - $|\text{Jet } \cos\theta| \leq 0.99$
- **jet pairing**
 - ✓ $\chi^2 \leq 11.2$ (16.5)
- **kinematics**
 - ✓ Leading 2 Jet Energy Sum
 - ✓ Lowest 3 Jet Energy Sum (for 8jets mode)
(Lowest 2 Jet Energy Sum (for 6jets mode))
 - ✓ Missing momentum > 20 GeV (for 6jtes mode)
- **reconstructed mass**
 - ✓ top candidate $M_{jjj} \geq 140$ GeV
 - ✓ higgs candidate $M_{jj} \geq 80$ GeV
 - ✓ $90\text{GeV} \leq h$ candidate $M_{jj} \leq 155\text{GeV}$

Result of ILD Full Simulation

$t\bar{t}H \rightarrow 8\text{jets}$



$t\bar{t}H \rightarrow l\nu + 6\text{jets}$



- $\sqrt{s} = 500 \text{ GeV}, 500 \text{ fb}^{-1}$
- $N_{\text{sig}} = 14.7$
- $N_{\text{bkgd}} = 24.5$
- $N_{\text{sig}}/\sqrt{N_{\text{sig}} + N_{\text{bkgd}}} = \underline{2.351},$

- $\sqrt{s} = 500 \text{ GeV}, 500 \text{ fb}^{-1}$
- $N_{\text{sig}} = 9.77$
- $N_{\text{bkgd}} = 13.4$
- $N_{\text{sig}}/\sqrt{N_{\text{sig}} + N_{\text{bkgd}}} = \underline{2.029},$

- $W \rightarrow e, \mu, \tau + \nu$ inclusive analysis

Significance and Precision of top-Yukawa coupling measurement with Systematic Uncertainties

- $M_h=125$ GeV, $\sqrt{s} = 500$ GeV, 500 fb^{-1}
- systematics: b tag eff. $\pm 1,3\%$, JESF $\pm 1,3\%$
Br 1%, L 0.1%, pol 0.1%

tth \rightarrow 8 Jets

with systematics	significance	$ \Delta g_t/g_t $
0% (stat. only)	2.351	22.11%
1% (b, JESF)	2.343	22.19%
3% (b, JESF)	2.240	23.2%

tth \rightarrow lv+6jets

with systematics	significance	$ \Delta g_t/g_t $
0% (stat. only)	2.029	25.62%
1% (b, JESF)	2.019	25.75%
3% (b, JESF)	1.958	26.55%

Rough estimation of

significance and $|\Delta g_{tth}/g_{tth}|$

@ $\sqrt{s} = 500\text{-}550 \text{ GeV}$, 500 fb^{-1}

Combined result of
8jets and 6jets mode

(* syst. error is not included)

\sqrt{s} : $S/\sqrt{S+B}$: $|\Delta g_{tth}/g_{tth}|$ %

500 : 3.105 : 16.74

520 : 5.113 : 10.16

550 : 7.403 : 7.023

cross section (fb)

\sqrt{s} : tth(total) : ttz : ttbb : tbw

500 : 0.485 : 1.974 : 1.058 : 979.8

520 : 0.981 : 2.753 : 1.151 : 953.5

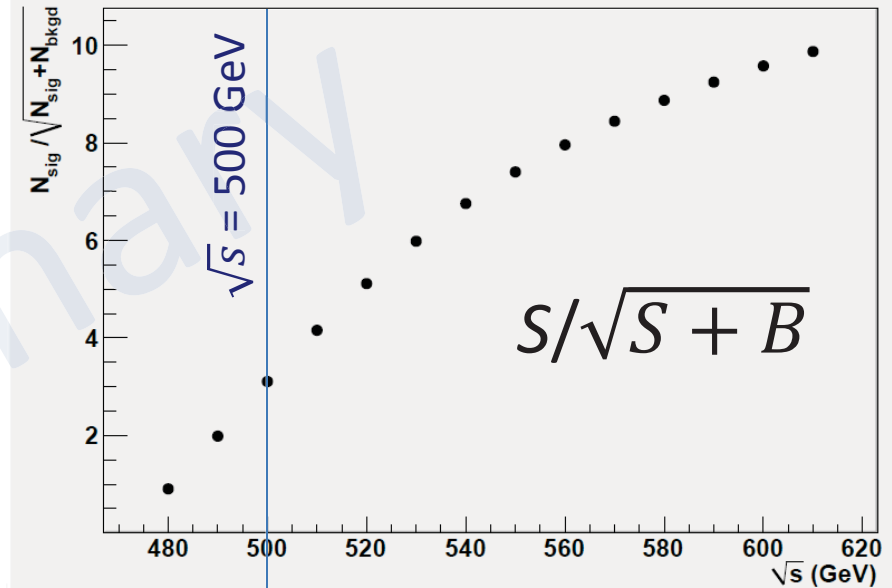
550 : 1.743 : 3.806 : 1.285 : 909.5

- ILC 1600fb^{-1} at $\sqrt{s} = 500 \text{ GeV}$

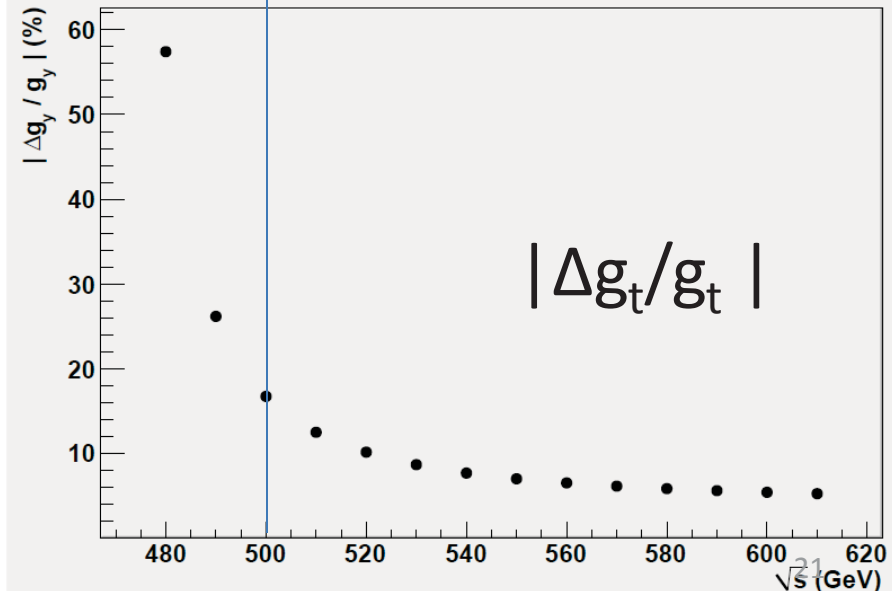
$|\Delta g_{tth}/g_{tth}| \sim 9.48\%$ (1% syst. included)

9.36% (stat. only)

