

$B \rightarrow \tau \nu_\tau$ and a hint for H+

Flavor Physics at the LHC ERA
(workshop 2005-7, 5 meetings,
YellowReport in prep.)

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Flavor Physics

Rare processes...

- Here inclusive and exclusive B to tau decays – data, SM and 2HDM
- W^+ contra H^+ :
in inclusive processes H^+ exchange small correction, in exclusive – maybe 2 times bigger than W^+ !!! – in agreement with data
- Consider H^+ in 2HDM (II) – large tan beta

Jeffrey Berryhill (FNAL) on behalf of other WG2 editors and contributors

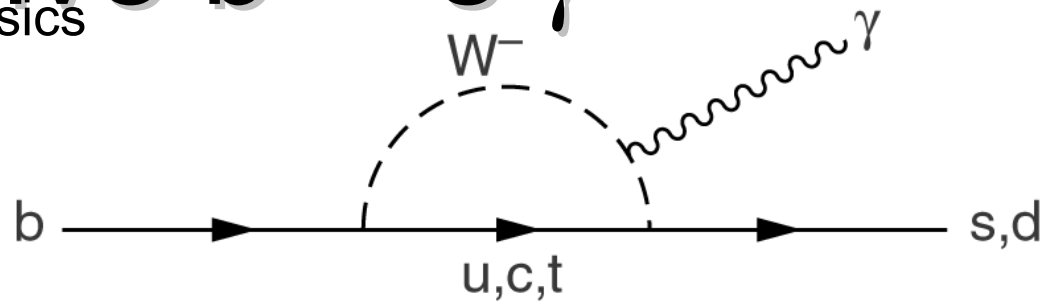
- Radiative penguin B decays ($b \rightarrow s, d \gamma$)
- Electroweak penguin B decays ($b \rightarrow s l^+ l^-$)
- Rare B decays with taus or neutrinos ($B \rightarrow \tau \nu, b \rightarrow s \nu \nu$)
- Very rare leptonic B decays ($B_s \rightarrow \mu \mu$)

Inclusive $b \rightarrow s \gamma$

A “Standard Candle” of flavor physics

Sensitive to top quark couplings

V_{td}, V_{ts}



Photon is DIS probe of B (shape function, m_B)

Broad sensitivity to new physics (2HDM, SUSY, LR, LED, little Higgs)

Misiak et al.

$$\mathcal{B}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}} = \mathcal{B}(\bar{B} \rightarrow X_c e \bar{\nu})_{\text{exp}} \left[\frac{\Gamma(b \rightarrow s \gamma)}{\Gamma(b \rightarrow c e \bar{\nu})} \right]_{\text{LO EW}} f \left(\frac{\alpha_s(M_W)}{\alpha_s(m_b)} \right) \times$$

$$\times \left\{ 1 + \underbrace{\mathcal{O}(\alpha_s)}_{\text{NLO } \sim 30\%} + \underbrace{\mathcal{O}(\alpha_s^2)}_{\text{NNLO } \sim 10\%} + \mathcal{O}(\alpha_{\text{em}}) \sim 4\% + \underbrace{\mathcal{O}\left(\frac{\Lambda^2}{m_b^2}\right)}_{\sim 1\%} + \mathcal{O}\left(\frac{\Lambda^2}{m_c^2}\right) \sim 3\% + \mathcal{O}\left(\frac{\Lambda}{m_b} \alpha_s\right) \sim 5\% \right\}$$

perturbative corrections
non-perturbative corrections
 (methods: Optical Theorem, Operator Product Expansion, Heavy Quark Effective Theory)

Need NNLO precision to compare with experiment

Recent estimate of NNLO decay rate!

Inclusive $b \rightarrow s \gamma$

Misiak et al.

In collaboration with: H.M. Asatrian, K. Bieri, M. Czakon, A. Czarnecki, T. Ewerth, A. Ferroglia, P. Gambino, M. Gorbahn, C. Greub, U. Haisch, A. Hovhannisyanyan, T. Hurth, A. Mitov, V. Poghosyan, M. Ślusarczyk and M. Steinhauser.

hep-ph/0609232

$$\mathcal{B}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}}^{\text{NNLO}} = (3.15 \pm 0.23) \times 10^{-4}$$

7.3% precision at NNLO

5% non-perturbative $\mathcal{O}\left(\alpha_s \frac{\Lambda}{m_b}\right)$ ← See e.g. hep-ph/0609224

3% parametric $(\alpha_s(M_Z), \mathcal{B}_{\text{semileptonic}}^{\text{exp}}, m_c, \dots)$ 3% m_c -interpolation ambiguity