Graph



Graph

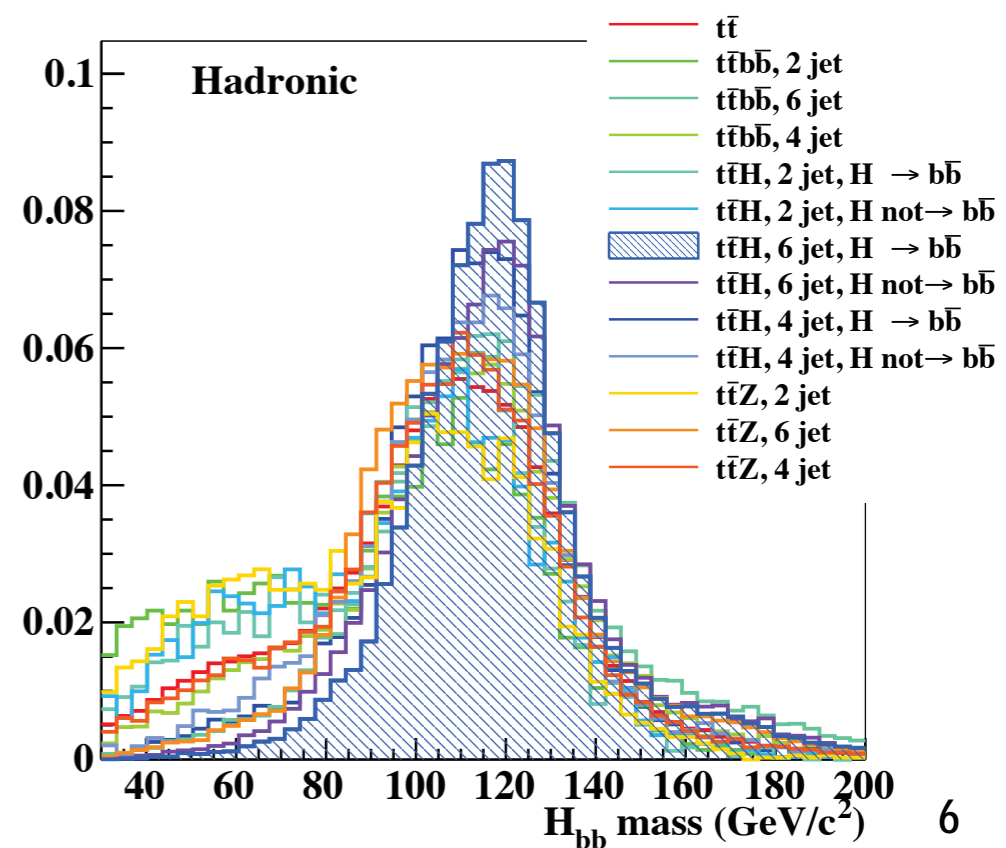
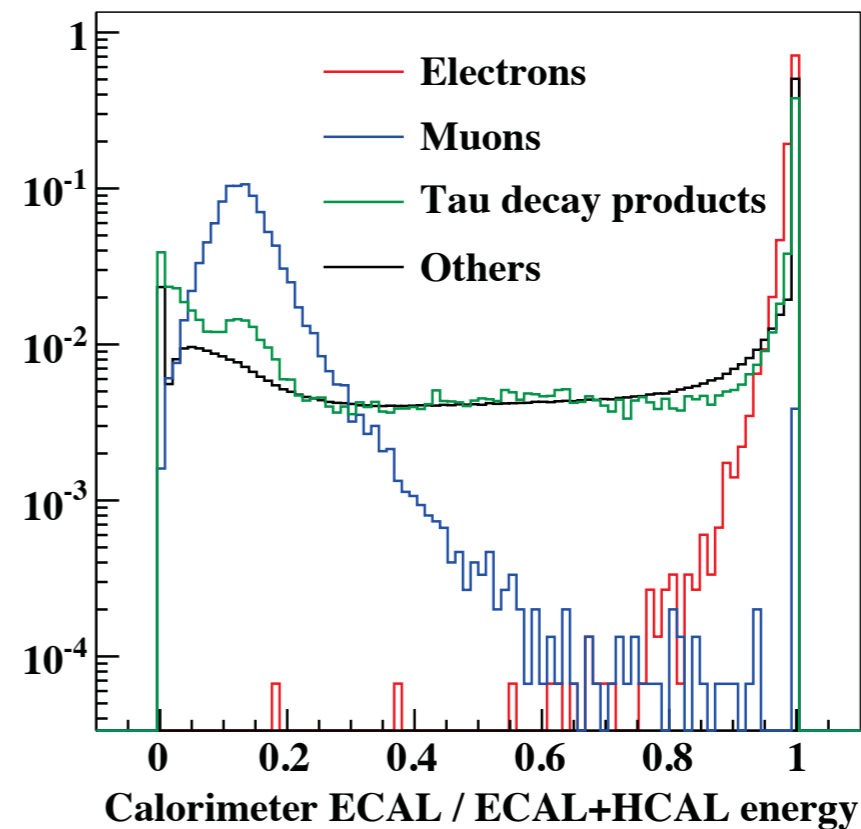


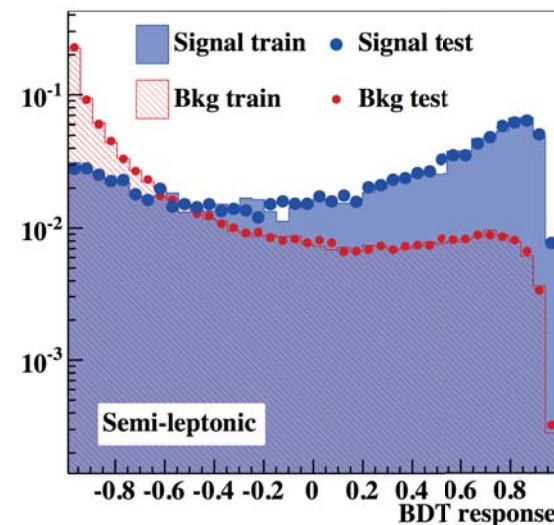
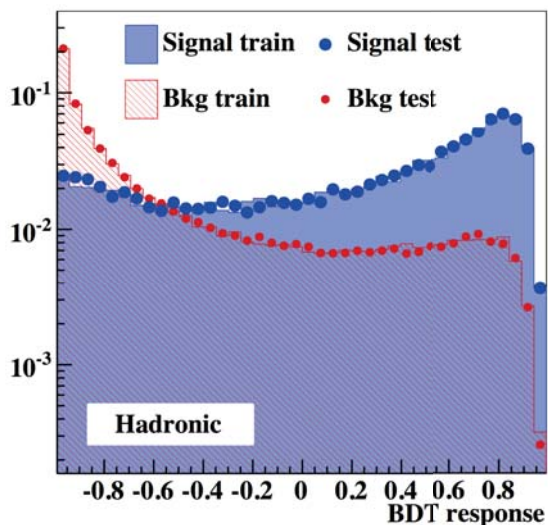
CLIC Analysis

- 1) Lepton finding
- 2) Jet clustering
- 3) Flavour tagging
 - 4 b jets!
- 4) Jet grouping
 - Choose permutation with smallest χ^2 :

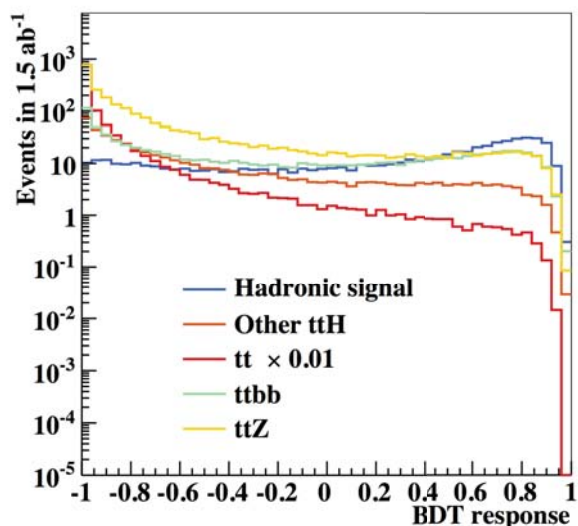
$$\chi^2 = \frac{(M_{12} - M_W)^2}{\sigma_W^2} + \frac{(M_{123} - M_t)^2}{\sigma_t^2} + \frac{(M_{45} - M_h)^2}{\sigma_h^2}$$

- 5) MVA selection on discriminating variables
- 6) Channel combination and compensation for Higgsstrahlung

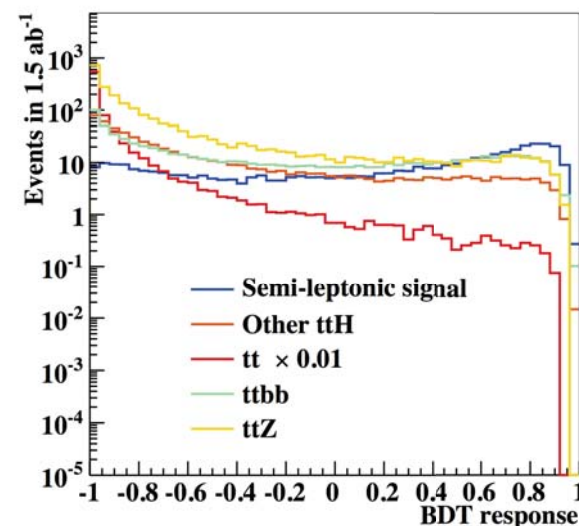




- Similar discrimination in both channels
- No sign of overtraining



- Cuts on BDT response providing maximum significance ($S / \sqrt{S + B}$) used for cross section extraction



Process	Evt in 1.5 ab^{-1}	Evt with 0 leptons	Evt pass Had BDT	Evt with 1 lepton	Evt pass SL BDT
$t\bar{t}H, 6 \text{ jet}, H \rightarrow b\bar{b}$	647	593	357 (60.2%)	49	9 (18.8%)
$t\bar{t}H, 4 \text{ jet}, H \rightarrow b\bar{b}$	623	178	62 (35.1%)	420	233 (55.3%)
$t\bar{t}H, 2 \text{ jet}, H \rightarrow b\bar{b}$	150	13	1 (10.7%)	61	20 (32.5%)
$t\bar{t}H, 6 \text{ jet}, H \not\rightarrow b\bar{b}$	473	306	38 (12.3%)	127	8 (6.52%)
$t\bar{t}H, 4 \text{ jet}, H \not\rightarrow b\bar{b}$	455	89	5 (5.81%)	246	19 (7.82%)
$t\bar{t}H, 2 \text{ jet}, H \not\rightarrow b\bar{b}$	110	6	0 (1.52%)	33	1 (3.66%)
$t\bar{t}b\bar{b}, 6 \text{ jet}$	824	737	287 (38.9%)	80	8 (9.75%)
$t\bar{t}b\bar{b}, 4 \text{ jet}$	794	222	44 (19.6%)	533	175 (32.9%)
$t\bar{t}b\bar{b}, 2 \text{ jet}$	191	16	1 (8.71%)	78	14 (18.1%)
$t\bar{t}Z, 6 \text{ jet}$	2,843	2,335	316 (13.5%)	322	12 (3.68%)
$t\bar{t}Z, 4 \text{ jet}$	2,738	711	49 (6.86%)	1,678	170 (10.2%)
$t\bar{t}Z, 2 \text{ jet}$	659	54	1 (2.03%)	248	13 (5.23%)
$t\bar{t}$	203,700	111,020	1,399 (1.26%)	77,110	523 (0.68%)

MVA selection efficiencies for the signal samples:

60% for the hadronic channel

55% for the semileptonic channel

Hadronic channel:

$$S / \sqrt{S + B} = 8.36$$

$$\Delta(\sigma(t\bar{t}H)) = 12.0\%$$

Combined:

$$\Delta(\sigma(t\bar{t}H)) = 8.1\%$$

$$\rightarrow \Delta(g_{t\bar{t}H}) = 4.3\%$$

Semileptonic channel:

$$S / \sqrt{S + B} = 9.17$$

$$\Delta(\sigma(t\bar{t}H)) = 10.9\%$$

$$L = 1.5 \text{ ab}^{-1}$$

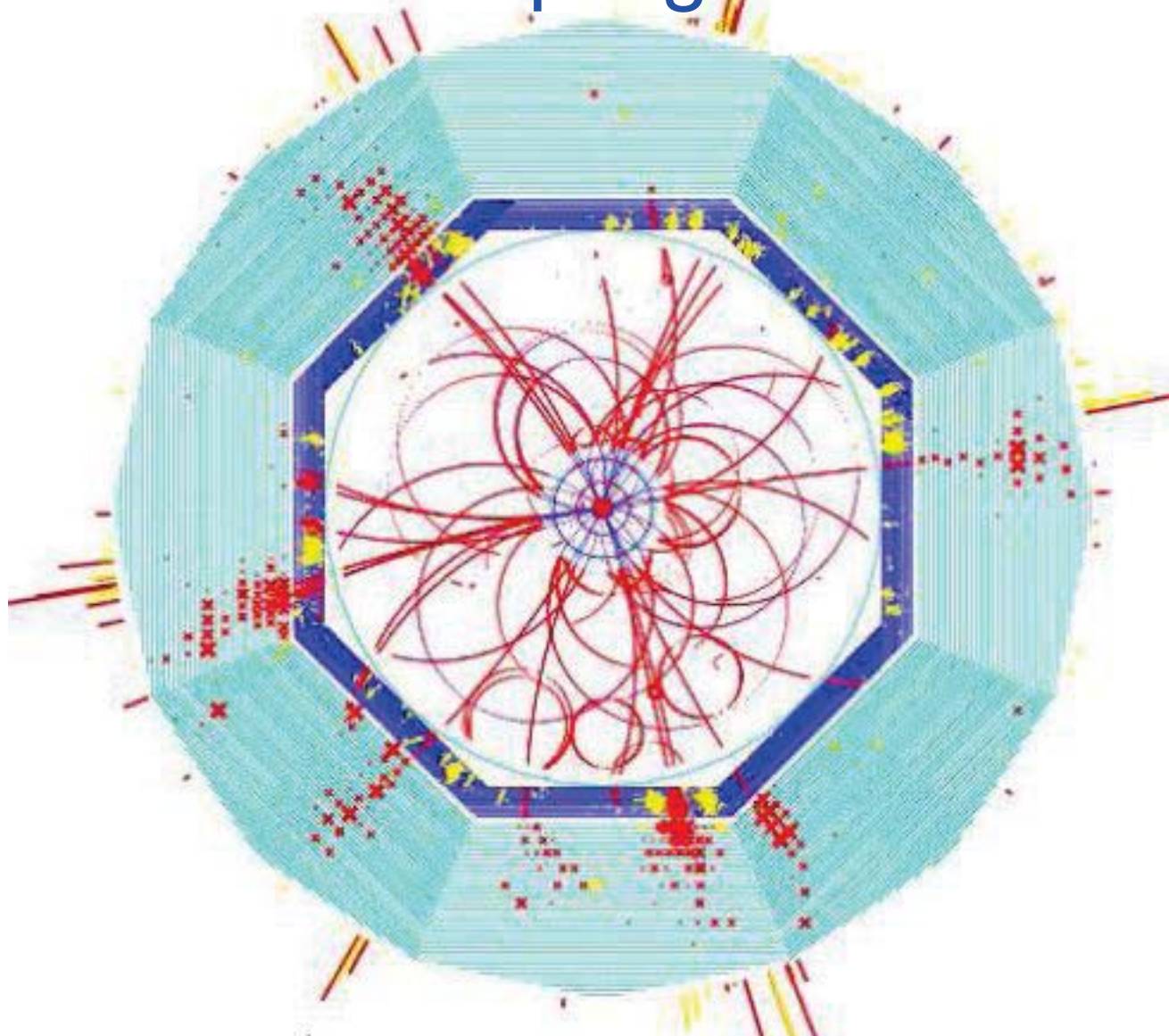
→ Precision on $g_{t\bar{t}H}$ would improve to better than 4% with -80% electron polarisation