2.0% 1.6% 1.1%

3% higher order $\mathcal{O}(\alpha_s^3)$

P.S. :Becher & Neubert, hep-ph/0610067: -5% shift from 2-loop corrections at intermediate and soft scales

$$\text{Br}(\bar{B} \rightarrow X_s \gamma) = (2.98 \pm 0.26) \cdot 10^{-4}$$

Inclusive $b \rightarrow s \gamma$

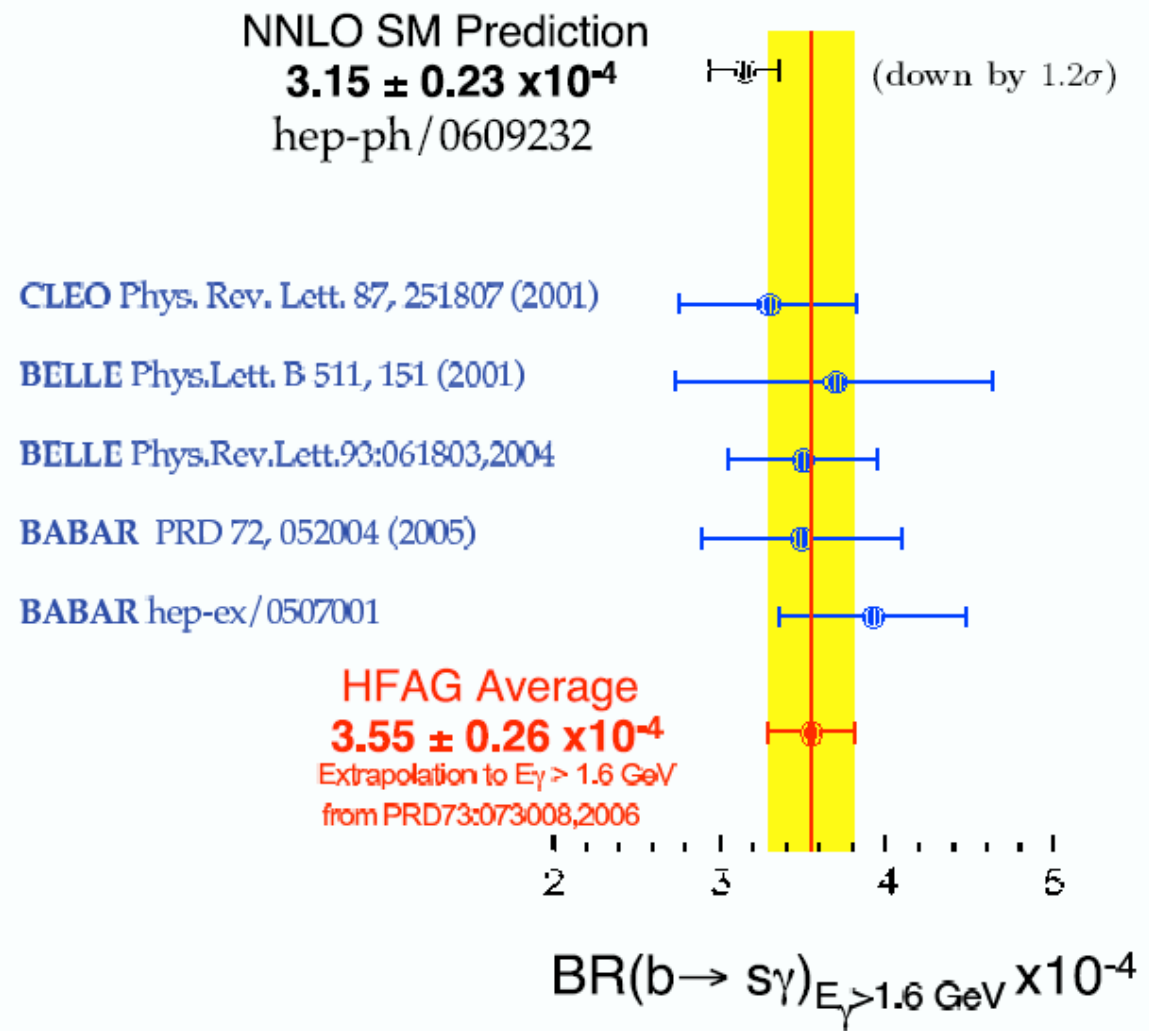
Experiment – NNLO Theory
= $+1.2\sigma$

B factories
are likely to improve
precision to 5%

Super B could push
down E_γ cutoff from current
1.8 GeV to 1.5 GeV

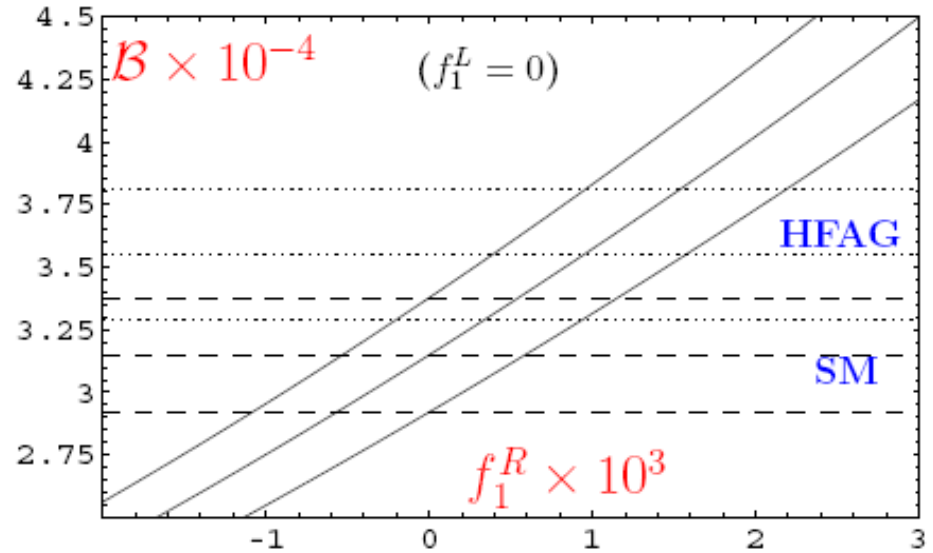
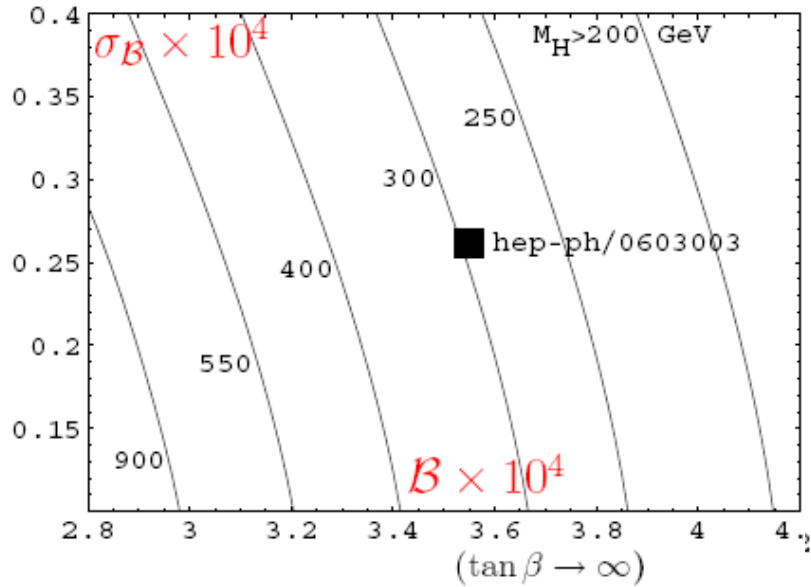
Super B incl. ACP precision:
0.9 % @ 5 ab^{-1}
0.3% @ 50 ab^{-1}

Super B can measure incl.
 $b \rightarrow d \gamma$ rate to
25% with 5 ab^{-1}



Inclusive $b \rightarrow s \gamma$

2HDM II Limit Right handed Wtb coupling



$f^R \sim 10^{-3}$

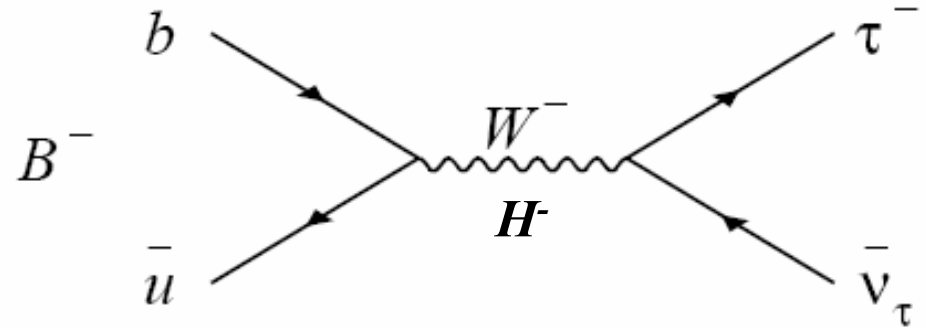
$M_{H^+} > 295 \text{ GeV @95\% C.L.}$

MSSM limits to be (re)evaluated: ideally NLO w/ minimal and general flavor violation