For comparison:

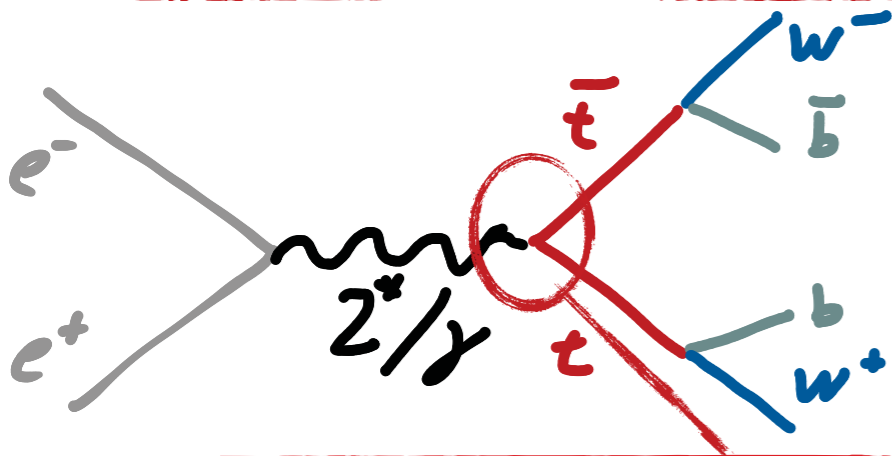
$\Delta(g_{t\bar{t}H}) = 4.3\text{-}4.5\%$ expected at 1 TeV ILC (from ILC TDR)

Top measurements

- Electroweak couplings



Electroweak Couplings of the Top Quark



- The production of top pairs provides direct access to electroweak couplings - axial and vector form factors

$$\Gamma_{\mu}^{ttX}(k^2, q, \bar{q}) = ie \left\{ \gamma_{\mu} \left(\tilde{F}_{1V}^X(k^2) + \gamma_5 \tilde{F}_{1A}^X(k^2) \right) + \frac{(q - \bar{q})_{\mu}}{2m_t} \left(\tilde{F}_{2V}^X(k^2) + \gamma_5 \tilde{F}_{2A}^X(k^2) \right) \right\}$$

X: Z, γ

A: axial coupling

V: vector coupling

- In total: 5 non-trivial CP-conserving form factors:

$$\begin{matrix} F_{1V}^{\gamma} & \boxed{F_{1A}^{\gamma}} & F_{2V}^{\gamma} \\ F_{1V}^Z & F_{1A}^Z & F_{2V}^Z \end{matrix} = 0 \text{ due to gauge invariance}$$

- Accessible through measurements of:

- **Total cross-section**

- **Forward-backward Asymmetry A_{FB}**

- **Helicity Angle λ** distribution (related to fraction of left- and right-handed tops)

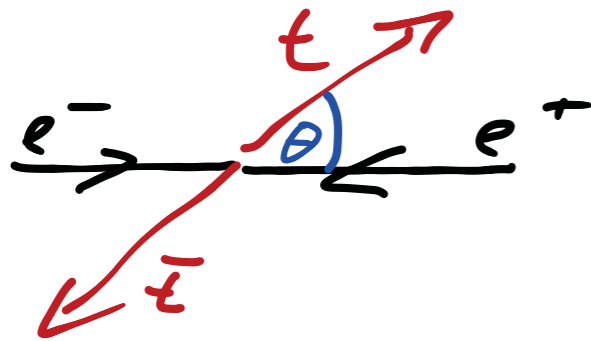
- For each: Two polarizations $e^-_L - e^+_R$, $e^-_R - e^+_L$

⇒ LC polarised beams crucial!

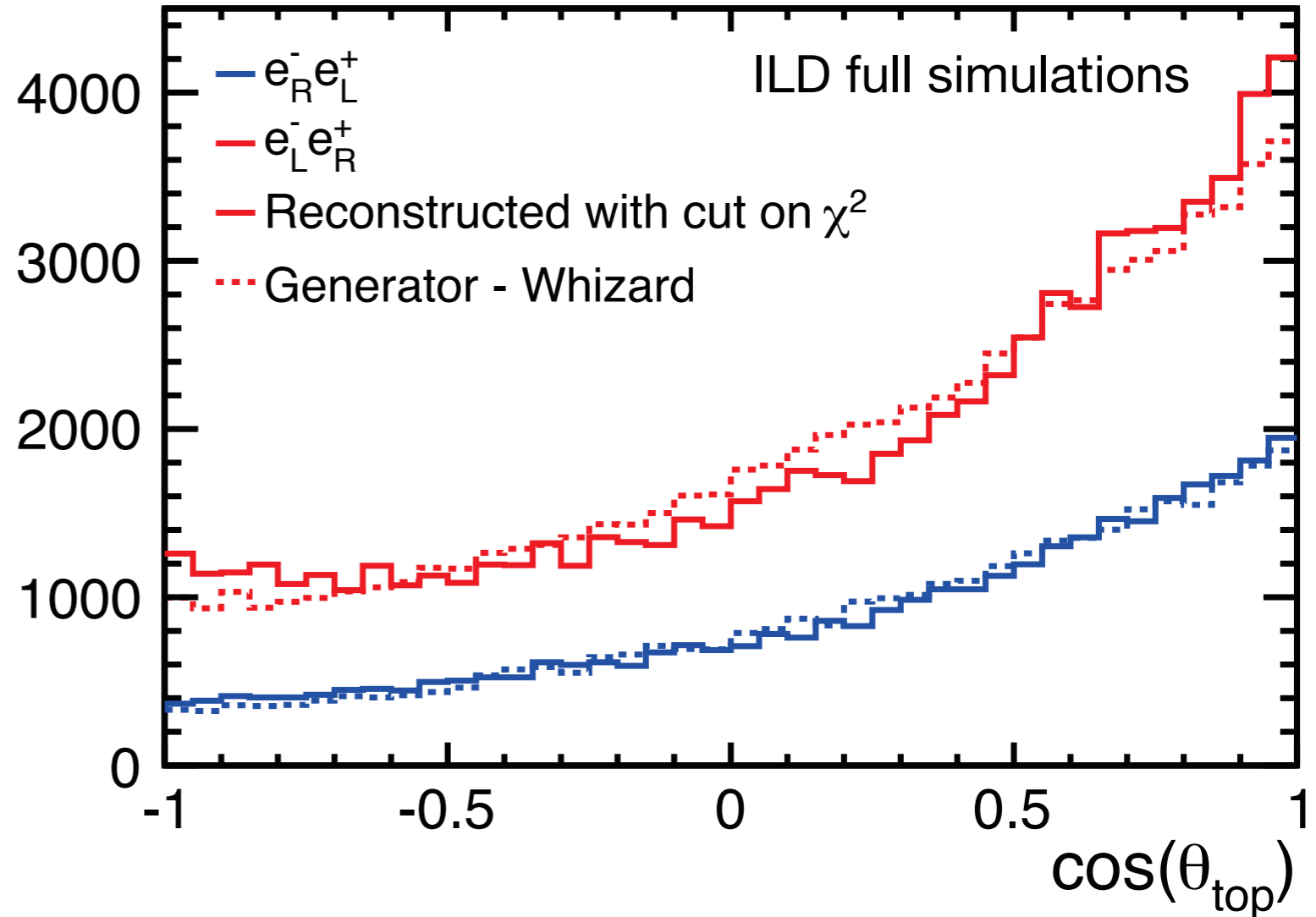
Accessing EW Couplings: Asymmetries & Angles