$B^+ \rightarrow \tau^+ \nu$

Simple decay through weak annihilation

Sensitive to B decay constant f_B or to charged Higgs boson



$$\mathcal{B}(B_u \rightarrow \tau \nu)^{\text{SM}} = \frac{G_F^2 m_B m_\tau^2}{8\pi} \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B = (1.59 \pm 0.40) \times 10^{-4}$$

$\tan^4 \beta$ modifications in 2HDM II model:

$$R_{B\tau\nu} = \frac{\mathcal{B}(B_u \rightarrow \tau \nu)}{\mathcal{B}(B_u \rightarrow \tau \nu)^{\text{SM}}} = r_H = \left[1 - \tan^2 \beta \frac{m_B^2}{m_{H^\pm}^2}\right]^2$$

f_B dependence can be removed via ratio with Δm_d , error shrinks 25% \rightarrow <13% (Isidori & Paradisi)

$$\left. \frac{\mathcal{B}(B_u \rightarrow \tau \nu)}{\tau_B \Delta M_{B_d}} \right|_{\text{SM}} = 1.77 \times 10^{-4} \left(\frac{|V_{ub}/V_{td}|}{0.464} \right)^2 \left(\frac{0.836}{\hat{B}_{B_d}} \right)$$

$B^+ \rightarrow \tau^+ \nu$ Belle PRL 97 (2006) 251802

Tag side reco of:

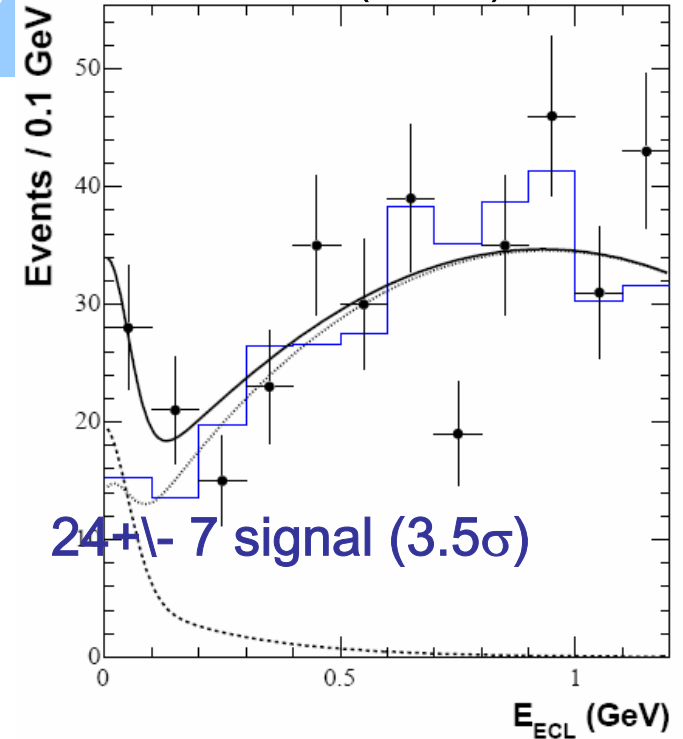
hadronic B decay (Belle, $\epsilon = 0.15\%$)
 or $D^0 / X \nu$ decay (BaBar, $\epsilon = 0.6\%$)

& signal side τ

(Belle: leptonic or 1- or 3-prong, $\epsilon = 16\%$)

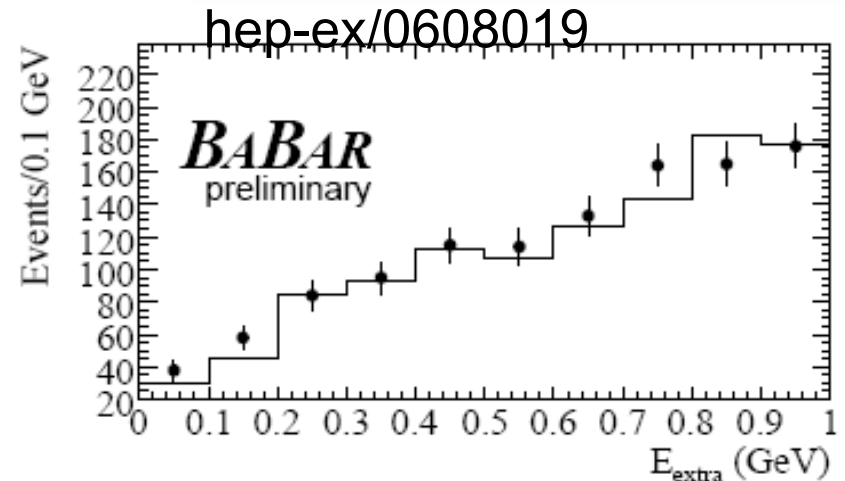
(BaBar: leptonic or 1-prong, $\epsilon = 13\%$)

& no other tracks & small extra ECAL energy



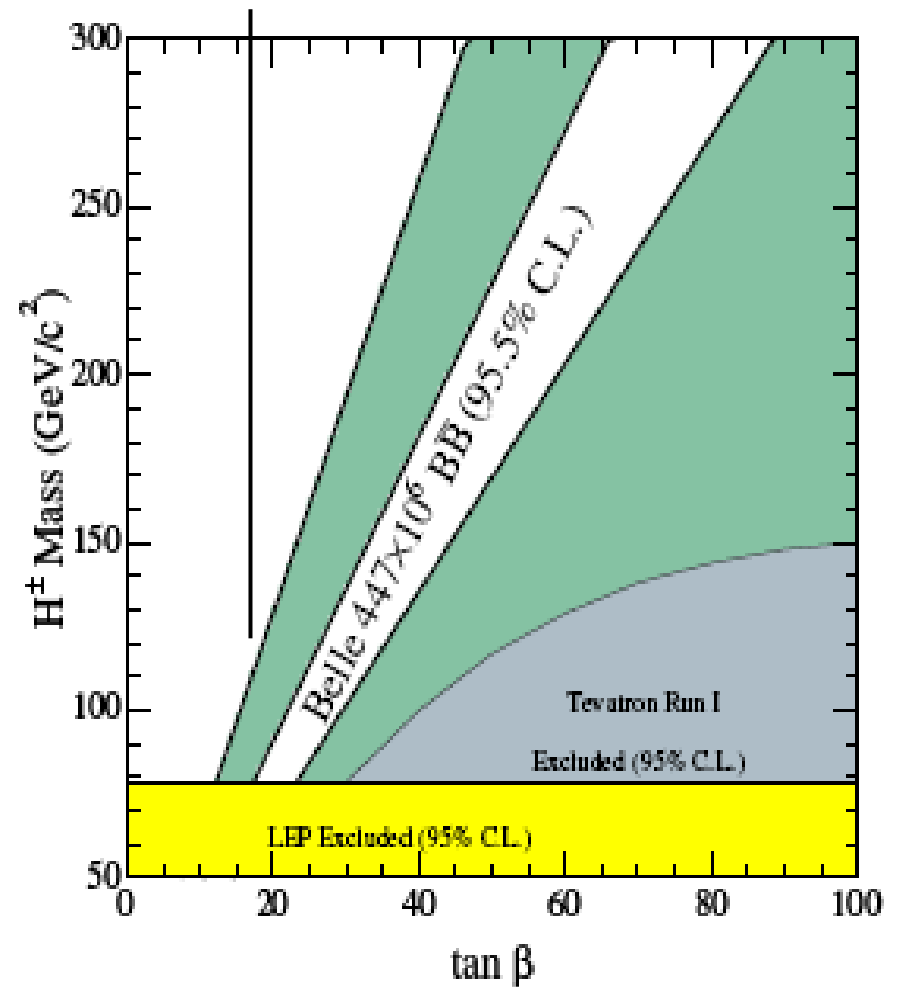
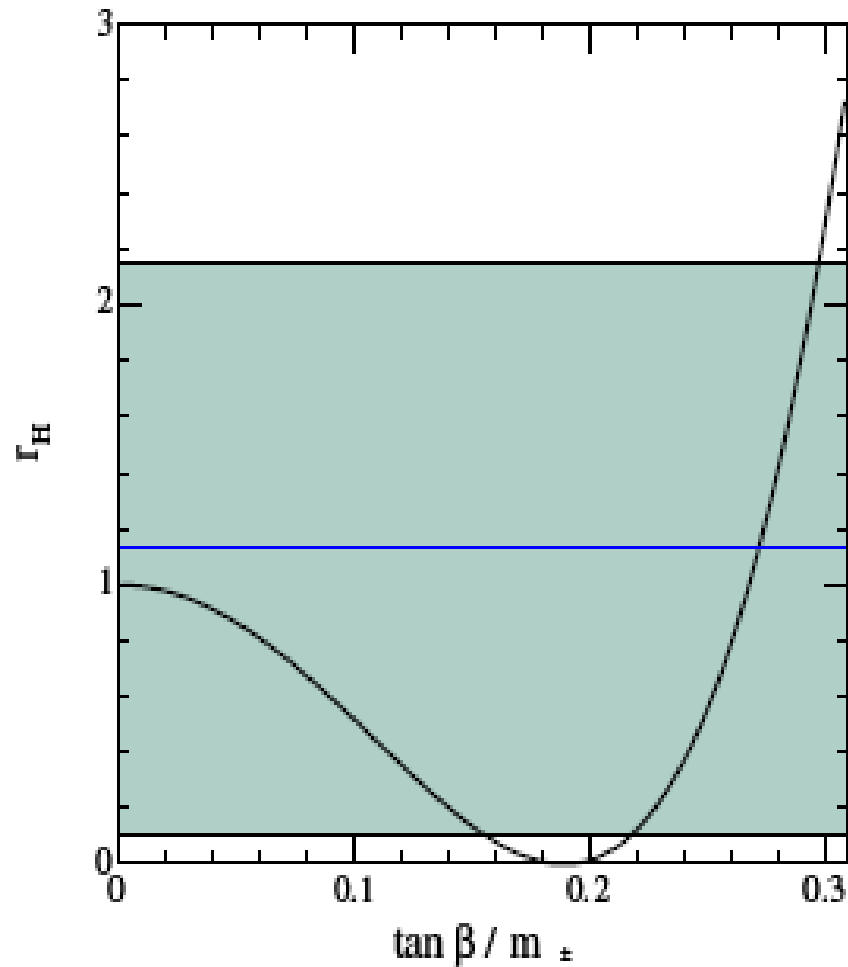
$$\text{HFAG } \text{BF}(B \rightarrow \tau \nu) = 1.34 \pm 0.48 \cdot 10^{-4}$$