- Forward-backward asymmetry:

$$A_{FB}^t = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)}$$



- ILC, 500 GeV, 500 fb⁻¹
- Two polarisation configurations:
 - e⁻_Re⁺_L: P(e⁻) -80%, P(e⁺) +30%
 - e⁻_Le⁺_R: P(e⁻) +80%, P(e⁺) -30%

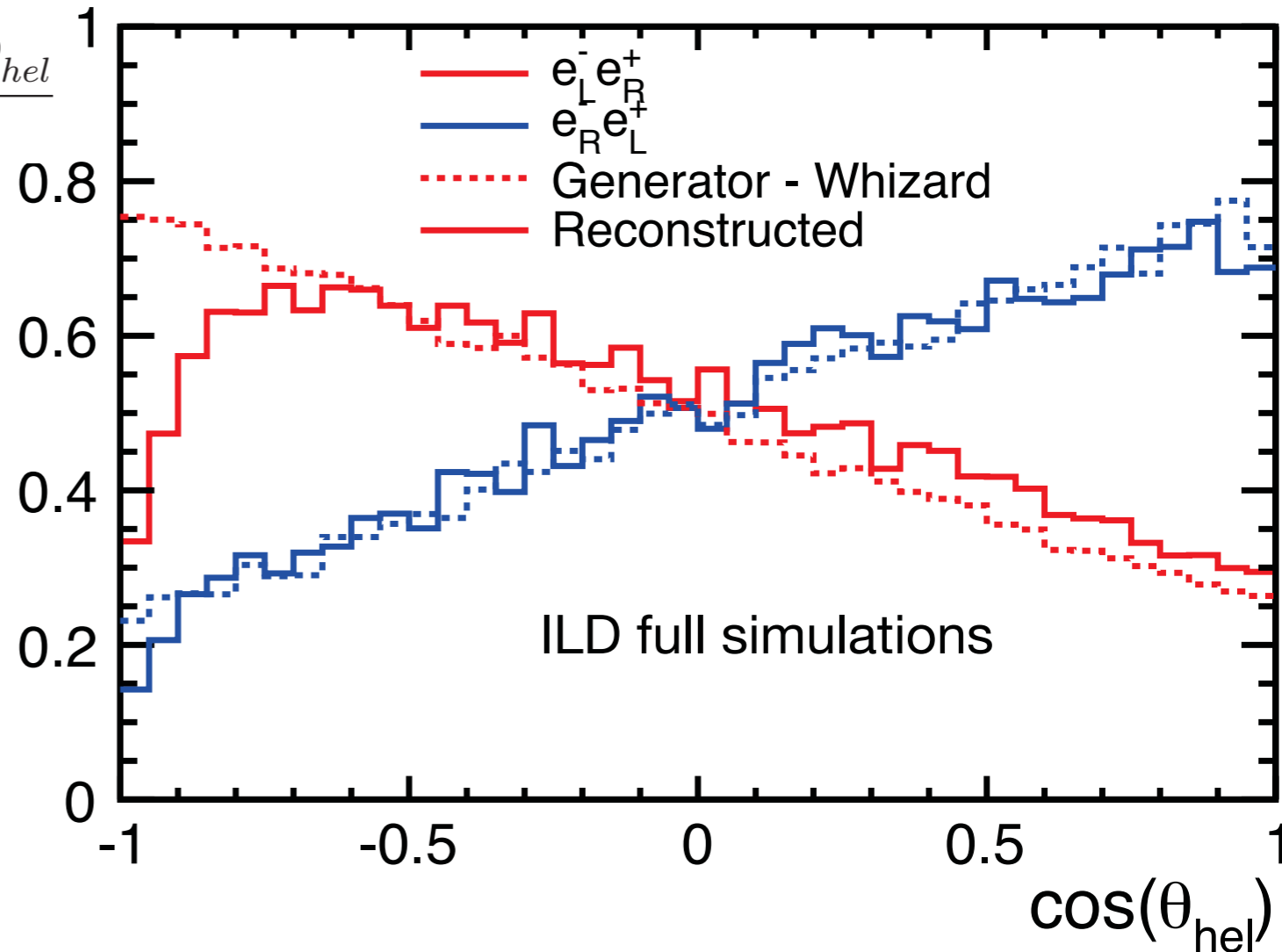
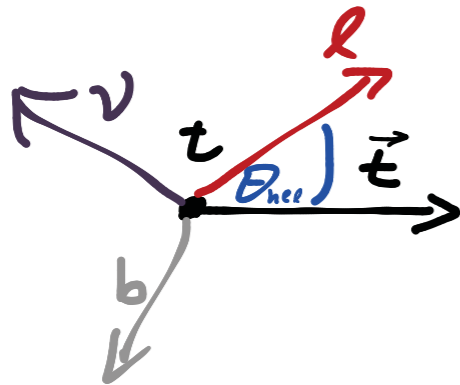


precision on asymmetry:
~2% (stat+ syst)

Accessing EW Couplings: Asymmetries & Angles

- Helicity Angle

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_{hel}} = \frac{1 + \lambda_t \cos\theta_{hel}}{2} = \frac{1}{2} + (2F_R - 1) \frac{\cos\theta_{hel}}{2}$$

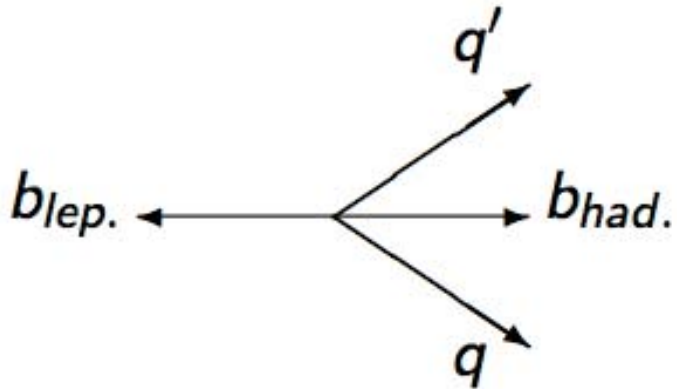


- ILC, 500 GeV, 500 fb⁻¹
- Two polarisation configurations:
 - e⁻_Re⁺_L: P(e⁻) -80%, P(e⁺) +30%
 - e⁻_Le⁺_R: P(e⁻) +80%, P(e⁺) -30%

precision on helicity angle:
~4% (stat+ syst)

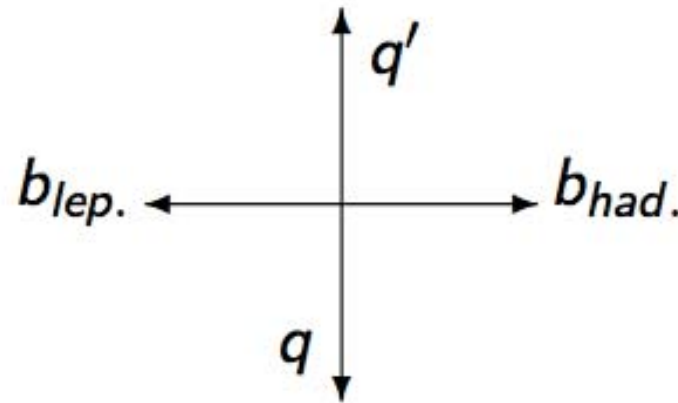
Experimental challenge b-charge reconstruction - Motivation

- To measure A_{FB} in fully hadronic decays there is no choice
- In semi-leptonic decays there is the charged lepton but



Right handed electron beam:

- mainly right handed tops
In final state (V-A)
- Hard W in flight direction of Top and soft b's
- Flight direction of t from flight direction of W