Consistent with SM.



$$B^+ \rightarrow \tau^+ \nu$$

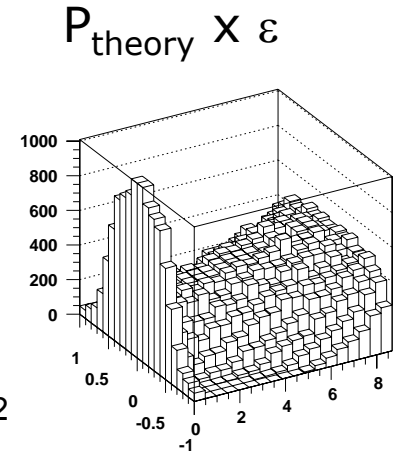
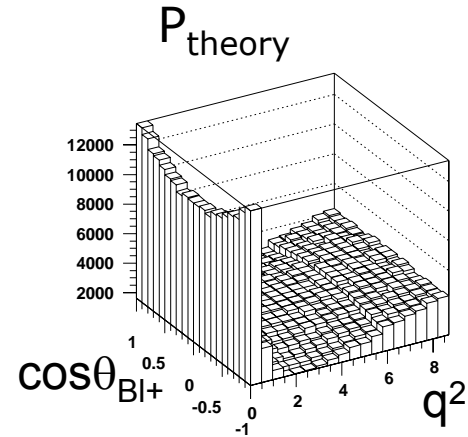
Belle result excludes (at $\tan \beta = 30$) $M(H^+) < 100$, $130 < M(H^+) < 190$ GeV



Probability Density Functions

$$\begin{aligned}
 \text{PDF} = & f_{\text{sig}} \times P_{\text{theory}}(A_7, A_9, A_{10}; q^2, \cos\theta) / N(A_7, A_9, A_{10}) \times \varepsilon(q^2, \cos\theta) \\
 & + (1 - f_{\text{sig}} - f_{\text{psi}} - f_{K^*hh}) \times P_{\text{dilepton}}(q^2, \cos\theta) \\
 & + f_{\text{psi}} \times P_{\text{psi}}(q^2, \cos\theta) \\
 & + f_{K^*hh} \times P_{K^*hh}(q^2, \cos\theta)
 \end{aligned}$$

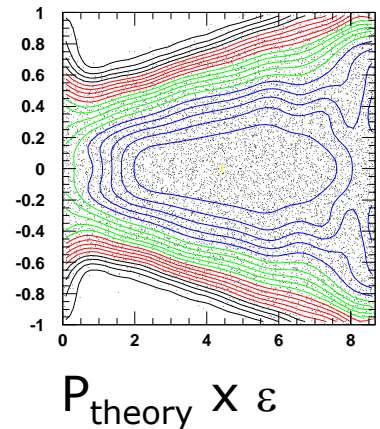
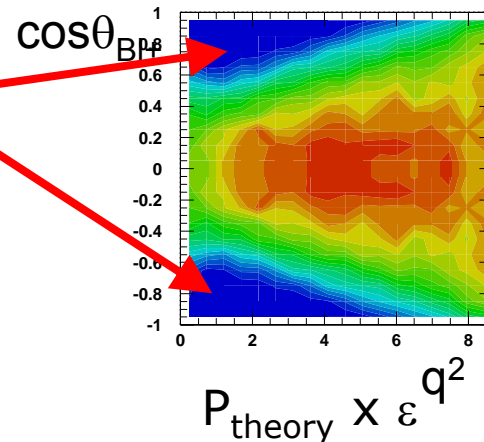
$K^*\mu\mu$ Signal PDF for q^2 below J/ψ veto window



We can not measure low q^2 and high $\cos\theta_{BI+}$ event since one of muon is low momentum.