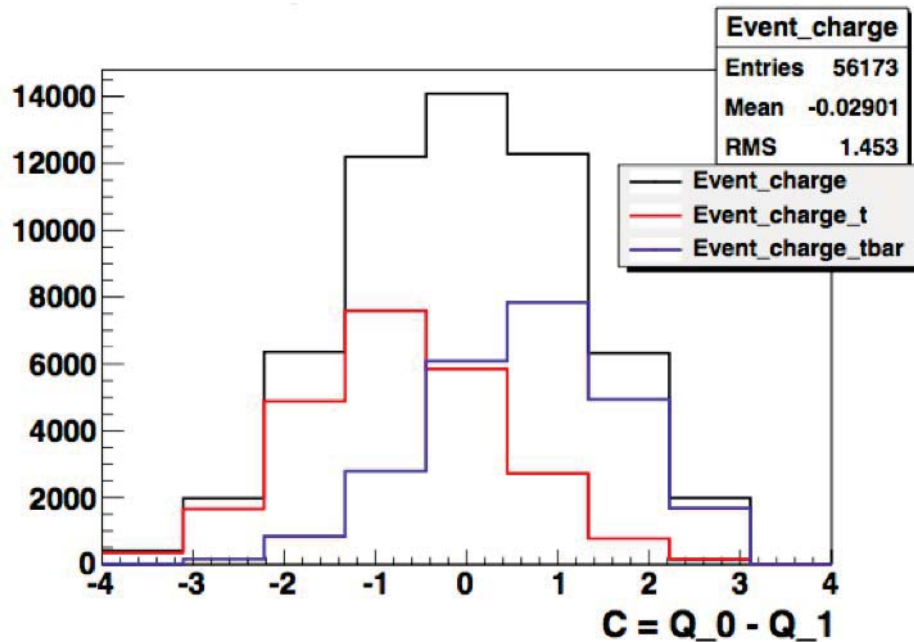
Left handed electron beam:

- mainly left handed tops
- Hard b in flight direction of Top and soft W's
- Flight direction of t from flight direction of b
=> Wrong association ↔ top flip

Measurement of b-charge to resolve ambiguities

Top polar angle using b charge

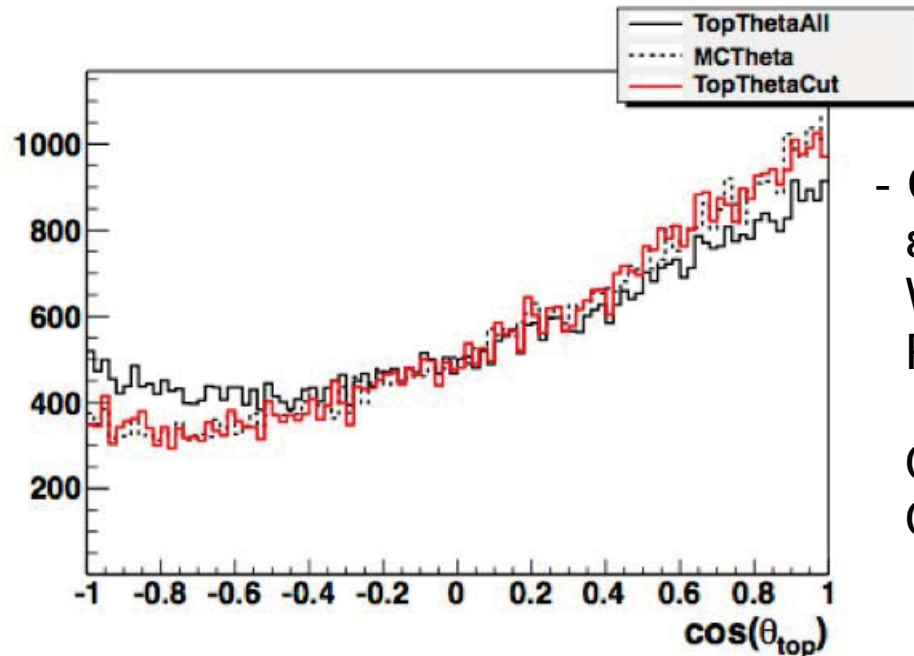
(SL Analysis)



Event charge $C = b_1 - b_2$

In SL can compare charge C with lepton charge to select clean sample

Use only events with correct C or C=0 (plus another cut on the Lorentz Factor)