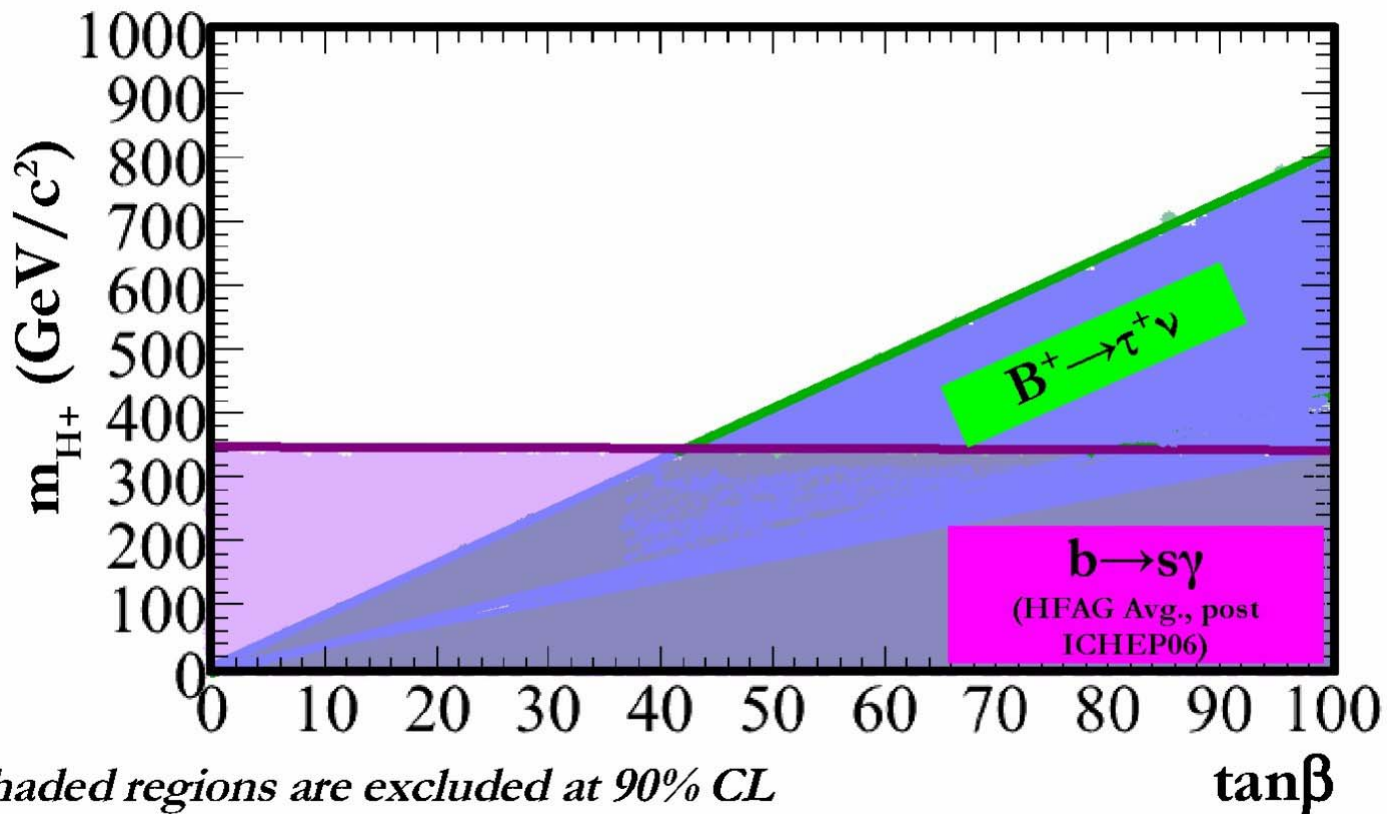
So we need muon detector for low momentum region.

(Acceptance for electron modes are a bit better.)