- Clean reconstruction of top quark direction
 $\epsilon \sim 30\%$

Will improve with improving charge Reconstruction

Can already be considered as independent Cross check of existing results

Results of full simulation study for DBD at $\sqrt{s} = 500$ GeV

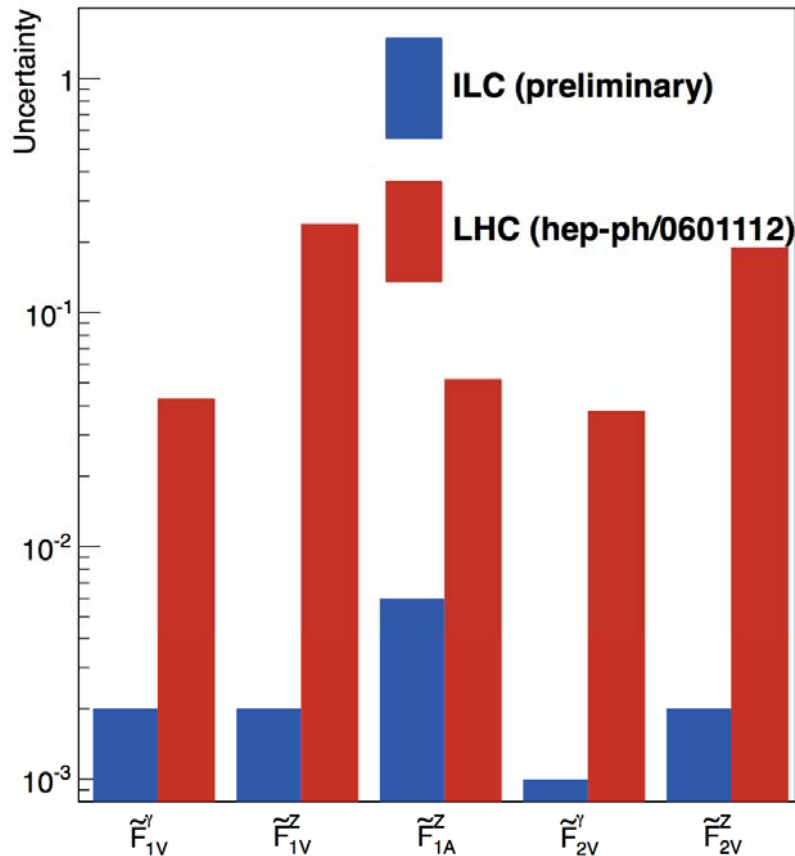
IFICLAL ArXiv: 1307.8102

Precision: cross section $\sim 0.5\%$,

Precision $A_{FB} \sim 2\%$,

Precision $\lambda_t \sim 3-4\%$

Accuracy on CP conserving couplings

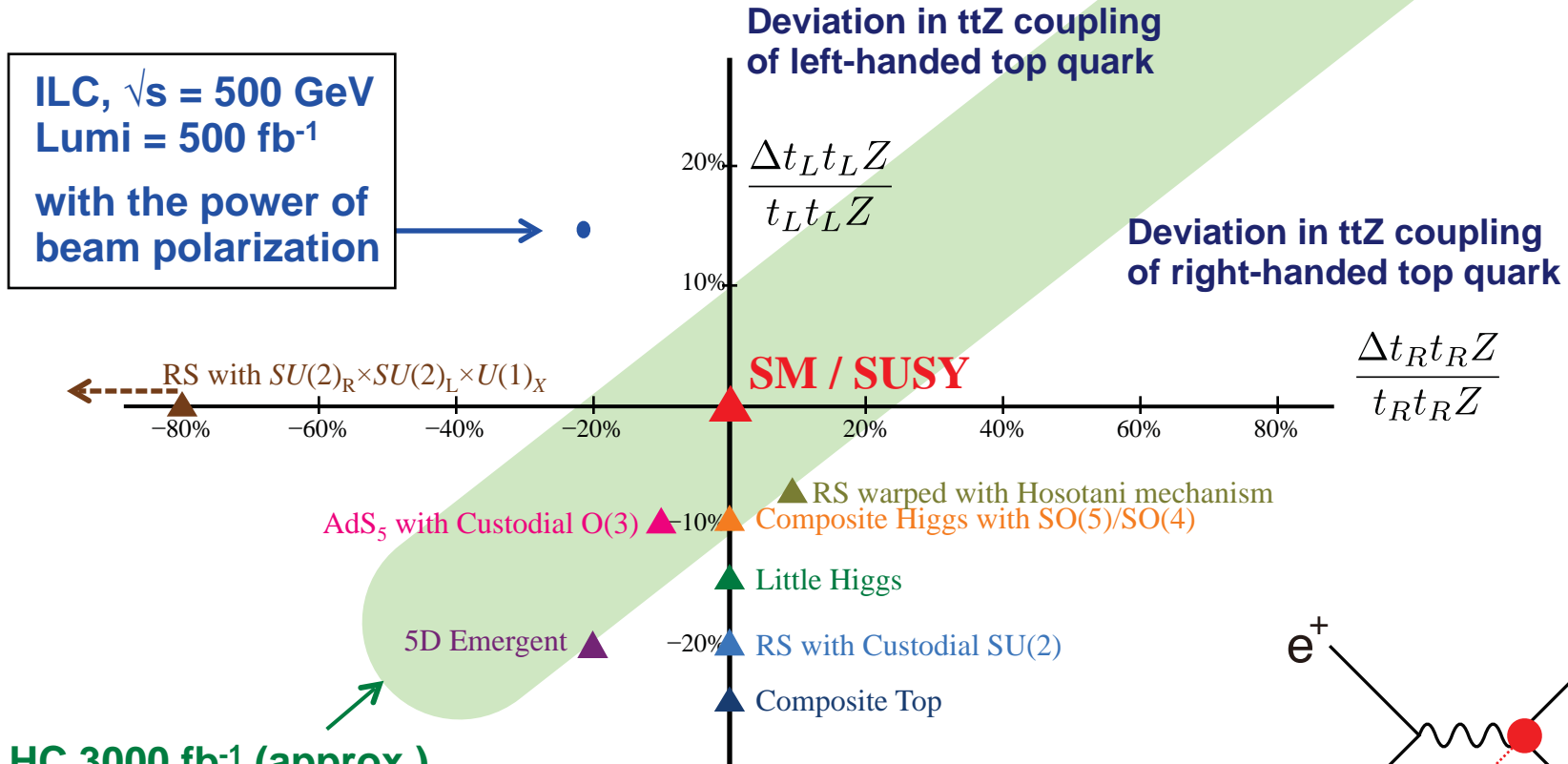


- ILC might be up to two orders of magnitude more precise than LHC ($\sqrt{s} = 14$ TeV, 300 fb^{-1})
Disentangling of vector/axial vector couplings for ILC
One variable at a time For LHC
However LHC projections from 8 years old study
- Need to control experimental (e.g. Top angle) and theoretical uncertainties (e.g. Electroweak corrections)
-> Dedicated work has started
- Potential for CP violating couplings at ILC under study
(However CP violation would rather show up at threshold)

ILC will be indeed high precision machine for electroweak top couplings

Impact of BSM on Top Sector

In composite Higgs models, it is often said that *the top quark is partially composite*, resulting in *form factors in ttZ couplings*, which can be measured at ILC. *Beam polarization is essential* to distinguish the *left- and right-handed couplings*.

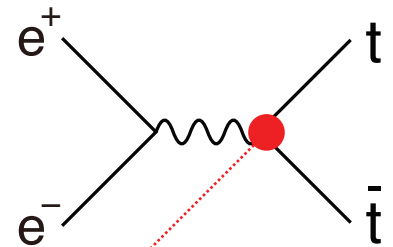


HL-LHC 3000 fb^{-1} (approx.)

Based on Baur, Juste, Orr, Rainwater, PRD71, 054013 (2005)

Deviations for different models for new physics scale at ~ 1 TeV.

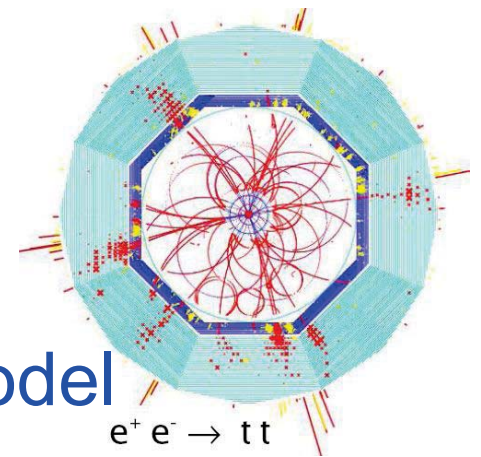
Based on F. Richard, arXiv:1403.2893



$$\Gamma_{\mu}^{ttZ}(k^2, q, \bar{q}) = ie \left\{ \gamma_{\mu} \left(\tilde{F}_{1V}^X(k^2) + \gamma_5 \tilde{F}_{1A}^X(k^2) \right) + \frac{(q - \bar{q})_{\mu}}{2m_t} \left(\tilde{F}_{2V}^X(k^2) + \gamma_5 \tilde{F}_{2A}^X(k^2) \right) \right\}$$

Conclusions

- ◆ Precise determination of top parameters is crucial for validation of the Standard Model (or any alternative BSM theory)
- ◆ Top threshold scan at the e^+e^- collider gives unique opportunities for precise mass, width and coupling determination
- ◆ Determination of Yukawa and electroweak couplings require running at higher beam energies
- ◆ Even in clean e^+e^- environment top event reconstruction is very challenging. Stringent requirements are imposed on detector performance.



The solid pillars of the LC physics program

Top quark



Discovered 1995 at Tevatron
LHC and ILC are/would be
Top factories

W Boson



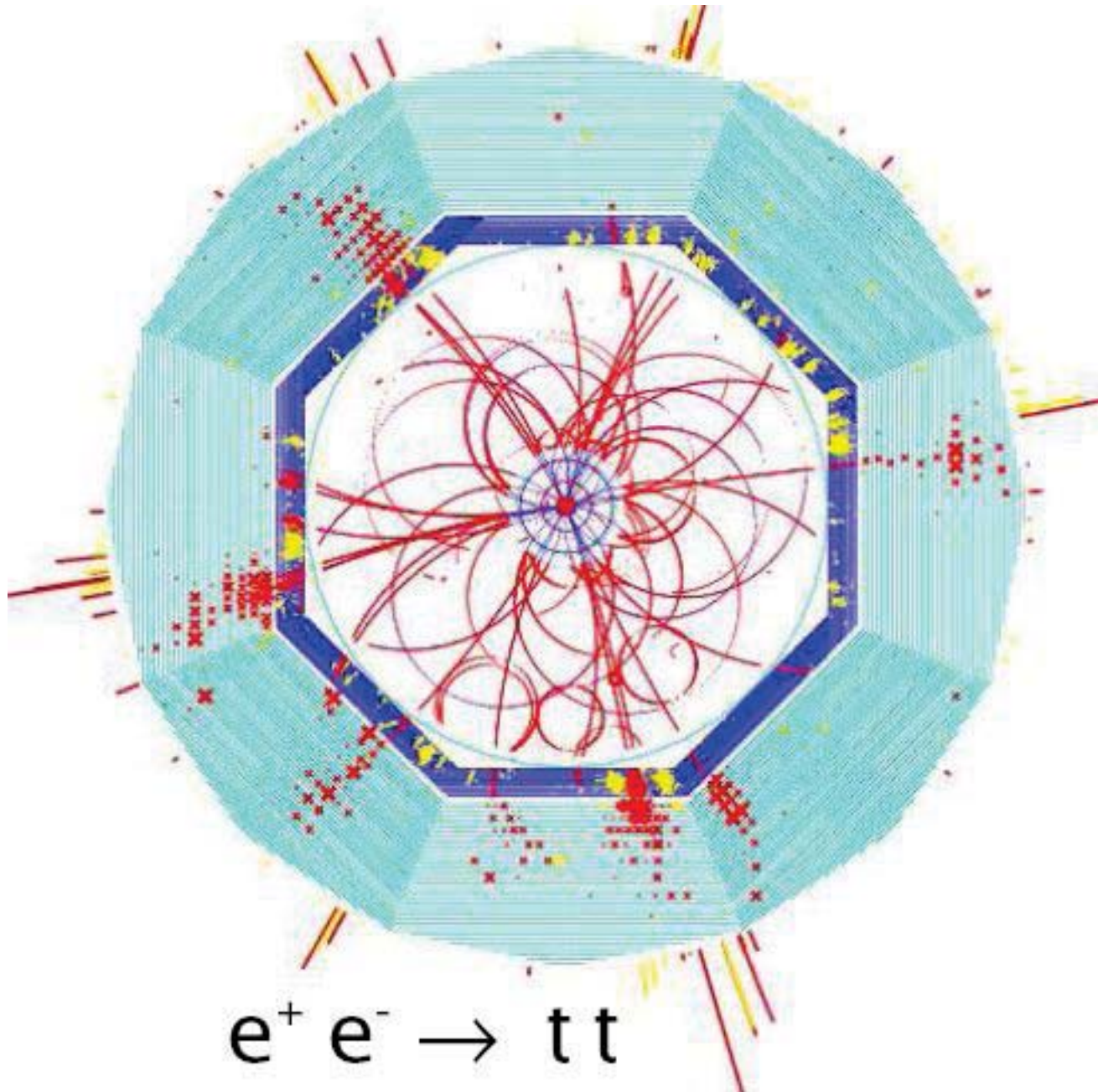
Discovered 1979 at SPS
Mass precisely at Tevatron
LHC and ILC are/would be
W factories

Higgs Boson



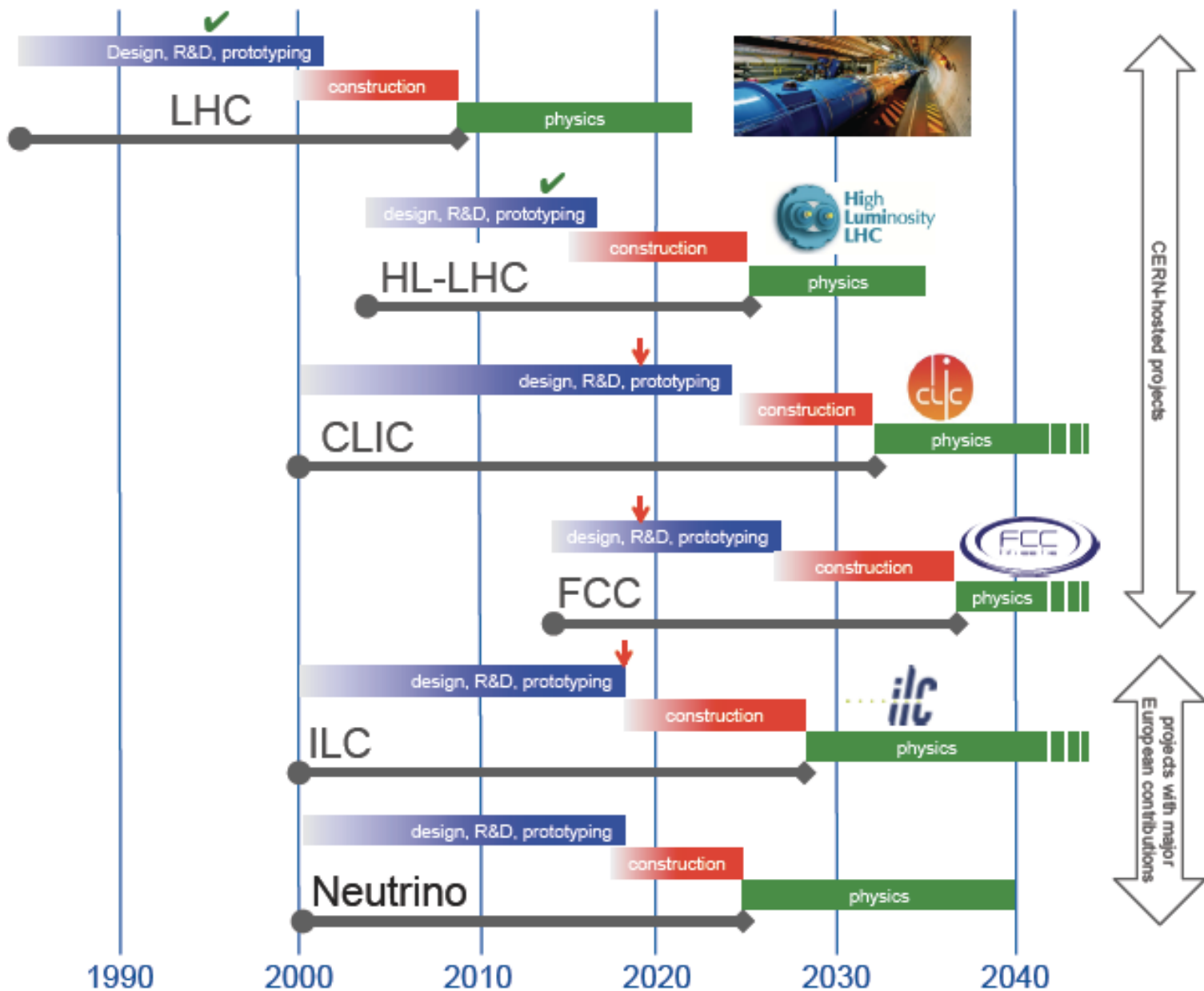
Discovered 2012 at LHC
ILC are/would be
Higgs factories
See talk by Klaus

Backup slides



$$e^+ e^- \rightarrow t t$$

Timeline: (HL-)LHC and future collider options



Jet reconstruction performance

The previous results in numbers: central value, width of the Z-boson mass peak and RMS_{90}

$\sqrt{s} = 500 \text{ GeV}$, no background overlay			
[GeV]	m_Z	σ_Z	RMS_{90}
Durham	90.6	5.4	13.8
long. inv. k_t	90.4	5.3	14.3
Valencia	90.3	5.2	12.5

$\sqrt{s} = 500 \text{ GeV}$, $0.3 \gamma\gamma \rightarrow \text{hadrons}$ events/BX			
[GeV]	m_Z	σ_Z	RMS_{90}
Durham	101.1	13.6	28.8
long. inv. k_t	92.0	9.0	17.2
Valencia	92.5	9.2	16.2

e^+e^- style algorithm can compete with hadron collider algorithm