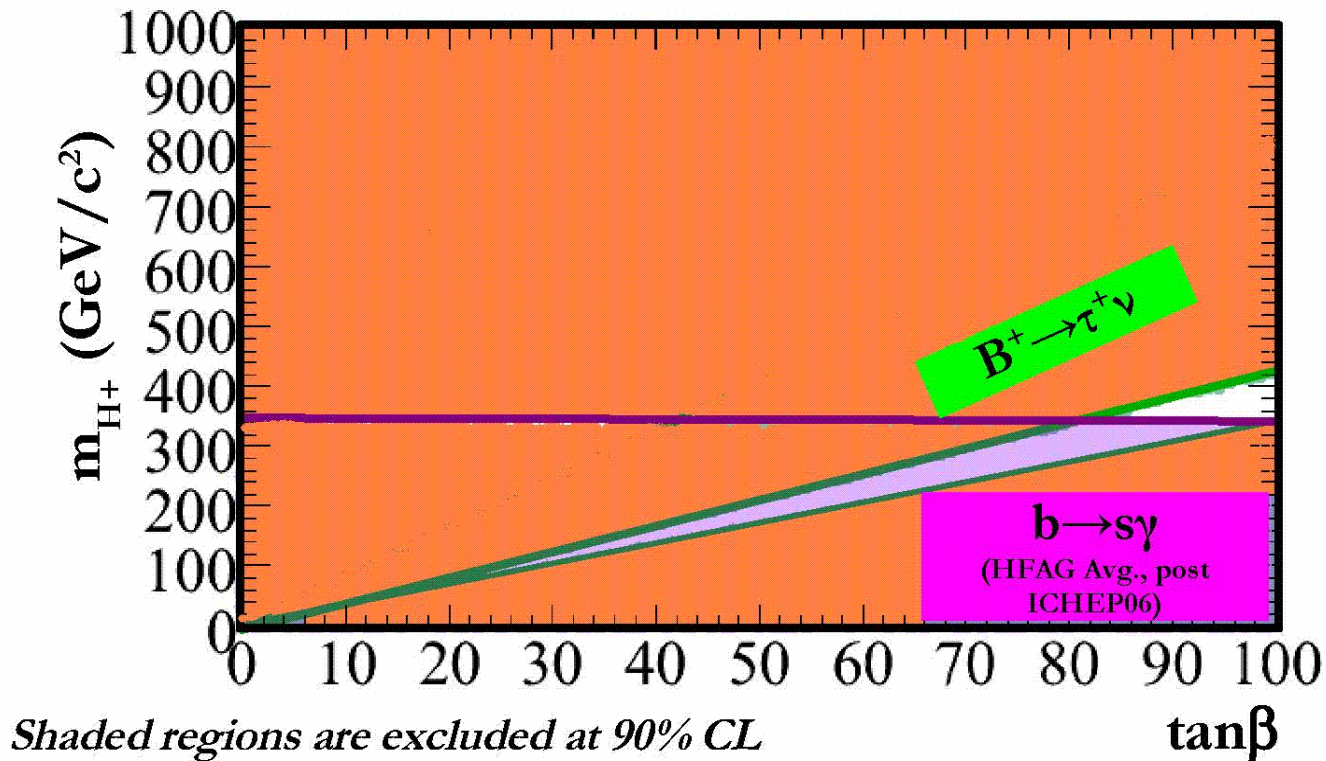
MK'07

Current Constraints on 2HDM



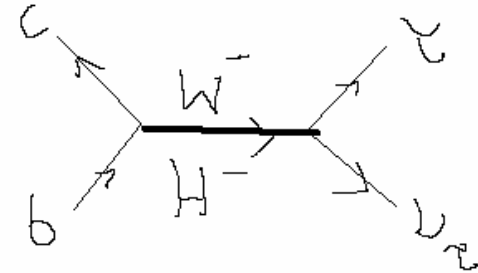
MK'07

Current Constraints on 2HDM



Compare with inclusive decay

$$B \rightarrow X_{c\tau\nu\tau}$$



The decay rate is different from the exclusive case

$$R_{bc\tau\nu} = 1 - 2m_\tau^2 \frac{\tan^2 \beta}{M_{H^\pm}^2} + m_b^2 m_\tau^2 \frac{\tan^4 \beta}{4M_{H^\pm}^4}$$

Corrections small, only small x solution

$$x = \tan^2 / M_{H^\pm}^2$$

$$R_{B\tau\nu} = \frac{\mathcal{B}(B_u \rightarrow \tau\nu)}{\mathcal{B}(B_u \rightarrow \tau\nu)^{\text{SM}}} = r_H = \left[1 - \tan^2 \beta \frac{m_B^2}{m_{H^\pm}^2} \right]^2$$

Why?

- In the exclusive case –
in SM helicity suppressed decay –
so W^- exchange prop. to m_τ/m_b
in H^+ exchange – there is also prop to m_τ
So, they are equally important and
 $(1-X)^2 = 1$ has two solutions $X=0$ and $X=2$
 $(X=x m_B^2)$
(H^+ gives 2 times bigger than W contribution)

Conclusion

- SM is great, in full agreement with data
- Nevertheless new physics effects in rare may be large, even larger than the SM – tree level prediction
- Be careful